

# *Real time dynamics of doped helium nanodroplets*

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Fausto CARGNONI,<sup>d</sup> Martí PI,<sup>c</sup>  
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<sup>b</sup> Universidad La Habana, Cuba

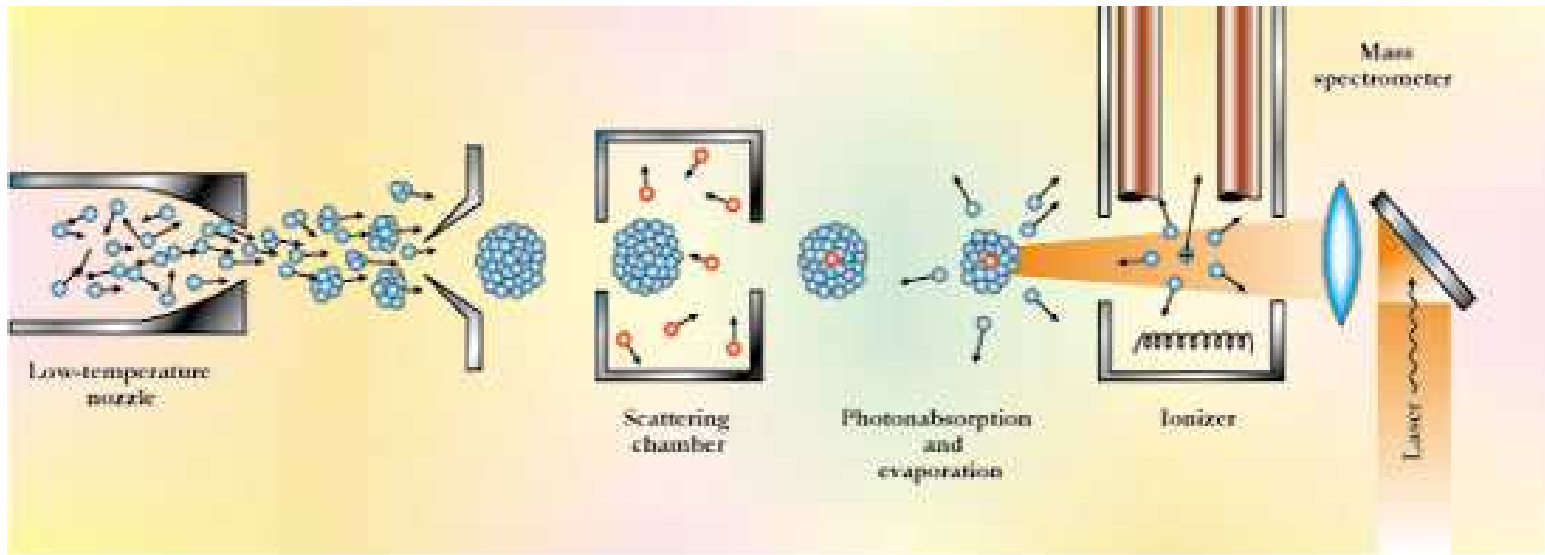
<sup>c</sup> Departament FQA and IN2UB, Universitat de Barcelona, Spain

<sup>d</sup> C N R - I S T M, Milano, Italy



# Helium nanodroplets:

- clusters of 500 to over  $10^8$  atoms

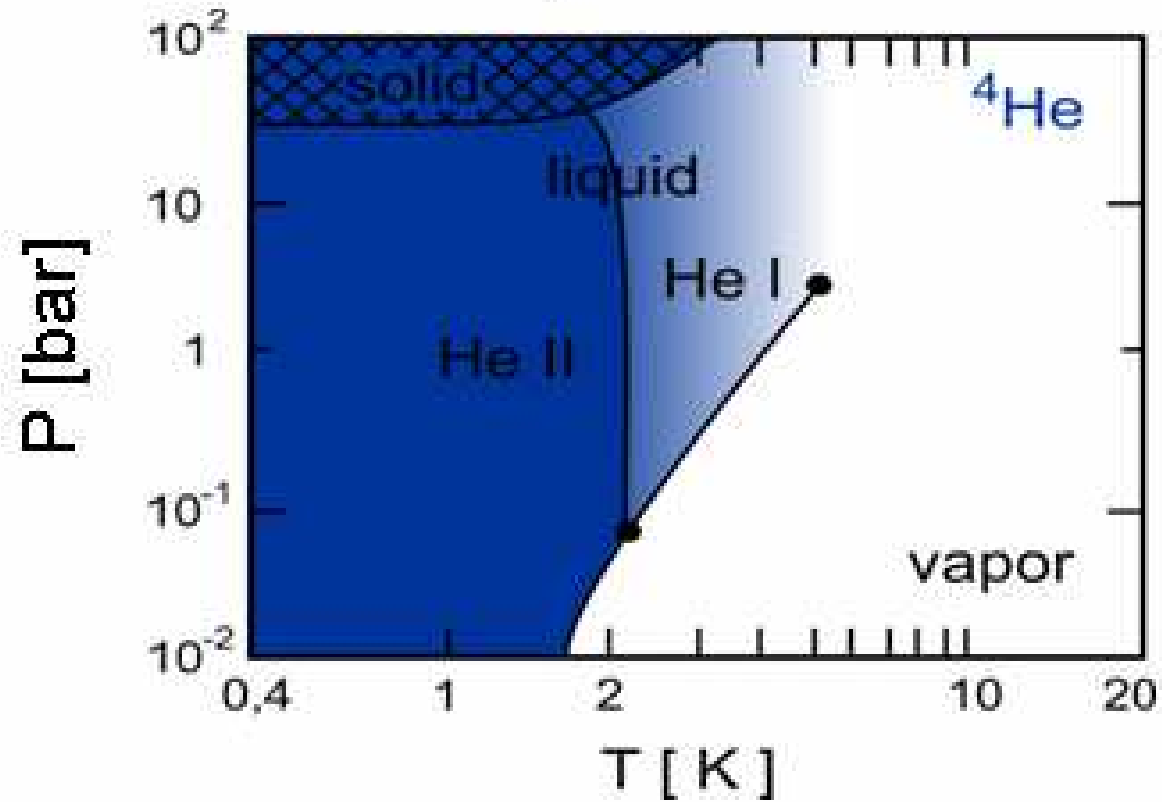


from Toennies, *Phys. Today* 2001

$T \simeq 0.4 \text{ K}$  from rotationally resolved spectra

# *Helium nanodroplets: Why study them?*

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***superfluidity*** in a finite size system?

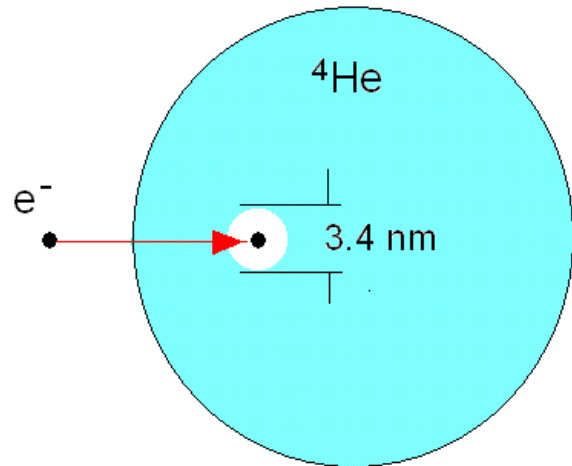
- “Frictionless” motion: molecular rotation, existence of a Landau critical velocity?
- Very high heat conductivity
- Zero-phonon lines separated from phonon wings
- Quantized vortices

# Quantum “solvent”

- Helium droplets will pick up “anything”
  - atoms, molecules...
  - $e^-$  metastable

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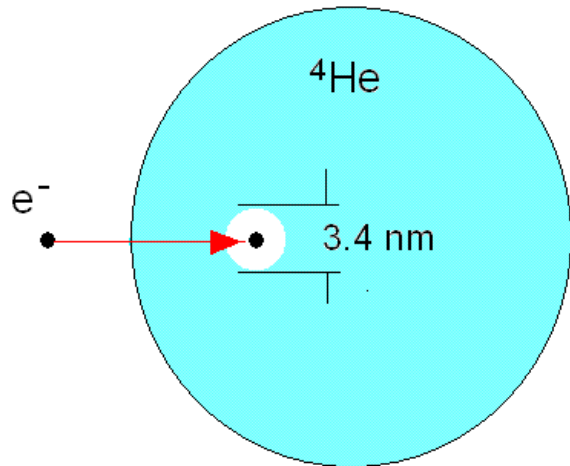
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He-electron interaction repulsive  
→ empty bubble around  $e^-$   
(*picture from G. Benedek*)

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He-electron interaction repulsive  
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- Location of the dopant:
  - M will sit near the center,
  - except alkalis, Ba, Sr, alkali clusters (which size?), A-Rg<sub>n</sub> to solvate them...



# *Nanolab's experiments*

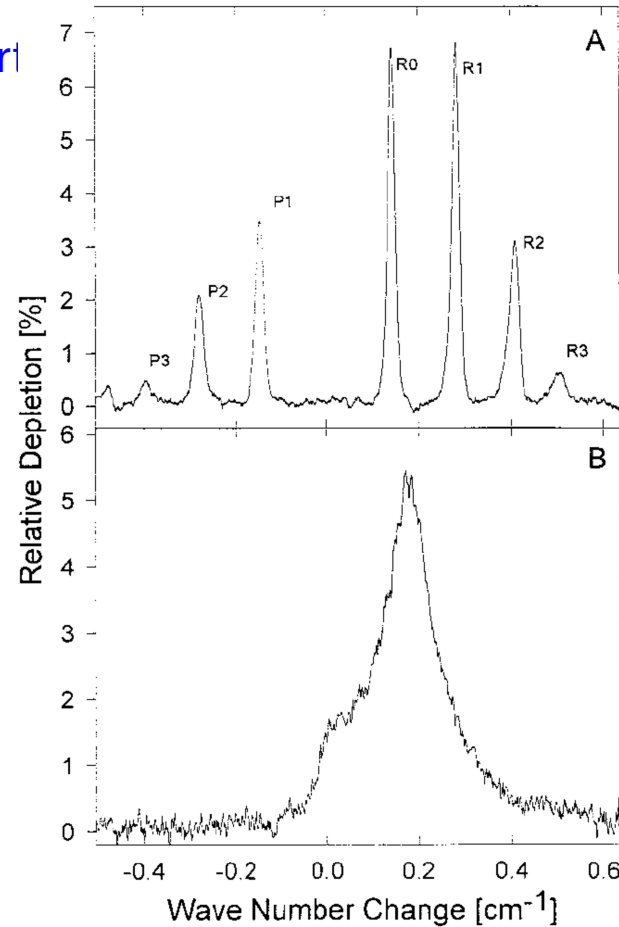
- $^4\text{He}_N$ : A “solvent” with very special properties:
    - extremely cold (0.4 K)
    - (super-)fluid
    - chemically inert
- high resolution spectroscopy

# Nanolab's experiments

## OCS rovibrational IR spectroscopy

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OCS rotational resolution in  $^4\text{He}$ !



*Grebenev, Toennies, Vilesov, Science 1998*

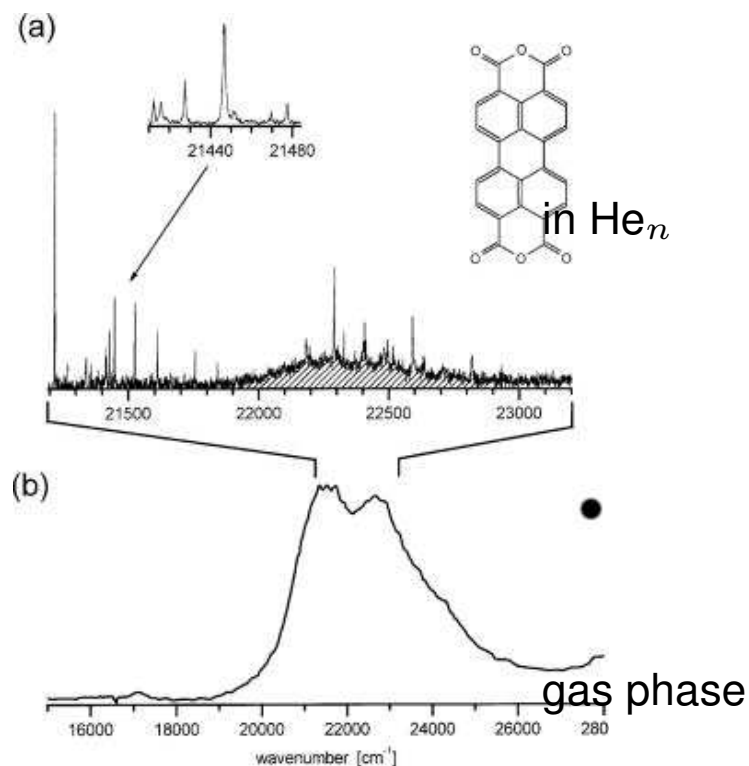
# Nanolab's experiments

## PTCDA electronic spectroscopy

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- (super-)fluid
- chemically inert

→ high resolution spectroscopy



(from Stienkemeier and Vilesov, *J.Chem.Phys.* 2001)

→ e.g., Gert Von Helden: *Catching proteins in liquid helium*

# Nanolab's experiments

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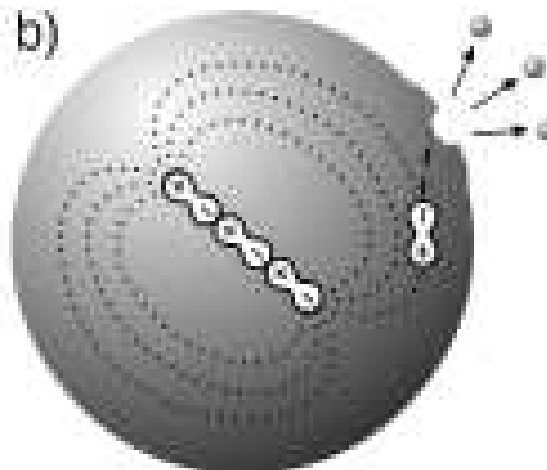
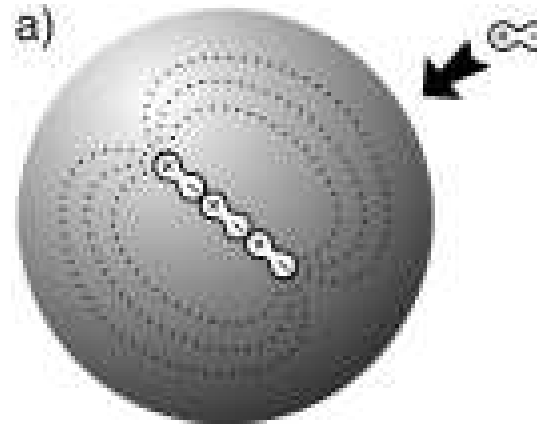
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→ high resolution spectroscopy

- very high heat conductivity → fast cooling

→ Exotic clustering

HCCCN clustering



(exp. Nauta, Moore and Miller 1999, from Toennies, *Angew.Chem.* 2004)

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→ Chemical reactions: highly reactive species, low barriers...

→ *e.g., Gary E. Douberly: Spectroscopy and reactions of hydrocarbon radicals in helium nanodroplets*

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**SUPERFLUID HELIUM**

## **The Ultimate Spectroscopic Matrix?**

K. K. Lehmann and G. Scoles

*Science 279, 2065 (1998)*

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But the influence of the superfluid helium environment is not well known...

# *Dynamics in the ground electronic state: Pickup and Clustering by superfluid $^4\text{He}$ droplets*

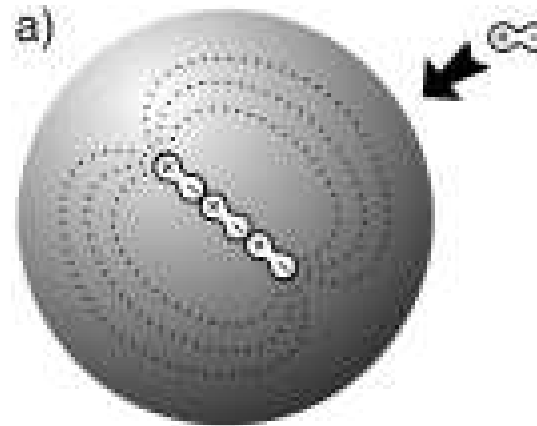
- Pickup by  $^4\text{He}_N$ :
  - extremely efficient:  $\sigma \sim \sigma_{\text{geometrical}}$



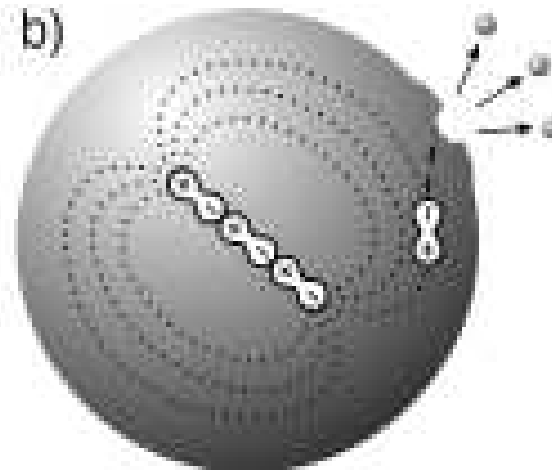
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HCCCN clustering



chain formation



driven by  
electrostatics

(exp. Nauta, Moore and Miller 1999, from Toennies, *Angew.Chem.* 2004)

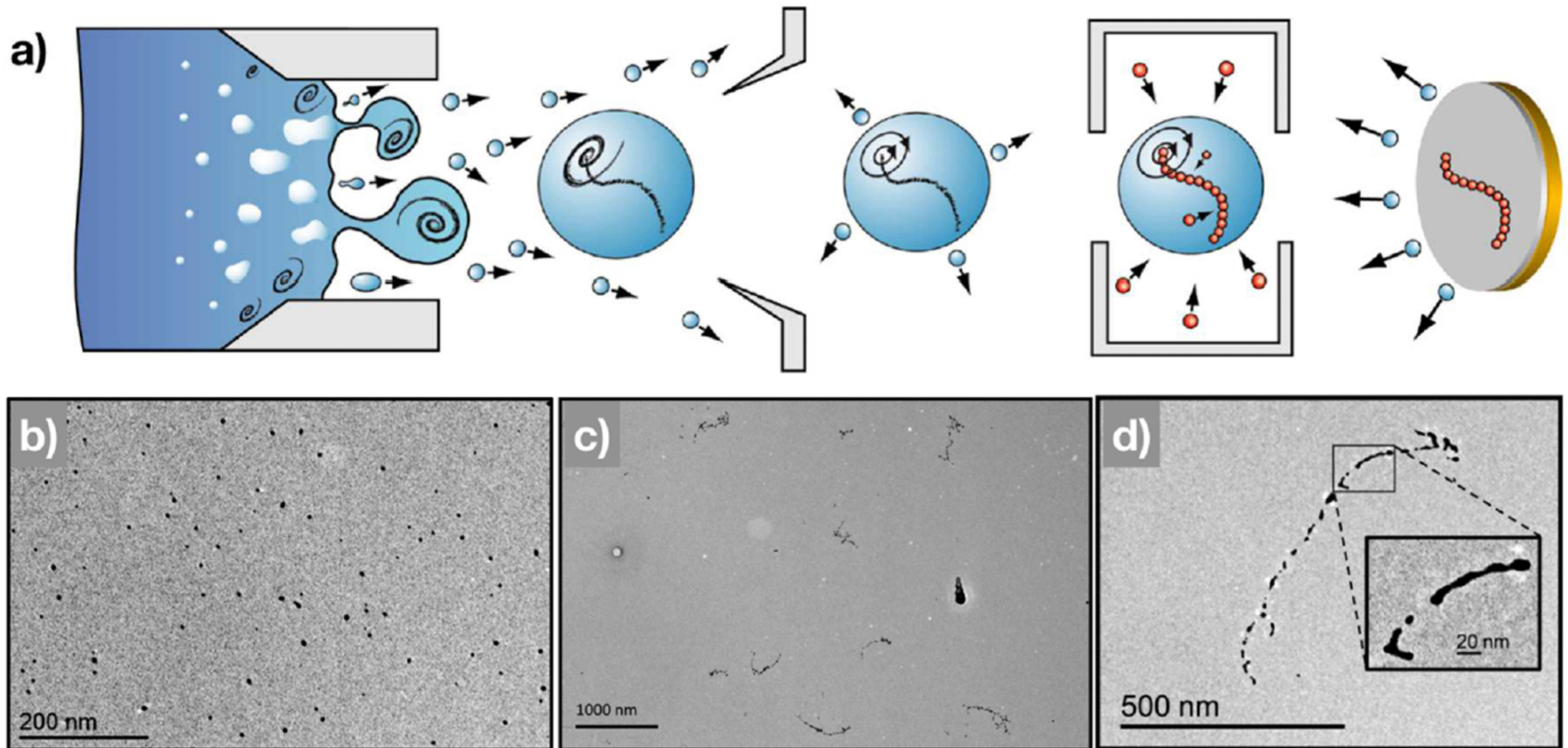
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  - visualize vortices with Xe atoms
  - filament-shaped clusters

# Dynamics in the ground electronic state: Pickup and Clustering by superfluid $^4\text{He}$ droplets



Ag-doped very large droplets  $\rightarrow$  filament-shaped clusters  
visualizing vortices

(Gessner and Vilesov, *Annu. Rev. Phys. Chem.* **70**, 173 (2019))

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  - Pickup + Clustering by  $\text{He}_N$  hosting Vortices
    - visualize vortices with Xe atoms
    - filament-shaped clusters
- Study pickup and clustering by superfluid helium droplets hosting or not vortices...

# Theoretical simulation: $^4\text{He}$ -DFT for $A@He_N$

Total energy of the system as a function of the  $^4\text{He}$  density  $\rho(\mathbf{r})$ :

$$E[\rho] = \int d\mathbf{r} \left\{ \frac{\hbar^2}{2m_{\text{He}}} |\nabla \sqrt{\rho(\mathbf{r})}|^2 + \mathcal{E}[\rho(\mathbf{r})] \right\} + \int d\mathbf{r} \rho(\mathbf{r}) V_A(|\mathbf{r}_{\text{Rb}} - \mathbf{r}|)$$

Orsay-Trento density functional : [F. Dalfovo, A. Lastrì, L. Pricapenko, S. Stringari & J. Treiner, PRB 1995; F. Ancilotto, M. Barranco, F. Caupin, R. Mayol & M. Pi, PRB 2005. ]

$$\begin{aligned} \mathcal{E}_c[\rho] = & \frac{1}{2} \int d\mathbf{r}' \rho(\mathbf{r}) V_{LJ}(|\mathbf{r} - \mathbf{r}'|) \rho(\mathbf{r}') + \frac{1}{2} c_2 \rho(\mathbf{r}) [\bar{\rho}(\mathbf{r})]^2 + \frac{1}{3} c_3 \rho(\mathbf{r}) [\bar{\rho}(\mathbf{r})]^3 \\ & - \frac{\hbar^2}{4m} \alpha_s \int d\mathbf{r}' F(|\mathbf{r} - \mathbf{r}'|) [1 - \tilde{\rho}(\mathbf{r})/\rho_{0s}] \nabla \rho(\mathbf{r}) \cdot \nabla' \rho(\mathbf{r}') [1 - \tilde{\rho}(\mathbf{r}')/\rho_{0s}] \\ & - \frac{m}{4} \int d\mathbf{r}' V_J(|\mathbf{r} - \mathbf{r}'|) \rho(\mathbf{r}) \rho(\mathbf{r}') [\mathbf{v}(\mathbf{r}) - \mathbf{v}(\mathbf{r}')]^2 + C \rho(\mathbf{r}) \{1 + \tanh[\beta (\rho(\mathbf{r}) - \rho_m)]\} \end{aligned}$$

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**Statics:** Minimize  $E[\rho]$  with respect to  $\Psi(\mathbf{r}) = \sqrt{\rho(\mathbf{r})}$ :  $\rightarrow$  Euler-Lagrange equation:

$$\left\{ -\frac{\hbar^2}{2m_{\text{He}}} \nabla^2 + \frac{\delta \mathcal{E}}{\delta \rho} + V_A(|\mathbf{r}_A - \mathbf{r}|) \right\} \Psi(\mathbf{r}) = \mu \Psi(\mathbf{r}), \quad \Psi(\mathbf{r}) = \sqrt{\rho(\mathbf{r})}$$

Sum of pairwise interactions with  $V_A$  from the literature (Aziz)

Imaginary time propagation on a 3-D grid  $\rightarrow \rho_0(\mathbf{r}) = \Psi_0^2(\mathbf{r})$

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$\Psi_{\text{He}}(\mathbf{r}, t)$  coupled to classical  $\mathbf{r}_A$  (mean field)

$$\begin{cases} i\hbar \frac{\partial}{\partial t} \Psi_{\text{He}} &= \left[ -\frac{\hbar^2}{2m_{\text{He}}} \nabla^2 + \frac{\delta \mathcal{E}}{\delta \rho} + V_A(\mathbf{r} - \mathbf{r}_A) \right] \Psi_{\text{He}} \quad \rightarrow \quad \rho_{\text{He}}(\mathbf{r}, t) = |\Psi_{\text{He}}(\mathbf{r}, t)|^2 \\ m_A \ddot{\mathbf{r}}_A &= -\nabla_{\mathbf{r}_A} \left[ \int d\mathbf{r} \rho(\mathbf{r}) V_\lambda(\mathbf{r} - \mathbf{r}_A) \right] = -\int d\mathbf{r} \nabla \rho(\mathbf{r}) V_A(\mathbf{r} - \mathbf{r}_A) \end{cases}$$



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With vortices: work in the co-rotating frame  $\omega$ :  $\{\mathcal{H}[\rho] - \omega \hat{L}_Z\} \Psi(\mathbf{r}) = \mu \Psi(\mathbf{r})$

*He-TDDFT simulation:  
Xe  $200 \text{ m s}^{-1}$  collision with He<sub>1000</sub>  
(thermal energy  $\leftrightarrow 240 \text{ m s}^{-1}$ )*

$$b = 0$$

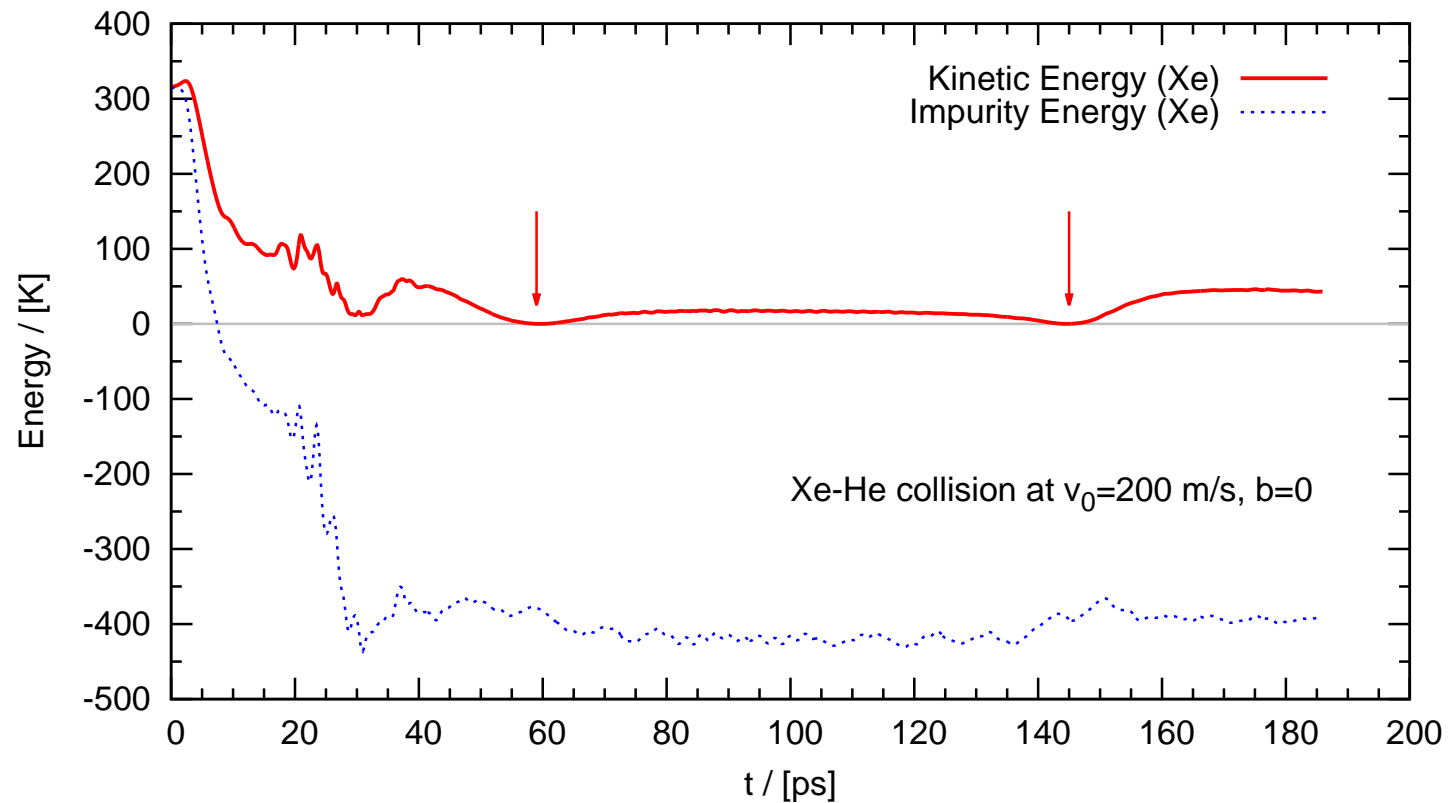
$$b = 22.2 \text{ \AA}$$

*He-TDDFT simulation:  
Xe  $200 \text{ m s}^{-1}$  collision with  $\text{He}_{1000}$   
(thermal energy  $\leftrightarrow 240 \text{ m s}^{-1}$ )*

If  $b >$  Droplet radius (22 Å), no capture

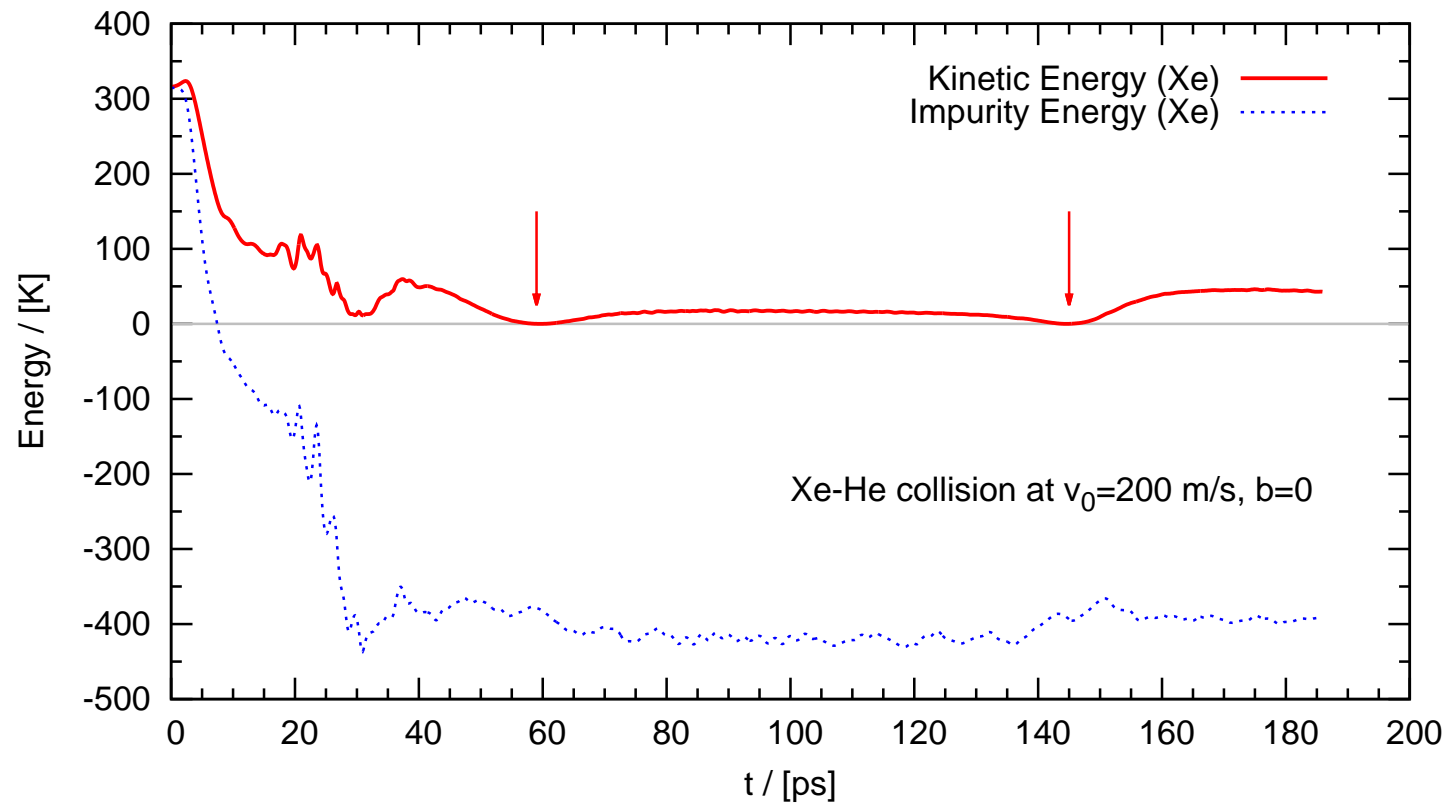
# He-TDDFT simulation: *Xe* 200 m s<sup>-1</sup> collision with He<sub>1000</sub>

(thermal energy  $\leftrightarrow$  240 m s<sup>-1</sup>)



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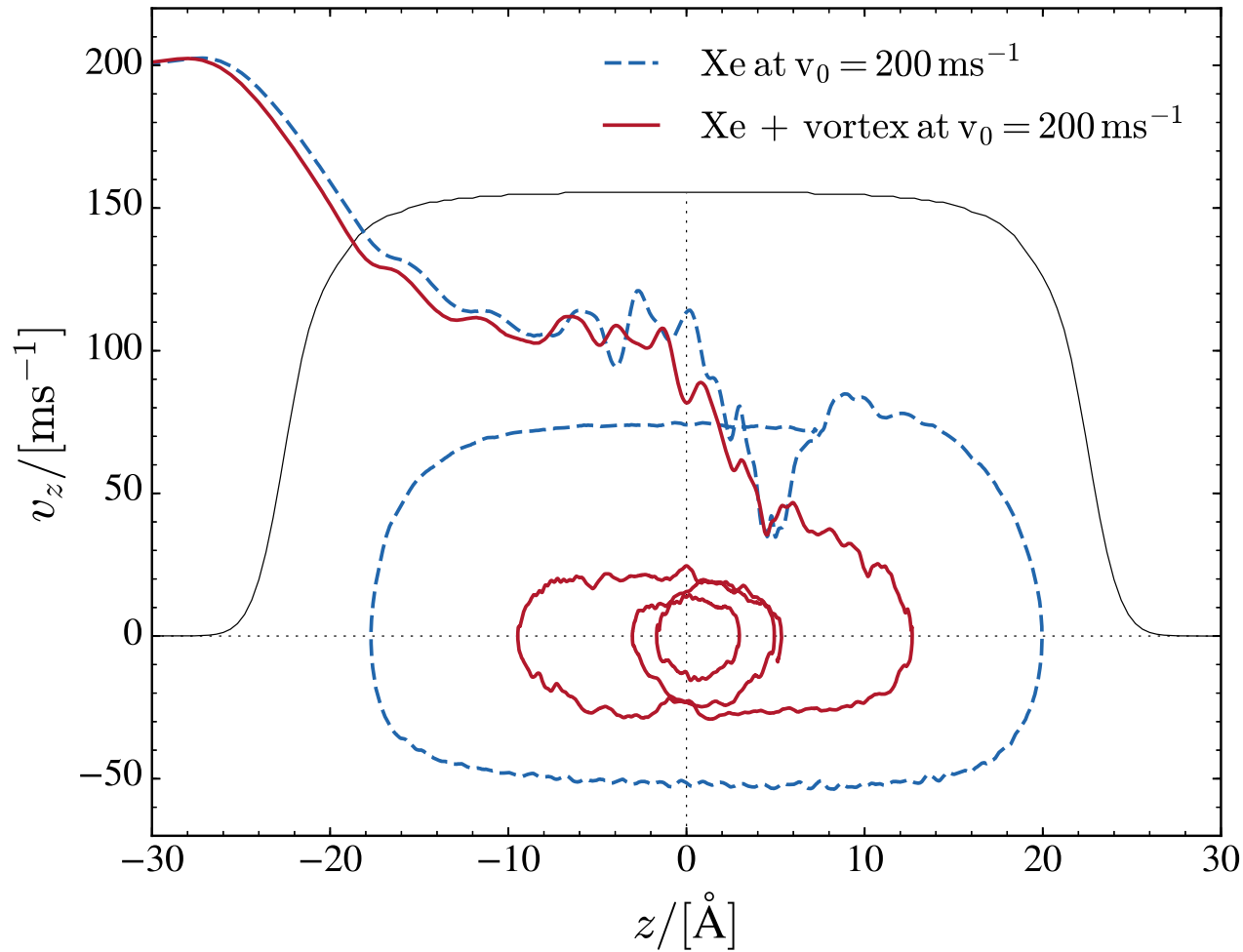
Capture cross section  $\simeq$  geometrical cross section

*Coppens, Leal, Barranco, Halberstadt, Pi, JLTP 187, 439 (2017);*

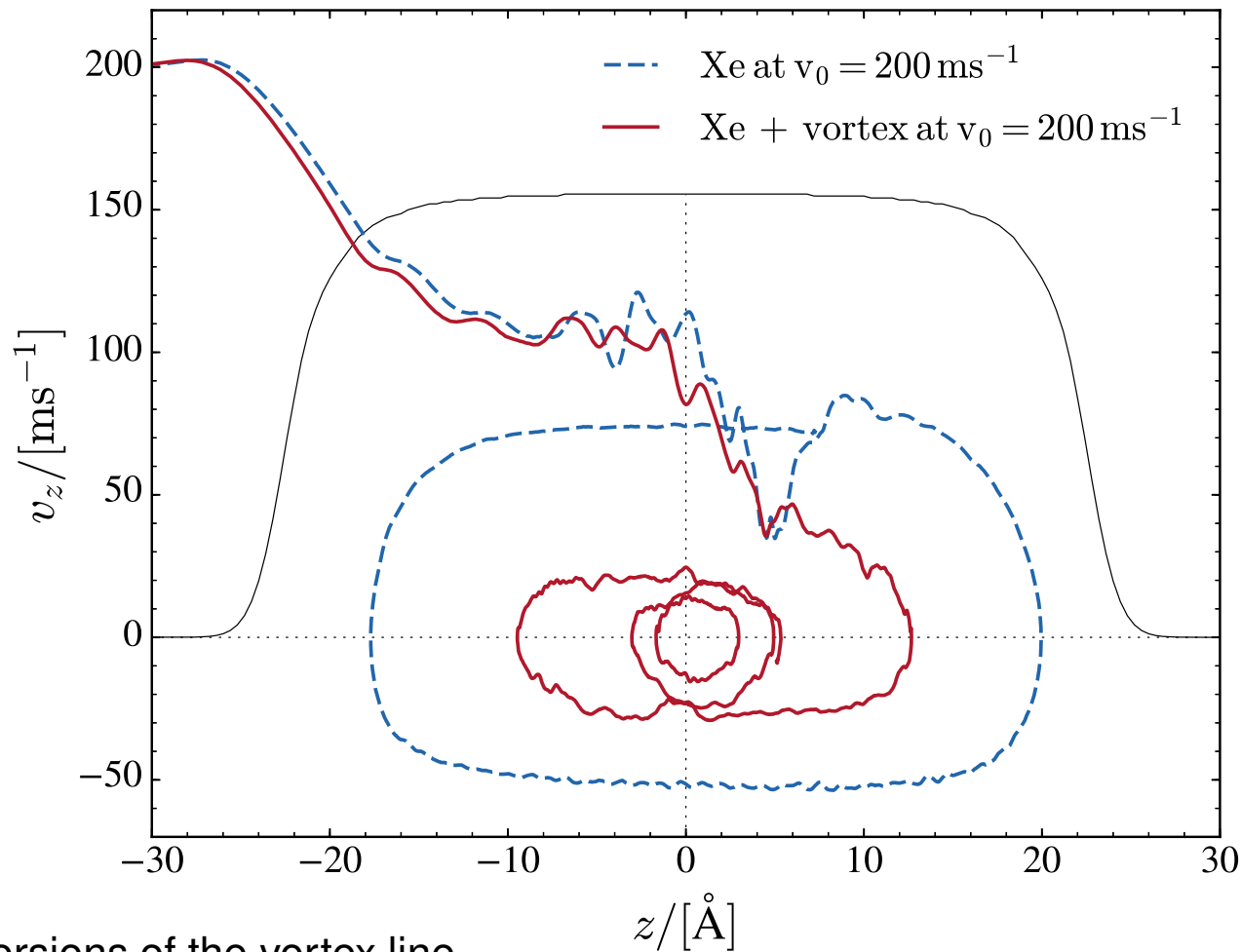
*With a vortex*  
*Xe  $200 \text{ m s}^{-1}$  collision with He<sub>1000</sub>*  
*(thermal energy  $\leftrightarrow 240 \text{ m s}^{-1}$ )*

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*With a vortex*  
*Xe  $200 \text{ m s}^{-1}$  collision with  $\text{He}_{1000}$*



Strong distortions of the vortex line

Xe ends up bound to the vortex

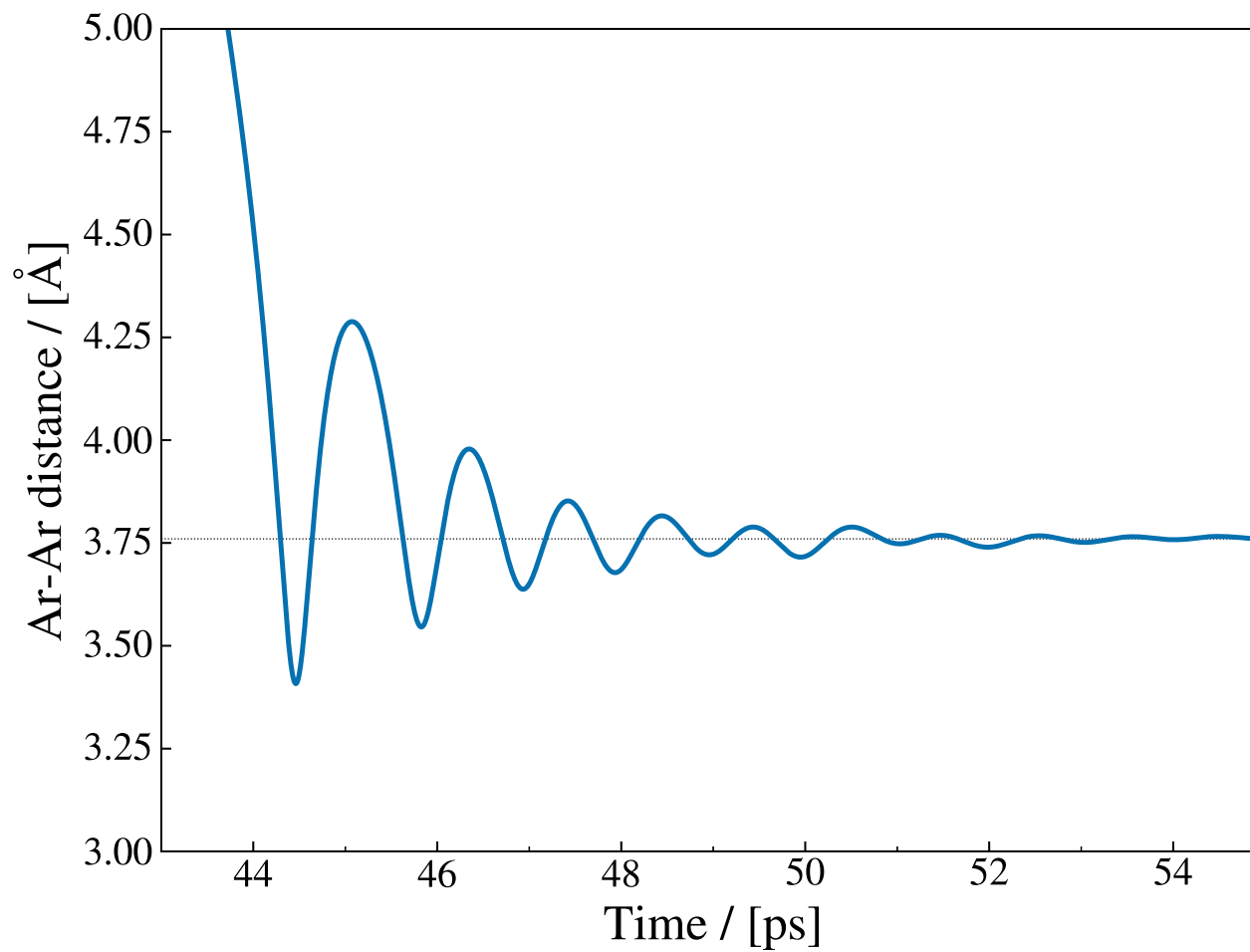
*Coppens et al., JLTP **187**, 439 (2017); Coppens et al., PCCP **19**, 24805 (2017)*



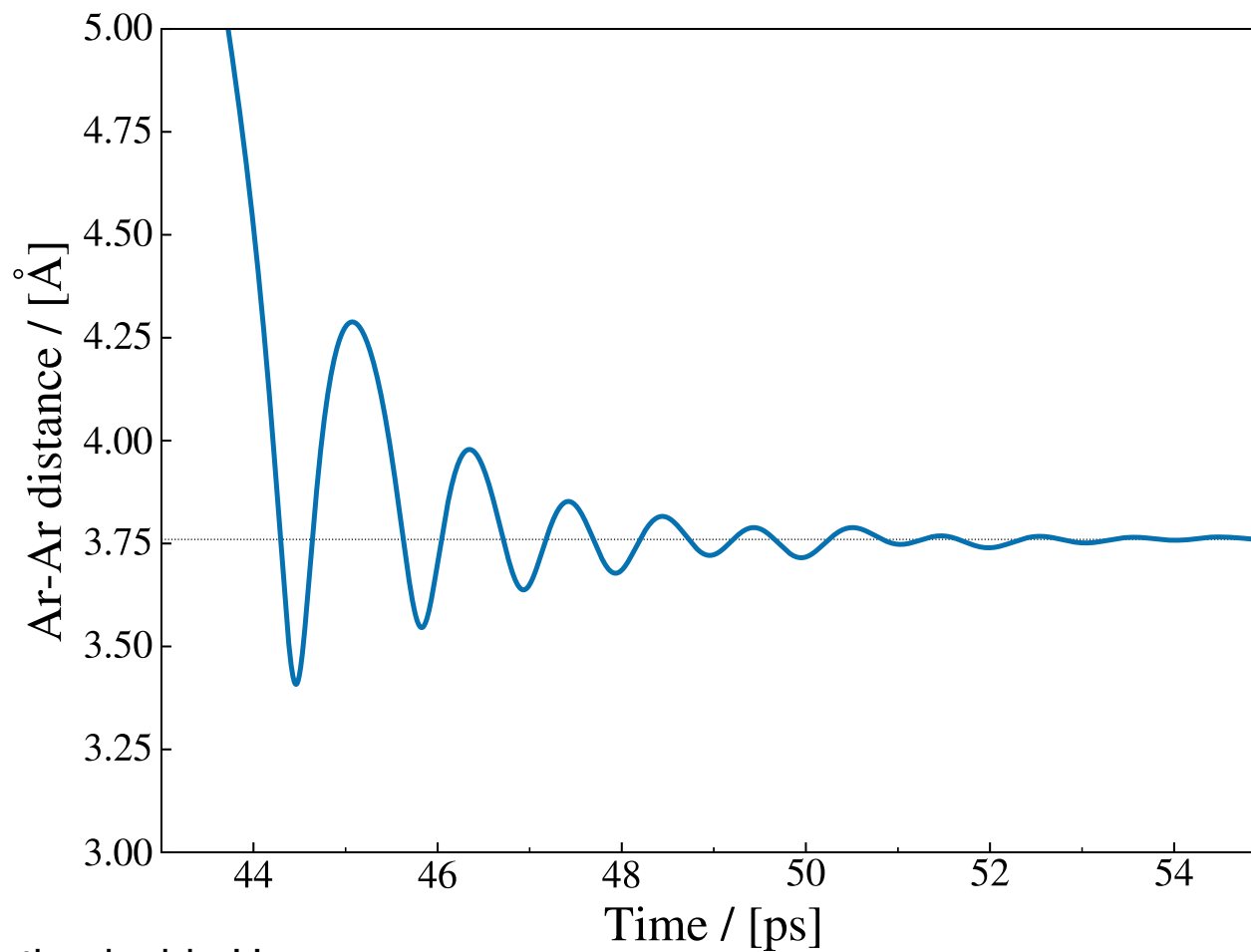
*Clustering: Ar<sub>2</sub>*  
*2 Ar, 360 m s<sup>-1</sup> collision with He<sub>5000</sub>*

$$b = 0$$

# Clustering: Ar<sub>2</sub>



## Clustering: Ar<sub>2</sub>



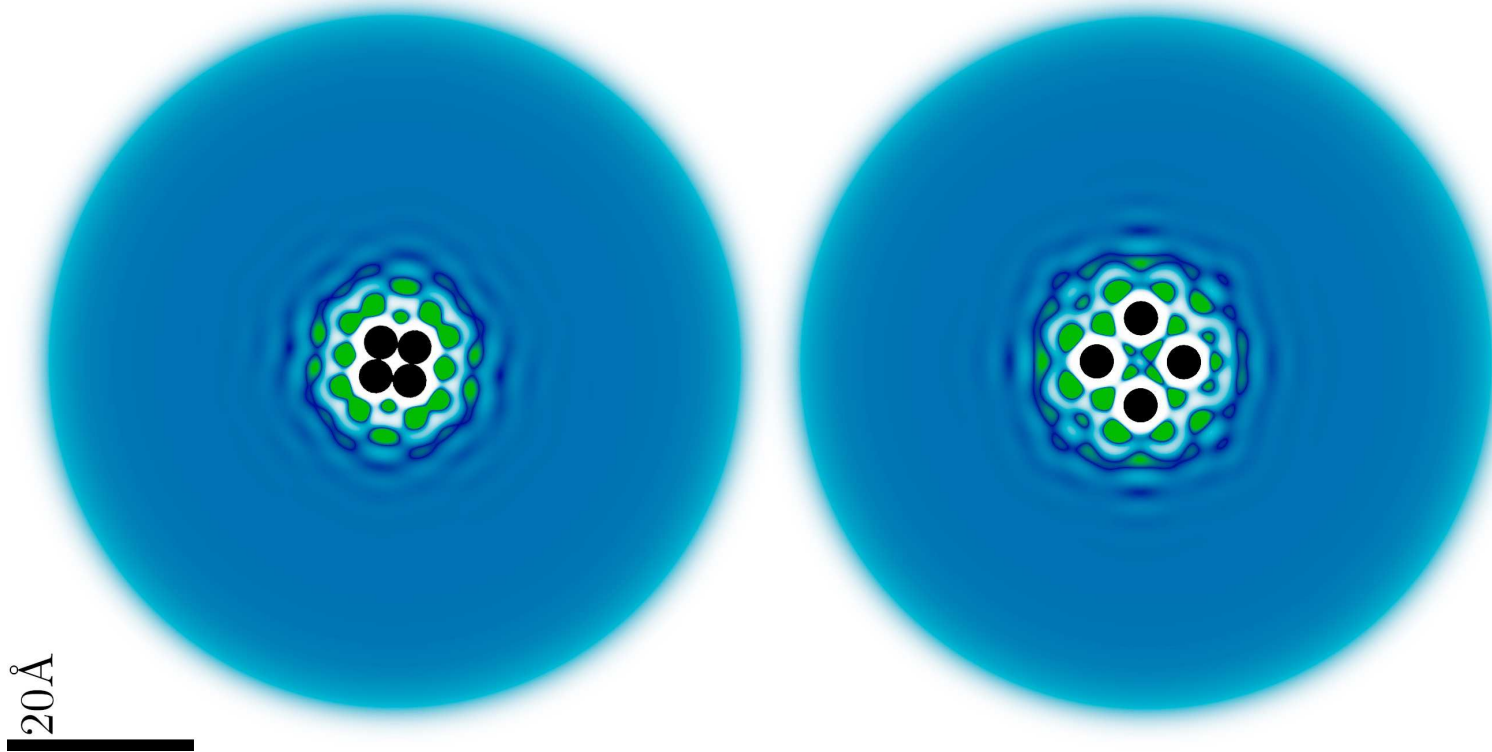
Dimer formation inside He<sub>5000</sub>

*Coppens et al., Phys.Chem.Chem.Phys.* **21**, 17423 (2019)

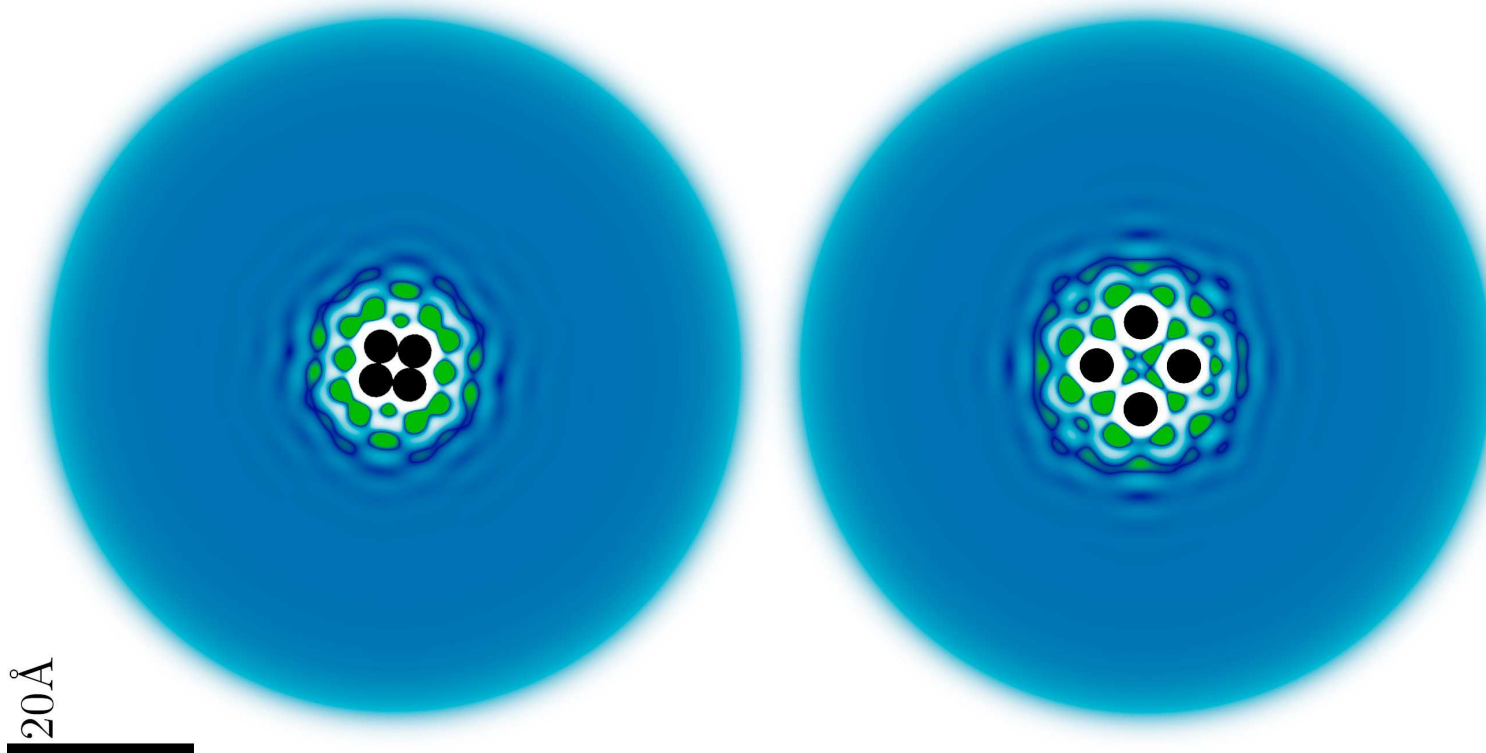
*Clustering: Ar<sub>6</sub>*  
*6 Ar 360 m s<sup>-1</sup> collision with He<sub>5000</sub>*

$$b = 0$$

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*Clustering: Ar<sub>6</sub>*  
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Ar<sub>6</sub> formation inside He<sub>5000</sub> hindered by He shell structure: “dilute” Ar<sub>6</sub>  
*Coppens et al., Phys.Chem.Chem.Phys.* **21**, 17423 (2019)

# *Clustering: $A_{r_6}$ with 6 vortices*

$$b = 0$$

# Clustering: $Ar_6$ with 6 vortices

$$b = 0$$

A dilute  $Ar_6$  cluster is formed, bound to the central cluster line  
*Coppens et al., Phys.Chem.Chem.Phys.* **21**, 17423 (2019)



## *Conclusion for cluster formation*

- Capture cross section  $\simeq$  geometrical cross section
- Capture and cluster formation release a lot of energy in the droplet  
→ necessity to describe the He droplet dynamics
- Clustering is well described by He-TDDFT (dimer formation)
- Clustering can be hindered by the He solvation structure
- Atoms are attracted to the vortex line(s) and end up bound to them
- Cluster formation along the vortex lines?

Would require longer time dynamics, with more Ar atoms and even larger droplets...

Also: Vary initial conditions

*Excited state dynamics:  
He-induced electronic transitions?*

He is chemically inert, at  $T \simeq 0.4$  K it looks “inoffensive”

# *Excited state dynamics: He-induced electronic transitions?*

He is chemically inert, at  $T \simeq 0.4$  K it looks “inoffensive”, and yet:

Laser-induced fluorescence of alkali atoms in He  
(*cold dense gas, superfluid liquid, solid, droplets*):

→ electronic relaxation

Also for molecules:  $\text{Cl}_2^*(\text{B})$

→ **study case:  $\text{Ak}^*$  at  $\text{He}_N$**

## *$Ak \rightarrow Ak^*$ Dynamics @ He droplets*

- many experiments: LIF, beam depletion, VMI... (Scoles, Stienkemeier, Ernst, Drabbels, Mudrich,...)
- LIF experiments: Ak sits in a “dimple” at the He droplet surface (shift/liquid He)
- Photoexcitation: leads to  $Ak^*$  or  $He_n Ak^*$  desorption [except for  $Rb^*(5\Pi_{1/2})$  and  $Cs^*(6\Pi_{1/2})$  low energy side]

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## **Theory**

- Pseudo-diatomic model (He<sub>N</sub>)—Ak remarkably valid for spectroscopy
- Ring Polymer MD for K\* He<sub>300</sub> (Takayanagi and Shiga, PCCP 2004): one of the first dynamics studies
- He-DFT, TDDFT (Ancilotto, Callegari, Hernando, Matteo, Leal, Barranco, Pi, Eloranta, Dalfovo,...)

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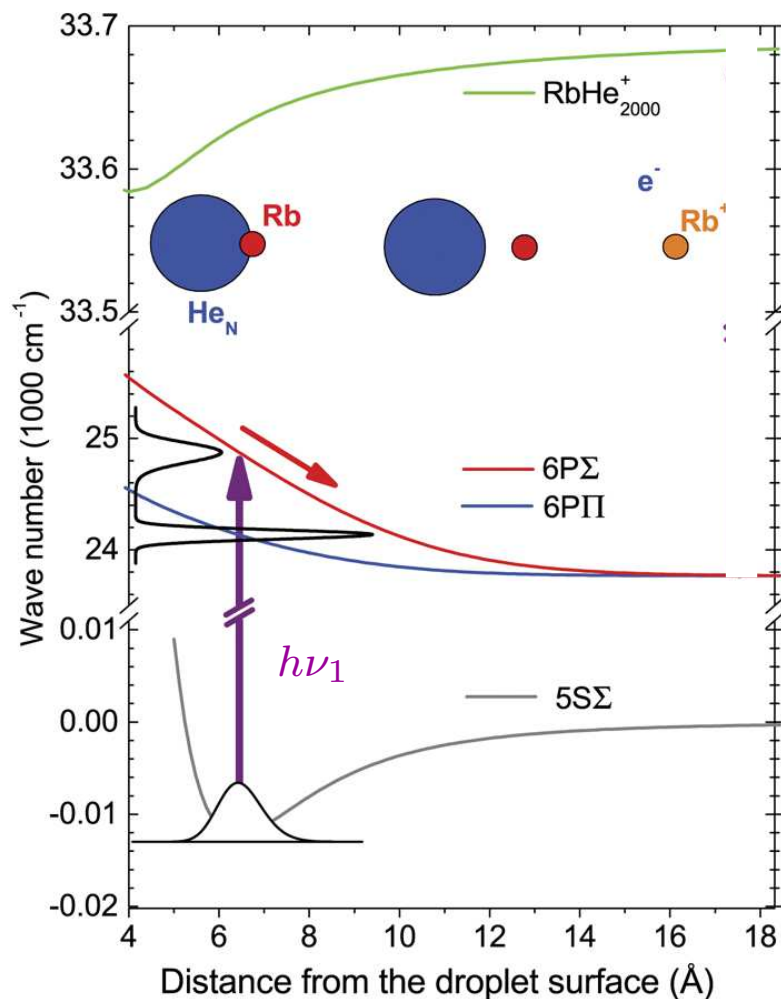
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**This work:** collaboration with Von Vangerow, Stienkemeier, Mudrich  
First direct time-dependent observation of the dynamics

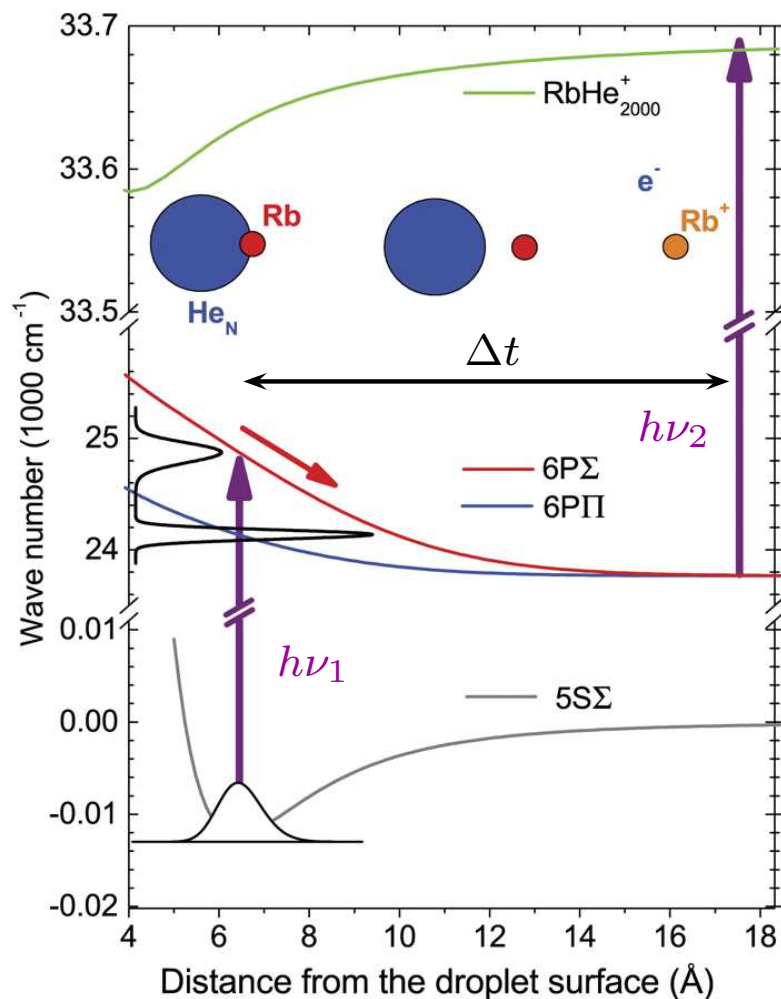
# Desorption dynamics of $Rb^*$ ( $5p, 6p$ )



Principle of the experiment

*J. von Vangerow, O. John, F. Stienkemeier, and M. Mudrich, JCP* **143**, 034302 (2015)

# Desorption dynamics of $Rb^*$ ( $5p, 6p$ )



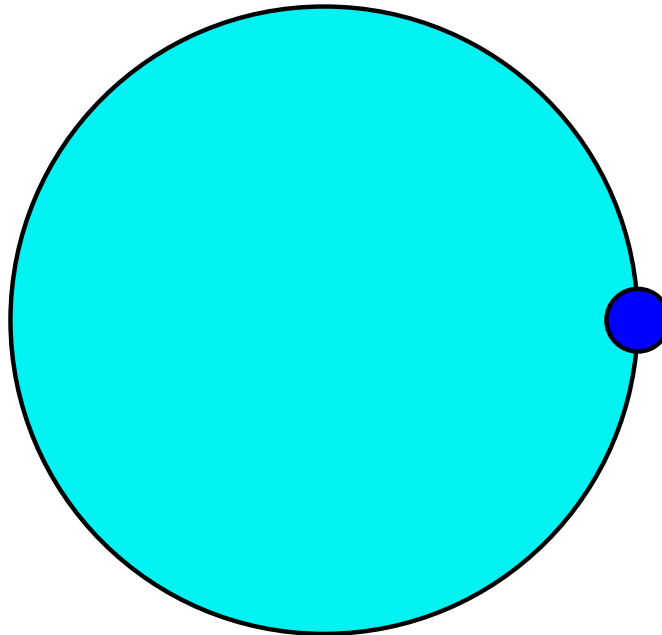
Principle of the experiment

*J. von Vangerow, O. John, F. Stienkemeier, and M. Mudrich, JCP **143**, 034302 (2015)*



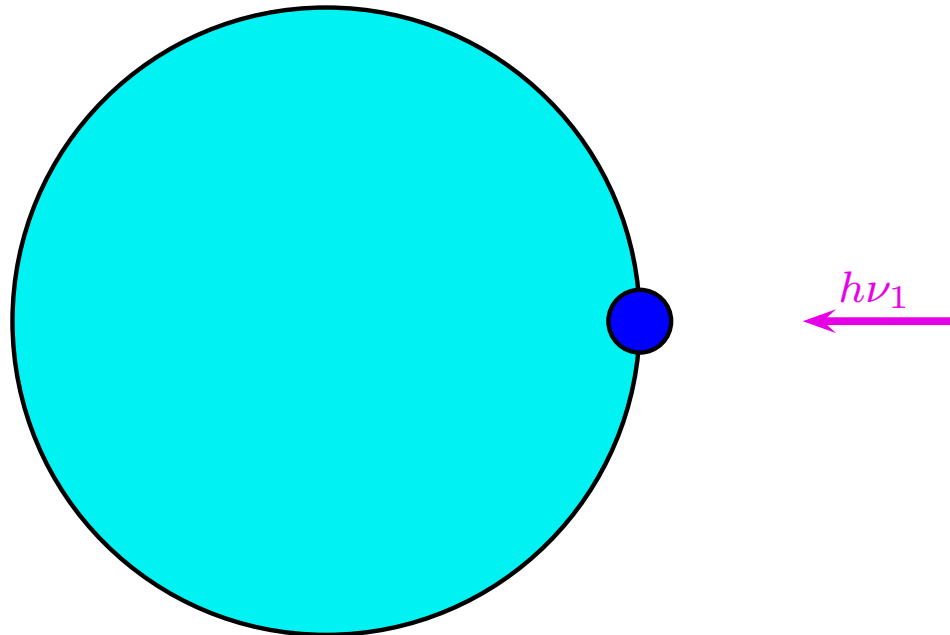
# *Principle of the experiment*

*initial state:  $Rb(5s, \Sigma)@He_N$*

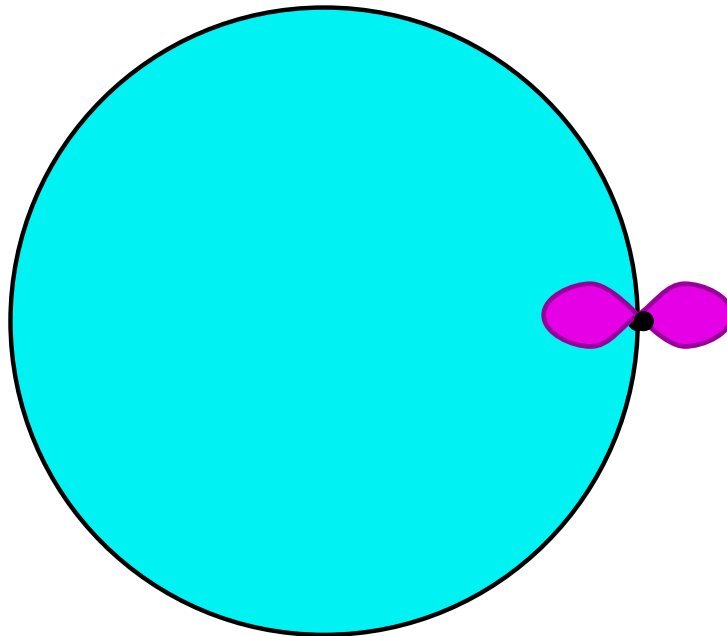


# *Principle of the experiment*

*laser 1 excitation*

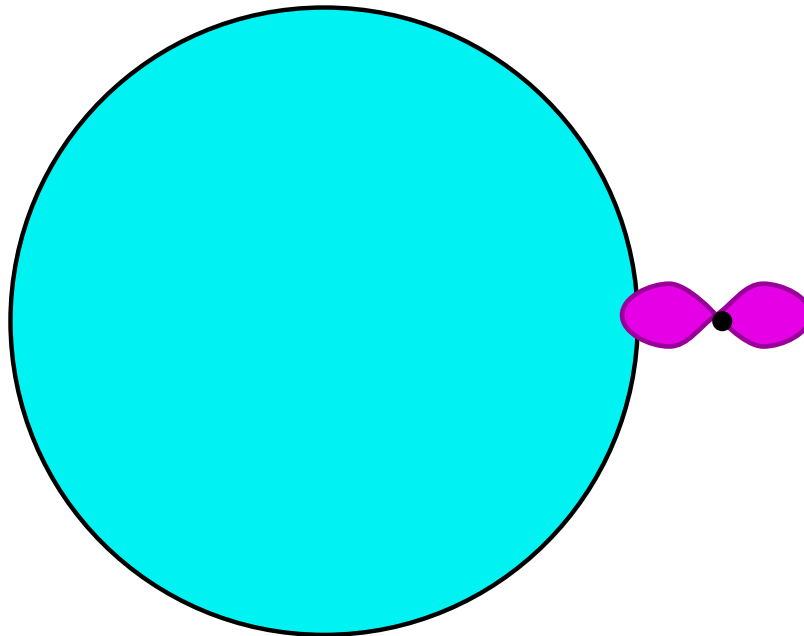


# Principle of the experiment



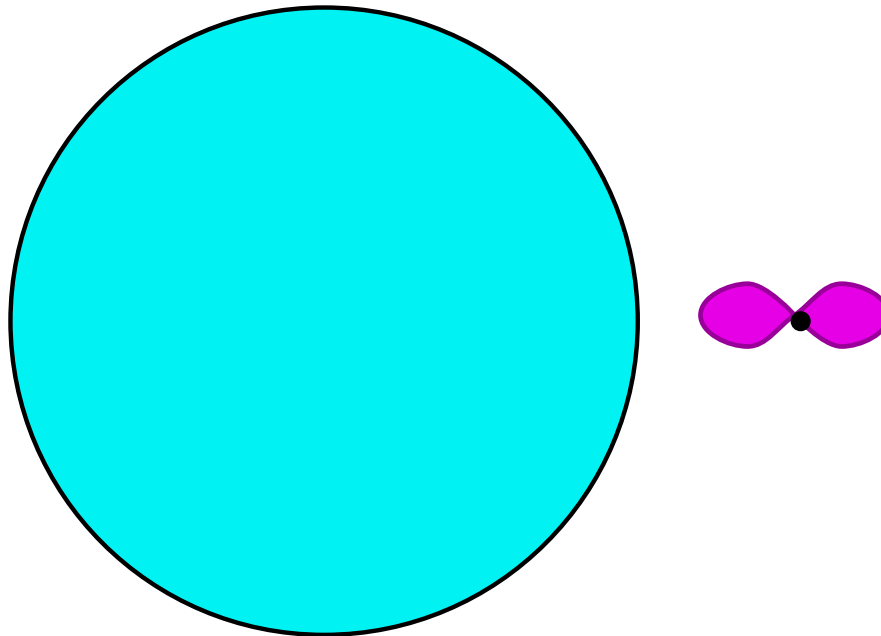
# *Principle of the experiment*

*ejecting  $Rb^*$  ( $np, \Sigma$ )*



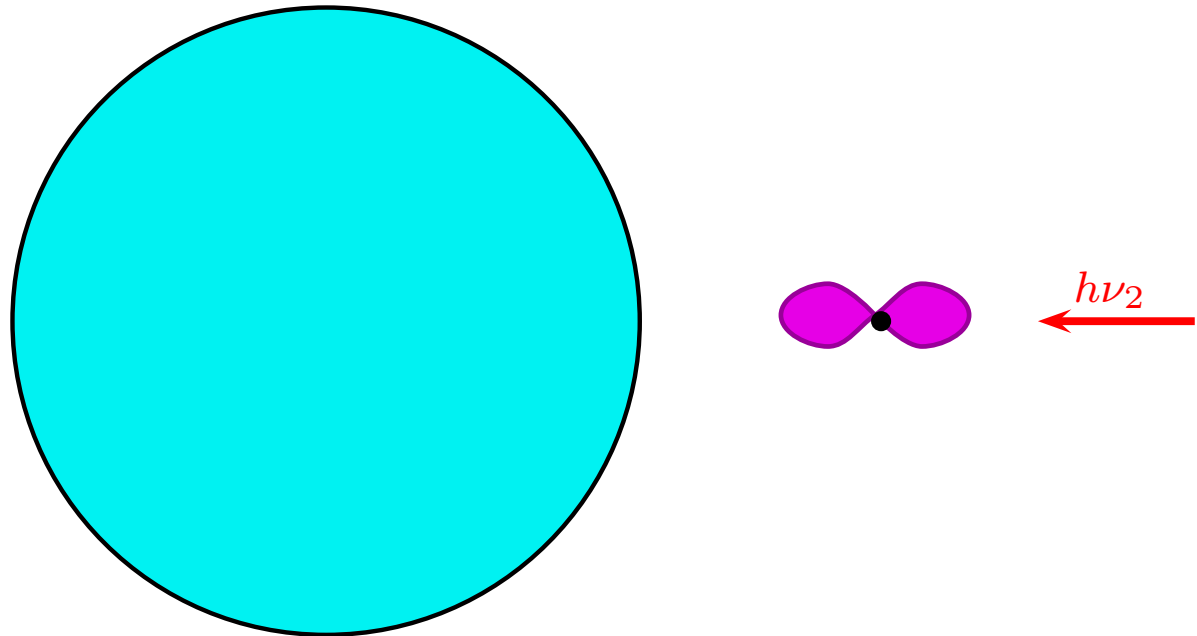
# *Principle of the experiment*

*ejecting  $Rb^*$  ( $np, \Sigma$ )*



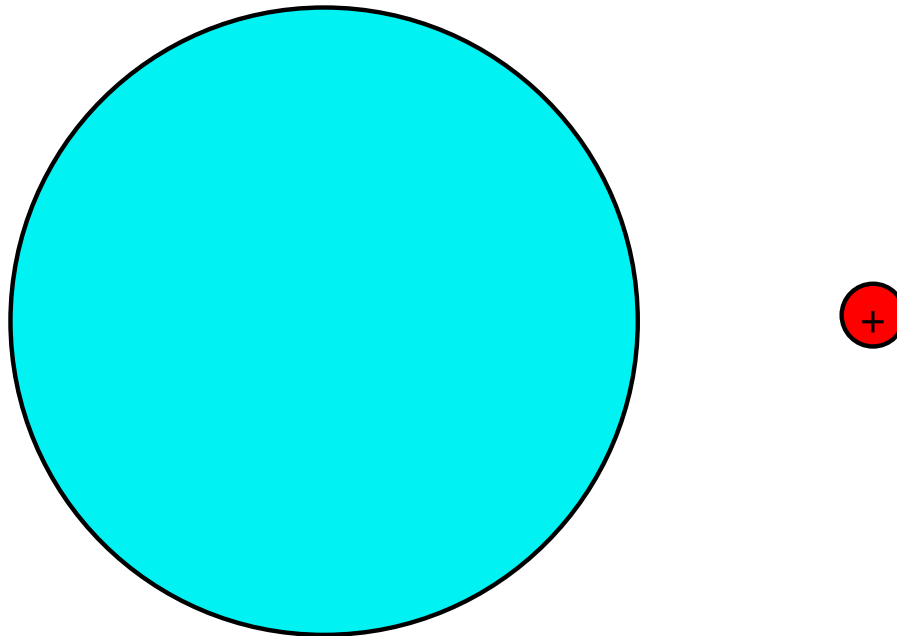
# Principle of the experiment

*laser 2 ionization*



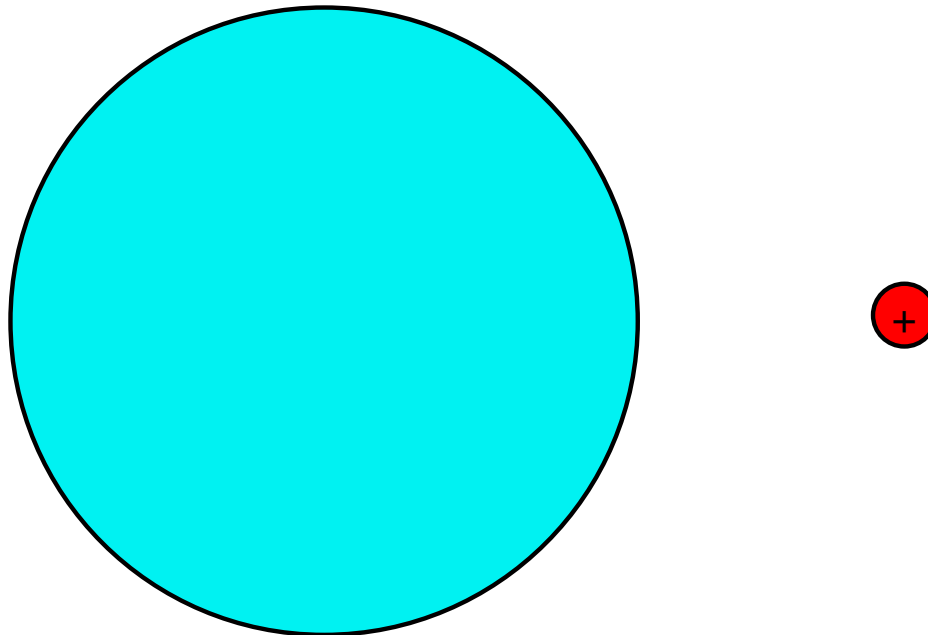
# *Principle of the experiment*

$Rb^+$



# *Principle of the experiment*

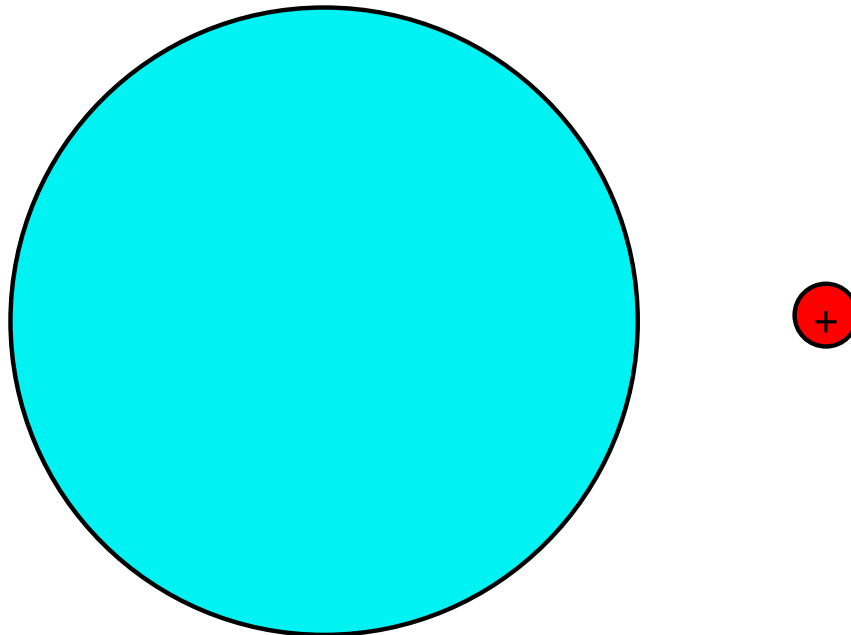
*slow down due to  $Rb^+$  -He attraction*





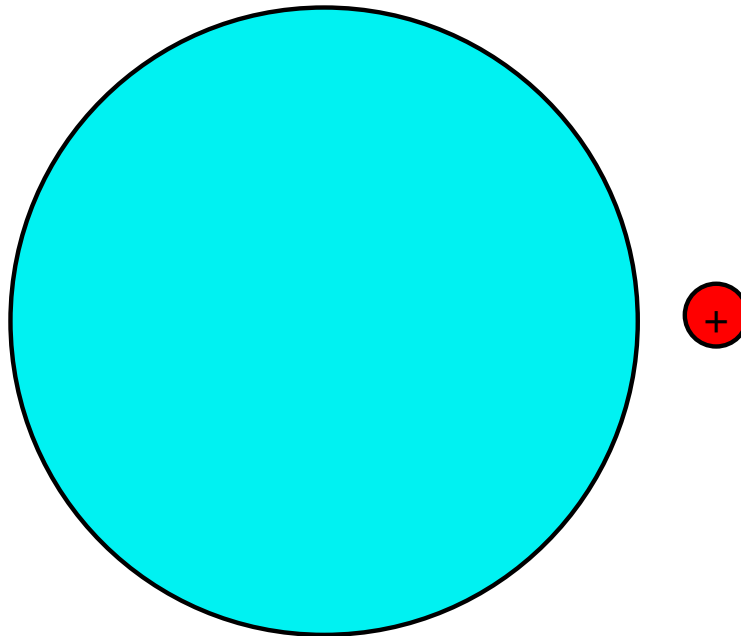
# *Principle of the experiment*

*Rb<sup>+</sup> possible turnaround and ...*



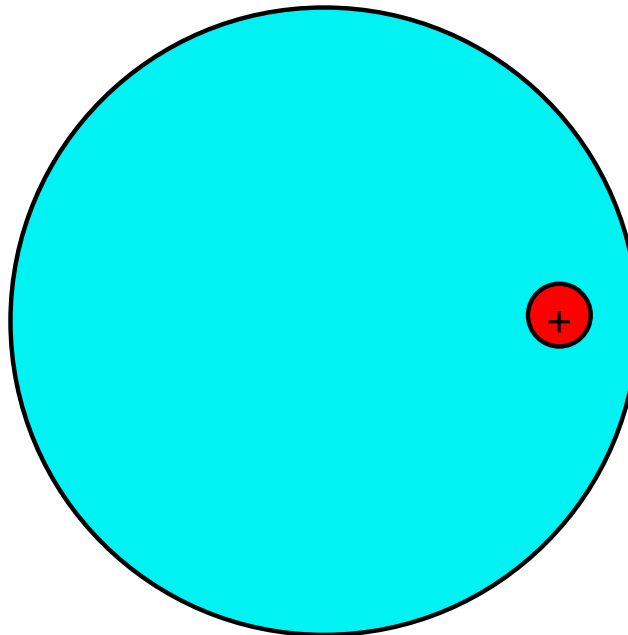
# *Principle of the experiment*

*Rb<sup>+</sup> possible turnaround and ...*



# *Principle of the experiment*

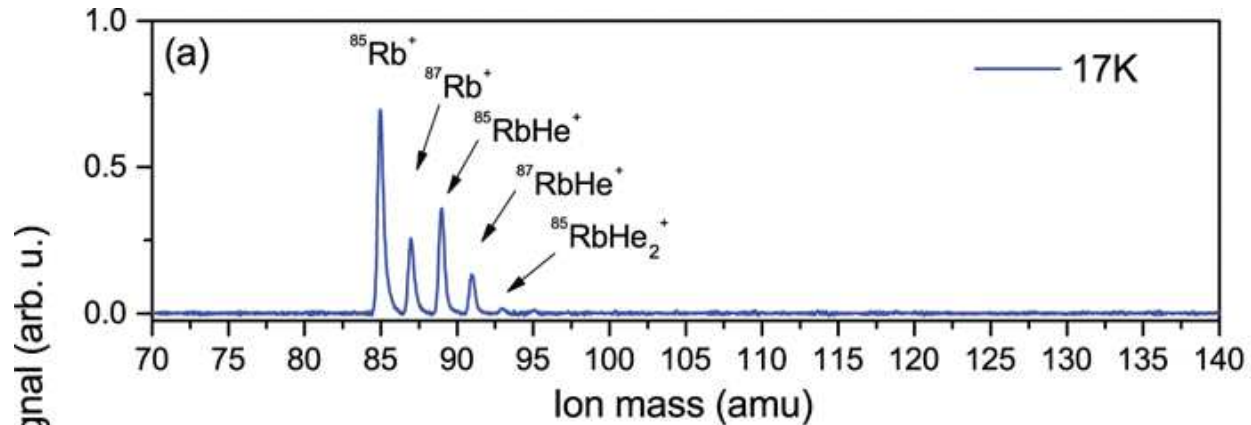
*Rb<sup>+</sup> possible turnaround and solvation*



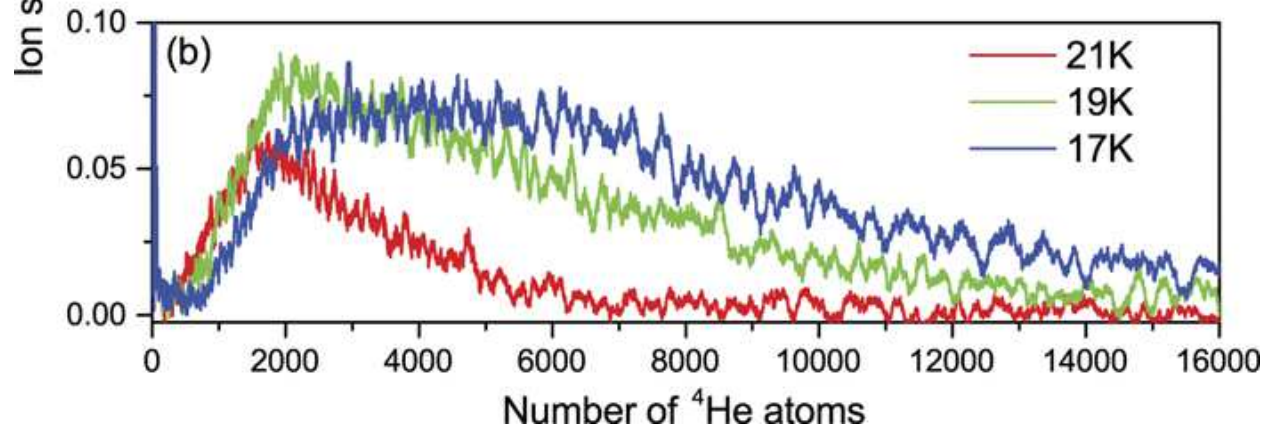
# Experimental mass spectra

415 nm ( $24100\text{ cm}^{-1}$ ) excitation (6pII)

$\Delta t = 5\text{ ps}$



single pulse  
( $\Delta t = 0$ )



*J. von Vangerow, O. John, F. Stienkemeier, and M. Mudrich, JCP* **143**, 034302 (2015)

# Theoretical simulation: $^4\text{He}$ -DFT for $A@He_{1000}$

Total energy of the system as a function of the  $^4\text{He}$  density  $\rho(\mathbf{r})$ :

$$E[\rho] = \int d\mathbf{r} \left\{ \frac{\hbar^2}{2m_{\text{He}}} |\nabla \sqrt{\rho(\mathbf{r})}|^2 + \mathcal{E}[\rho(\mathbf{r})] \right\} + \int d\mathbf{r} \rho(\mathbf{r}) V_X(|\mathbf{r}_{\text{Rb}} - \mathbf{r}|)$$

\* Orsay-Trento density functional [*F. Dalfovo, A. Lastrì, L. Pricauenko, S. Stringari & J. Treiner, PRB 1995; F. Ancilotto, M. Barranco, F. Caupin, R. Mayol & M. Pi, PRB 2005.*]

$$\begin{aligned} \mathcal{E}_c[\rho] = & \frac{1}{2} \int d\mathbf{r}' \rho(\mathbf{r}) V_{LJ}(|\mathbf{r} - \mathbf{r}'|) \rho(\mathbf{r}') + \frac{1}{2} c_2 \rho(\mathbf{r}) [\bar{\rho}(\mathbf{r})]^2 + \frac{1}{3} c_3 \rho(\mathbf{r}) [\bar{\rho}(\mathbf{r})]^3 \\ & - \frac{\hbar^2}{4m} \alpha_s \int d\mathbf{r}' F(|\mathbf{r} - \mathbf{r}'|) [1 - \tilde{\rho}(\mathbf{r})/\rho_{0s}] \nabla \rho(\mathbf{r}) \cdot \nabla' \rho(\mathbf{r}') [1 - \tilde{\rho}(\mathbf{r}')/\rho_{0s}] \\ & - \frac{m}{4} \int d\mathbf{r}' V_J(|\mathbf{r} - \mathbf{r}'|) \rho(\mathbf{r}) \rho(\mathbf{r}') [\mathbf{v}(\mathbf{r}) - \mathbf{v}(\mathbf{r}')]^2 + C \rho(\mathbf{r}) \{1 + \tanh[\beta (\rho(\mathbf{r}) - \rho_m)]\} \end{aligned}$$

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**Statics:** Minimize  $E[\rho]$  with respect to  $\Psi(\mathbf{r}) = \sqrt{\rho(\mathbf{r})}$ :  $\rightarrow$  Euler-Lagrange equation:

$$\left\{ -\frac{\hbar^2}{2m_{\text{He}}} \nabla^2 + \frac{\delta \mathcal{E}}{\delta \rho} + V_X(|\mathbf{r}_{\text{Rb}} - \mathbf{r}|) \right\} \Psi(\mathbf{r}) = \mu \Psi(\mathbf{r}), \quad \Psi(\mathbf{r}) = \sqrt{\rho(\mathbf{r})}$$

Sum of pairwise interactions with  $V_X$  from the literature (Patil 1991)

Imaginary time propagation on a 3-D grid  $\rightarrow \rho_0(\mathbf{r}) = \Psi_0^2(\mathbf{r})$

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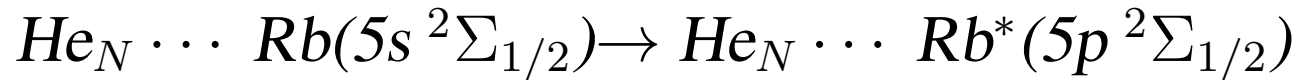
**Dynamics:** Time-dependent DFT Minimize the action /  $\Psi_{\text{He}}(\mathbf{r}, t)$

$\Psi_{\text{He}}(\mathbf{r}, t)$  coupled to **electronic wave packet**  $|\lambda\rangle$  and to **classical**  $\mathbf{r}_{\text{Rb}^*}$  (mean field)

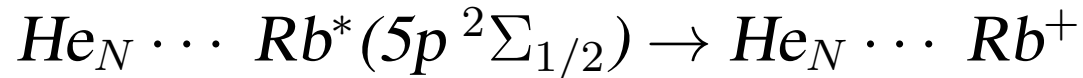
$$\left\{ \begin{array}{l} i\hbar \frac{\partial}{\partial t} \Psi_{\text{He}} = \left[ -\frac{\hbar^2}{2m_{\text{He}}} \nabla^2 + \frac{\delta \mathcal{E}}{\delta \rho} + V_\lambda(\mathbf{r} - \mathbf{r}_{\text{Rb}^*}) \right] \Psi_{\text{He}} \rightarrow \rho_{\text{He}}(\mathbf{r}, t) = |\Psi_{\text{He}}(\mathbf{r}, t)|^2 \\ i\hbar \frac{\partial}{\partial t} |\lambda\rangle = \mathcal{H}_{\text{el}} |\lambda\rangle = (\mathcal{H}_{\text{DIM}} + \mathcal{H}_{\text{SO}}) |\lambda\rangle \\ \text{where } |\lambda\rangle = \sum_{i_s} \lambda_{i_s}(t) |p_i, m_s\rangle \rightarrow V_\lambda = \langle \lambda | \mathcal{H}_{\text{el}} | \lambda \rangle \\ m_{\text{Rb}} \ddot{\mathbf{r}}_{\text{Rb}^*} = -\nabla_{\mathbf{r}_{\text{Rb}^*}} \left[ \int d\mathbf{r} \rho(\mathbf{r}) V_\lambda(\mathbf{r} - \mathbf{r}_{\text{Rb}^*}) \right] = -\int d\mathbf{r} \nabla \rho(\mathbf{r}) V_\lambda(\mathbf{r} - \mathbf{r}_{\text{Rb}^*}) \end{array} \right.$$

## He-TDDFT simulation

$t_1$ : 1<sup>st</sup> laser excitation



$t_2$ : 2<sup>nd</sup> laser ionization



$$t_2 = 55 \text{ ps}$$

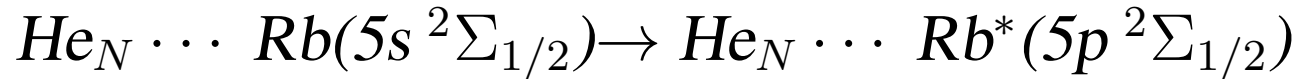
$$t_2 = 20 \text{ ps}$$

*Von Vangerow, Coppens, Leal, Pi, Barranco, Halberstadt, Stienkemeier and Mudrich, JPCL 8, 307 (2017)*

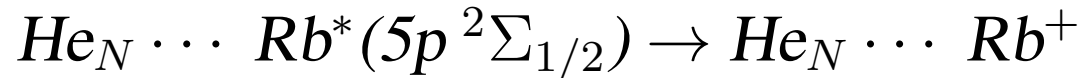


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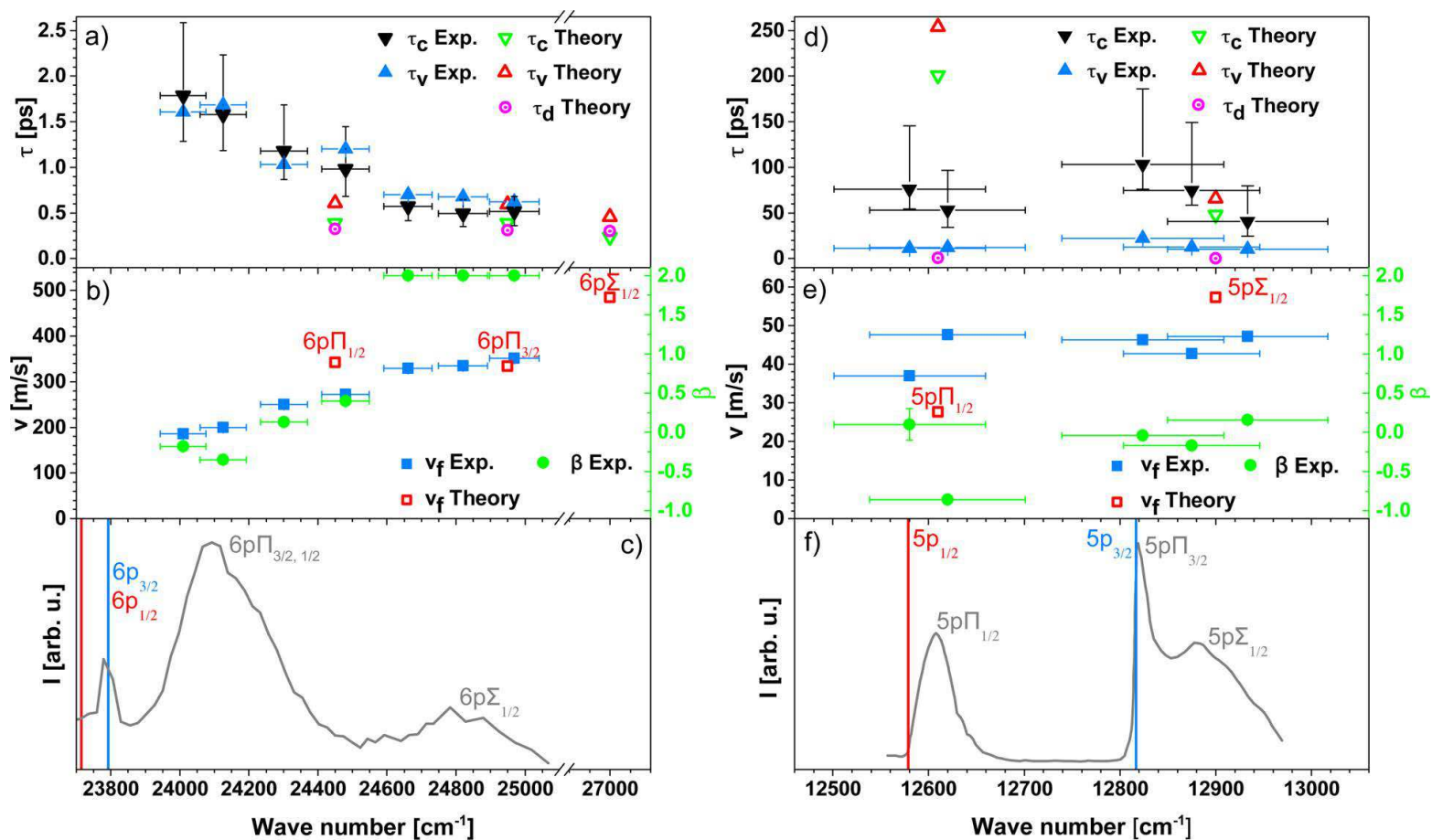


Critical time delay  $\tau_c$ :

- if  $t_2 - t_1 > \tau_c$ ,  $\text{Rb}^+$  escapes;
- if  $t_2 - t_1 < \tau_c$ ,  $\text{Rb}^+$  turns around and gets solvated.

*Von Vangerow, Coppens, Leal, Pi, Barranco, Halberstadt, Stienkemeier and Mudrich,  
JPCL 8, 307 (2017)*

# Comparison with experiments



Von Vangerow et al., *J. Phys. Chem. Lett.* **8**, 307 (2017)

## Everything works...

- He-TDDFT gives a good qualitative description of  $\text{Rb}^*(6p, {}^2\Sigma_{1/2}, {}^2\Pi_{3/2}, {}^2\Pi_{1/2})$  and  $\text{Rb}^*(5p, {}^2\Sigma_{1/2}, {}^2\Pi_{1/2})$  photodissociation from  $\text{He}_N$
- 2 orders of magnitude difference for critical fall-back times between 5p and 6p excitation well reproduced

### except for $\text{Rb}^*(5 {}^2\Pi_{3/2})$ :

- ejected exciplex in experiment
- bound exciplex in our simulations

- initial Rb position? (zero-point delocalization or  $T = 0.4$  K):
- initial electronic excitation could be mixed?
- non-adiabatic transition to  ${}^2\Pi_{1/2}$  ? (*Moroshkin et al. 2006 solid He; several Ak\* experiments*)

→ existence of **He-induced electronic relaxation** (non-radiative transitions)?

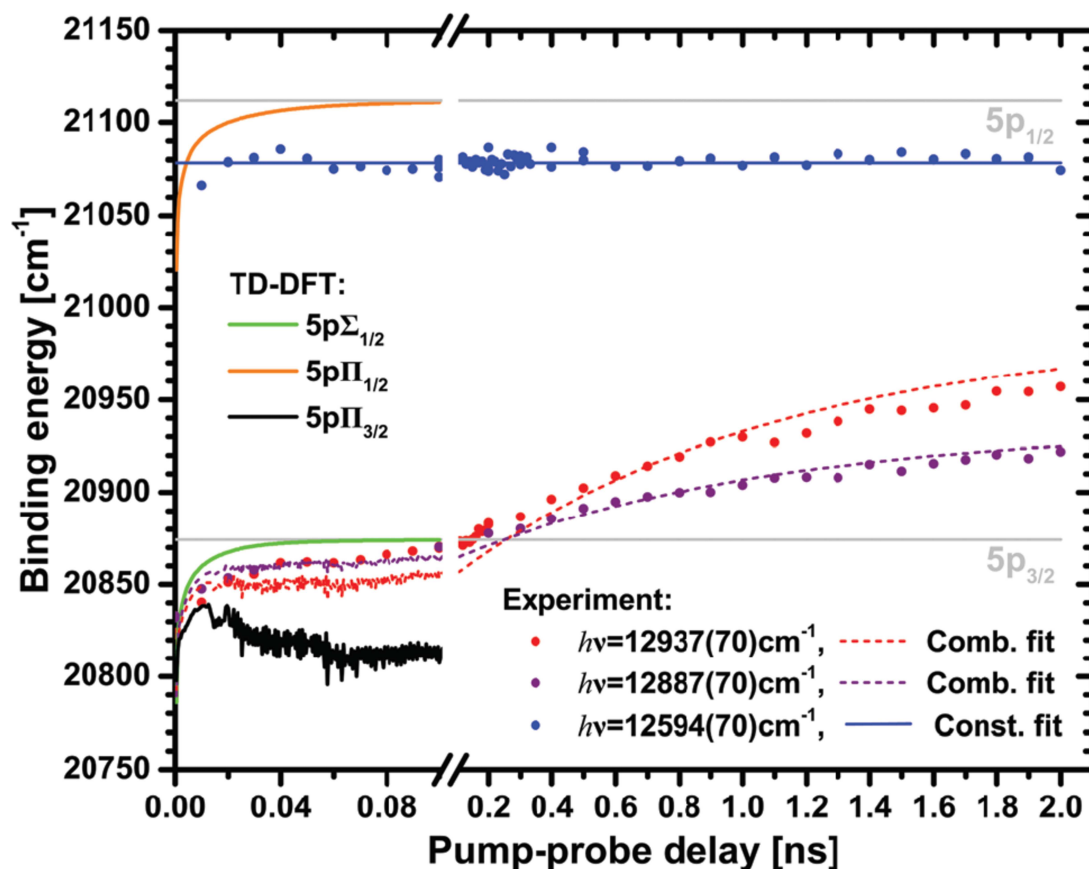
# *Including spin-orbit relaxation*

## **Model:**

- He-TDDFT for Rb( $5p^2\Pi_{3/2}$ )  
→ bound exciplex
- after 60 ps, jump down to  $^2\Pi_{1/2}$ ;  
transfer  $x\%$  of  $\Delta E$  to Rb  
→ He-Rb ejection ?

He-TDDFT for  $x=12.5$

## Including spin-orbit relaxation



12.5% of  $\Delta E=165\text{ cm}^{-1}$  is enough to eject the exciplex

Fit to the slow electrons binding energies as a function of time:

${}^2\Pi_{3/2} \rightarrow {}^2\Pi_{1/2}$  relaxation time = 700 ps

Coppens et al., *Phys. Chem. Chem. Phys.* **20**, 9309 (2018)

*He-induced electronic transitions:*

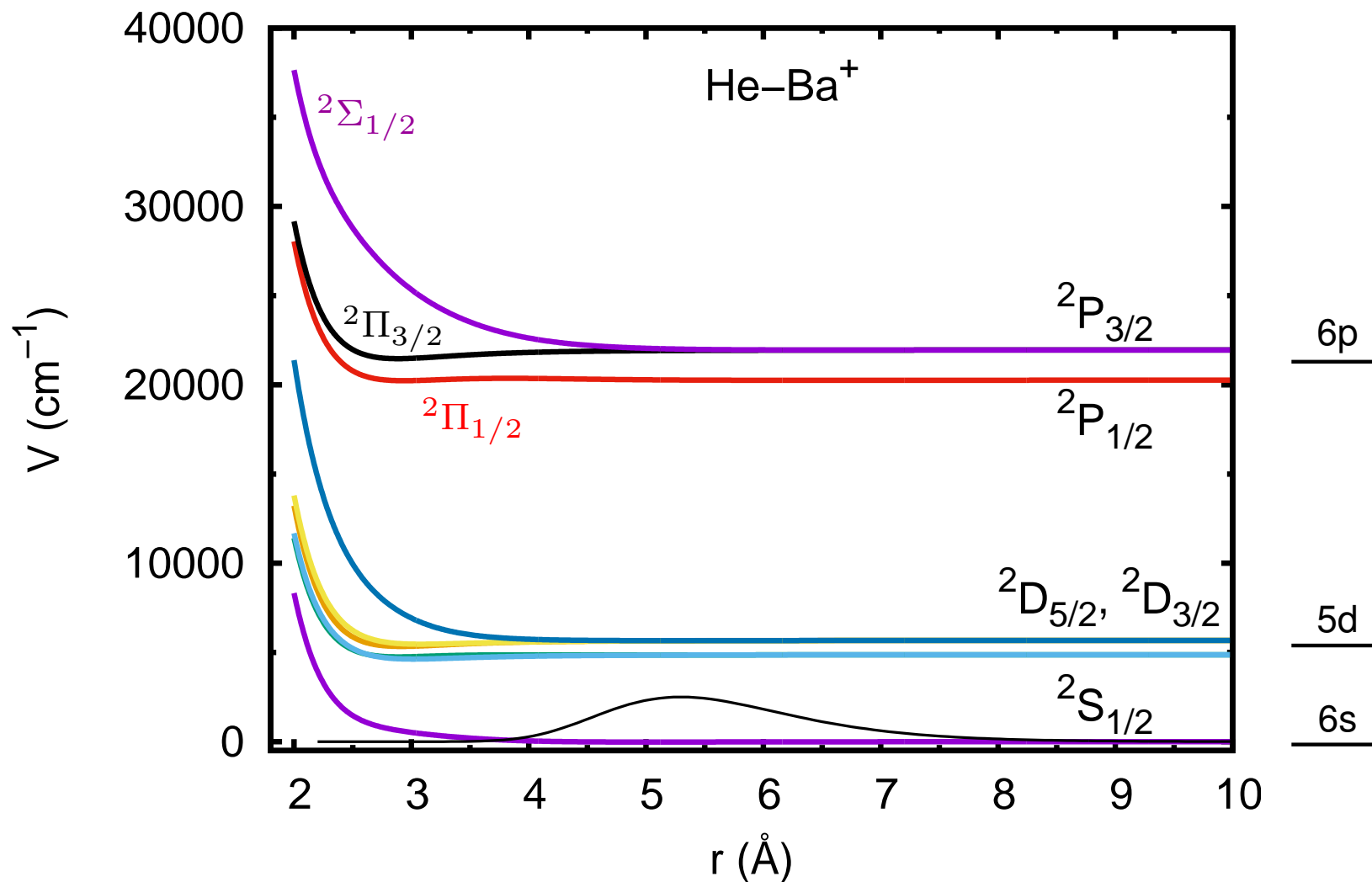
*How to predict them?*

## *He-induced electronic transitions:*

### *How to predict them? The case of $Ba^+$*

- He-TDDFT gives very good results for  $Ba^+$  solvation upon ionization,
- also for its absorption and emission spectrum once it is solvated
- but Zhang and Drabbels observed  $Ba^+$  ejection upon  $6p \leftarrow 6s$  excitation which could not be reproduced

# He-Ba<sup>+</sup> potential energy curves

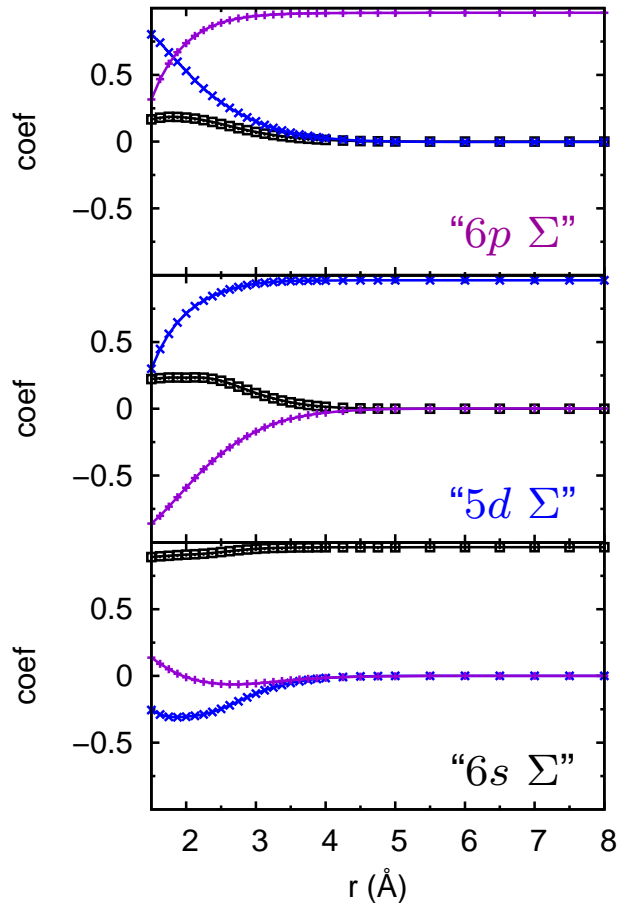


PEC from M. Mella and MF. Cargnoni, *JPCA* **118**, 6473 (2014)



# Model for potential energies and couplings

$\Lambda = 0$  mixing coefficients



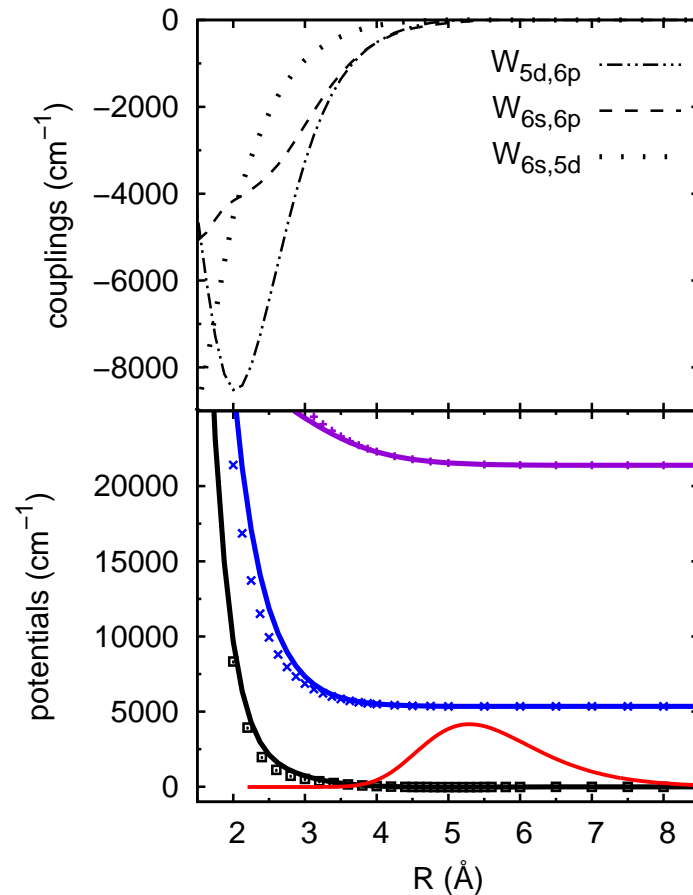
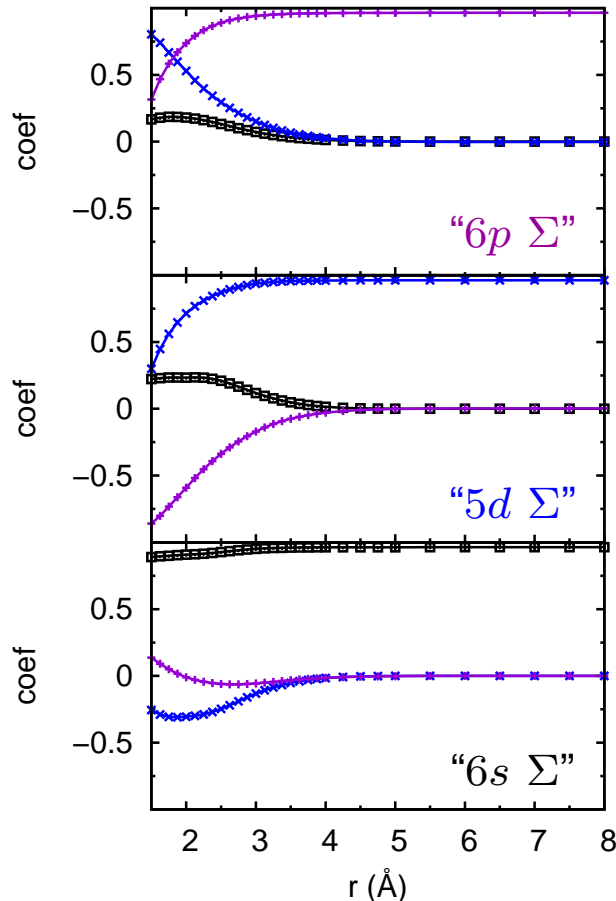
MRCI electronic structure calculation; Eigenvectors projected onto asymptotic ones:  
 $Ba^+$  6s, 5d, 6p

# Model for potential energies and couplings

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$\Lambda = 0$  PEC and couplings



MRCI electronic structure calculation; Eigenvectors projected onto asymptotic ones:

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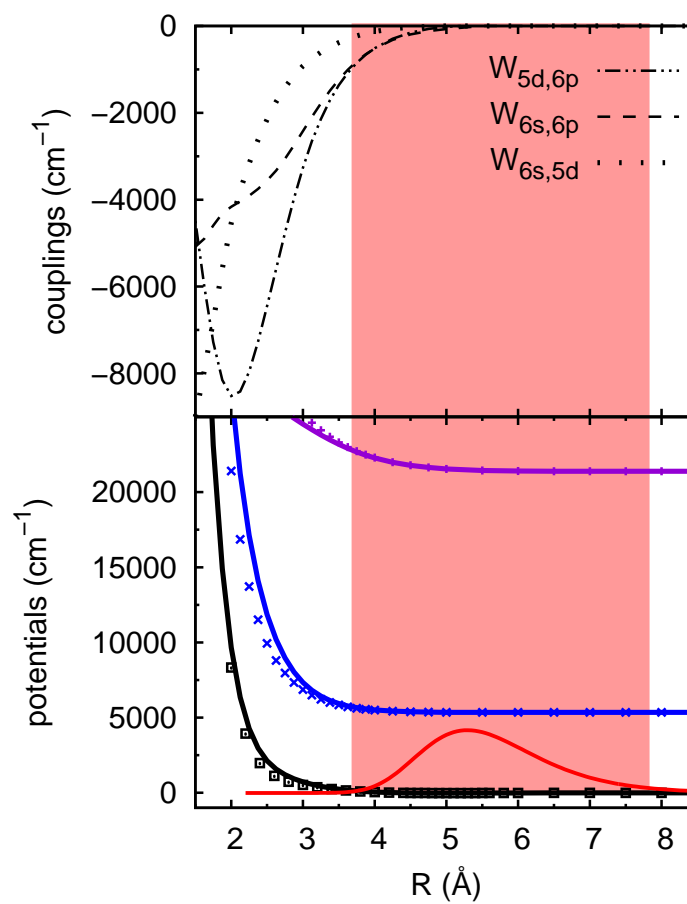
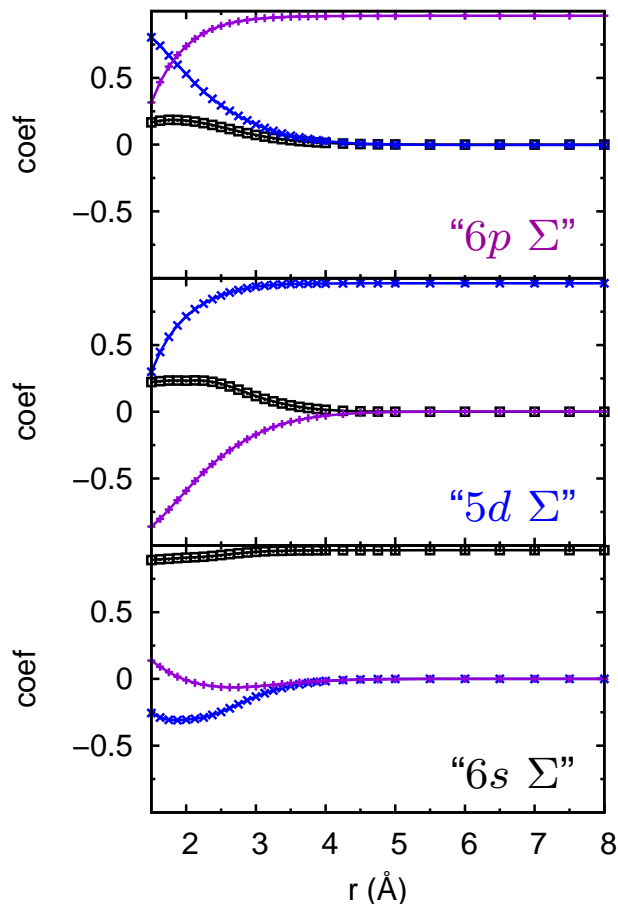
Eigenvector coefficients = Transformation matrix to  $Ba^+$  states → diabatic curves and couplings

# Model for potential energies and couplings

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MRCI electronic structure calculation; Eigenvectors projected onto asymptotic ones:

$Ba^+$  6s, 5d, 6p

Eigenvector coefficients = Transformation matrix to  $Ba^+$  states → diabatic curves and couplings

## *So from this potential model:*

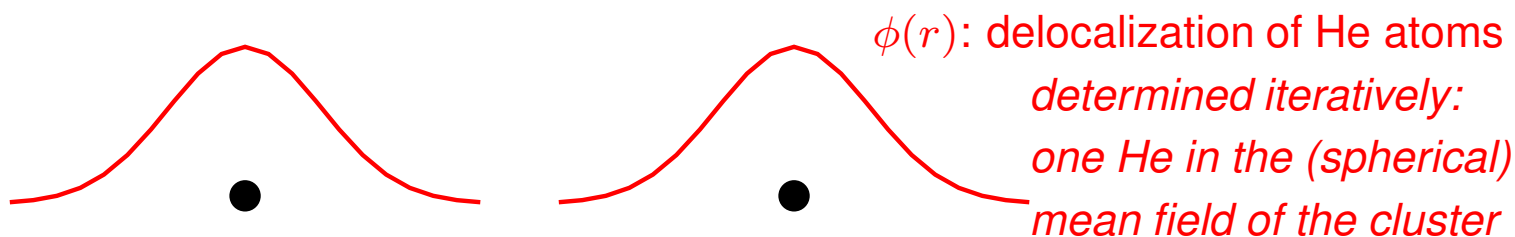
- Model for He-Ba<sup>+</sup> electronic Hamiltonian → diabatic representation (potentials, couplings) for 6s, 5d, 6p mixing
- Couplings are too weak to induce electronic relaxation in HeBa<sup>+</sup>
- But they can induce electronic relaxation in configurations with more He atoms
- More relaxation when getting out of cylindrical symmetry (vibrations)

*Vindel Zandbergen et al., J. Phys. Chem. **148**, 144302 (2018);*

→ Dynamics study !

# Influence of the Helium nanodroplet environment: Explicit Zero-Point Averaged He dynamics (ZPAD)

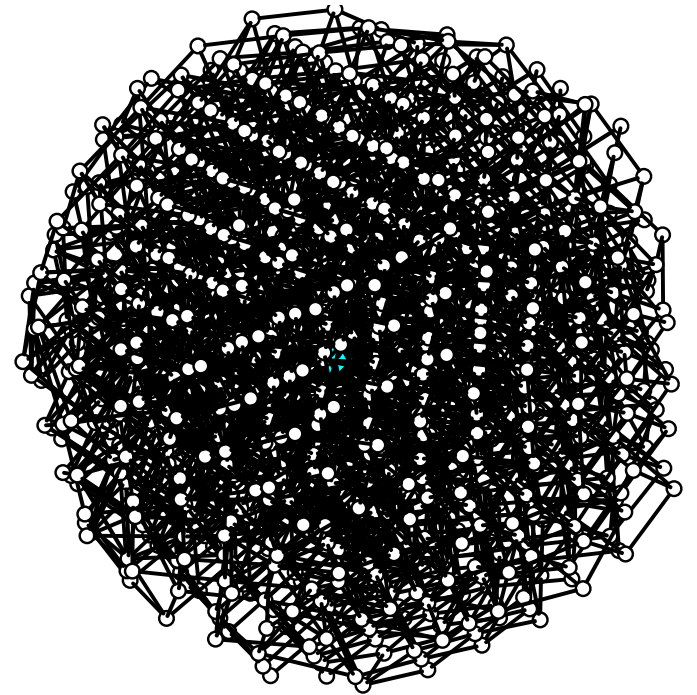
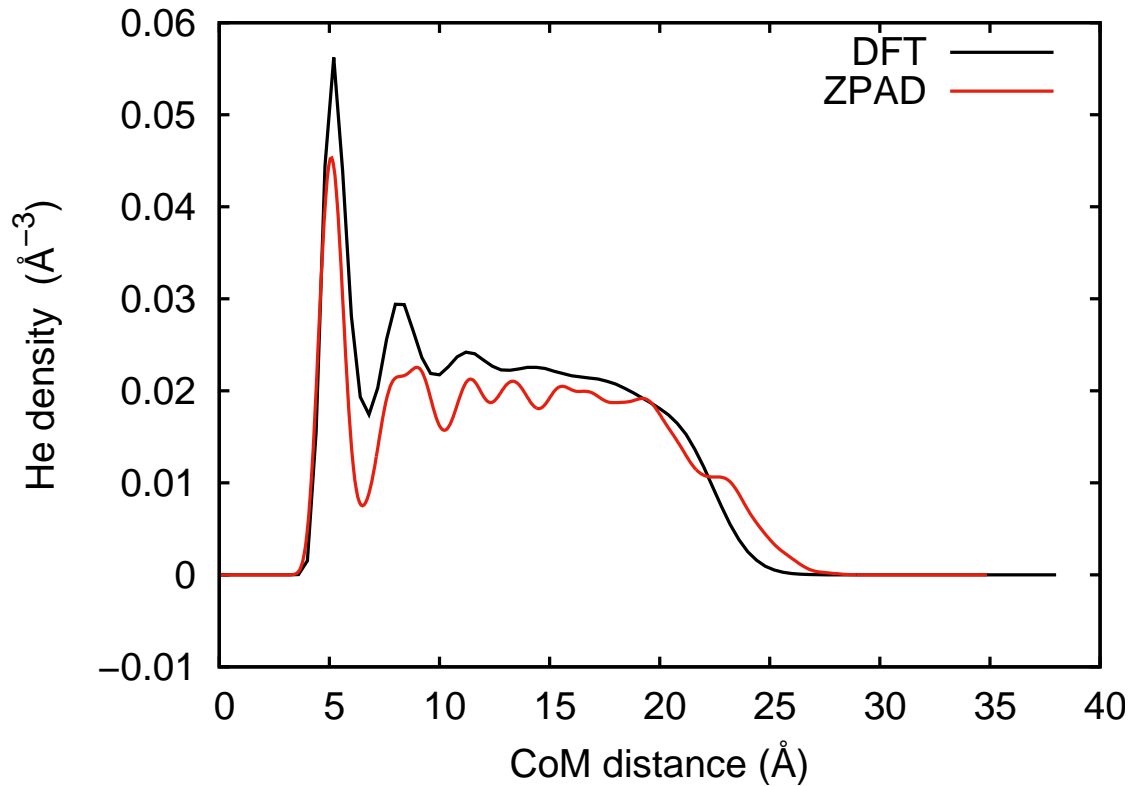
- ZPAD Dynamics *cf.* David Bonhommeau's thesis on  $Rg_n^+ He_N$ 
  - $(Ba^+)^* - He_N$  DIM + SO; Molecular Dynamics with Quantum Transitions (Tully)
  - $He_N$ : classical (adiabatic) taking the ZPE effects into account on average



- Potentials
  - He ... He: ZPAD equivalent to **classical dynamics** with **effective potential**  
(He-He potential averaged over the He wave function)  
(Portwich 1995, Slavíček et al., JPCA 2003; Sterling et al., JCP 1995)
  - He ...  $Ba^+$ , He ...  $(Ba^+)^*$   $\Sigma$ ,  $\Pi$ : analytic forms fitted to F.Cargnoni's curves  
averaged over the He wave function

(TDDFT; semiclassical path integral centroid MD by Takayanagi and Shiga, multiple hard sphere collisions by Drabbels, CBF by Zillich; but most are mean field)

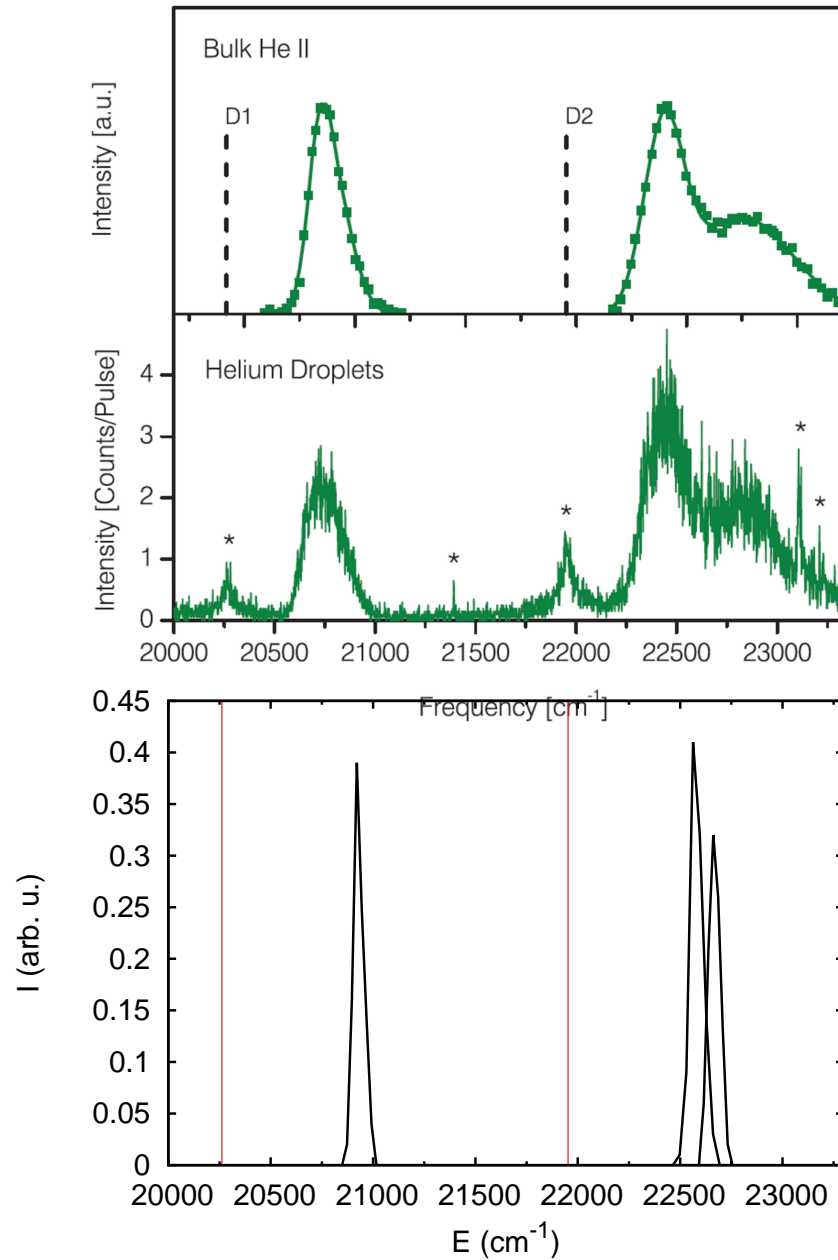
## Reasonable initial conditions



*He-DFT from Leal, Mateo, Hernando, Pi, Barranco, Ponti, Cargnoni and Drabbels, PRB*  
**90**, 224518 (2014).

*He<sub>1000</sub>Ba<sup>+</sup> ZPAD structure calculated by David Bonhommeau*

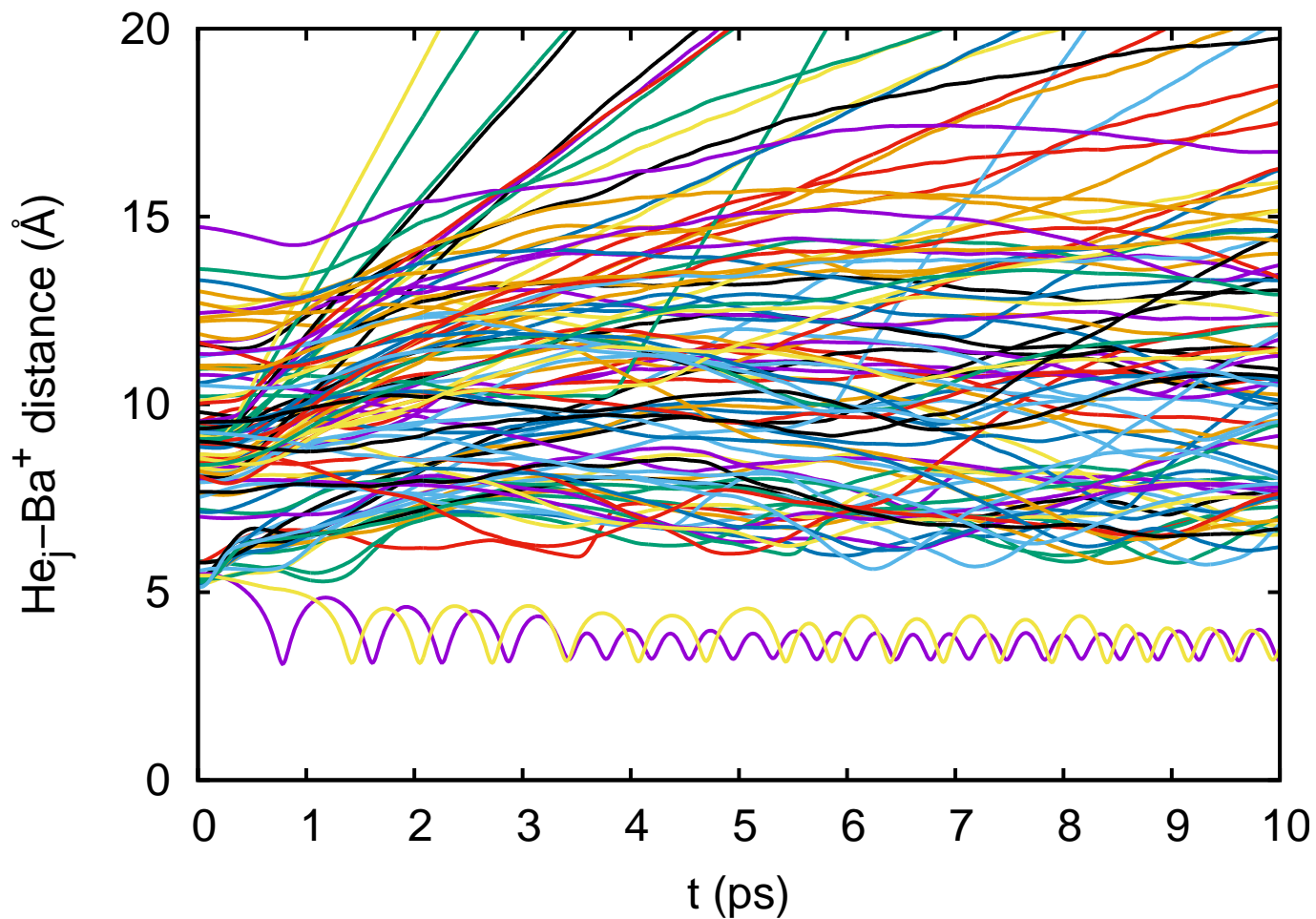
# Absorption spectrum



Zhang and Drabbels,  
*JCP* **137**, 051102 (2012)

shift  $\sim$  correct,  
width too small  
missing quantum fluctuations

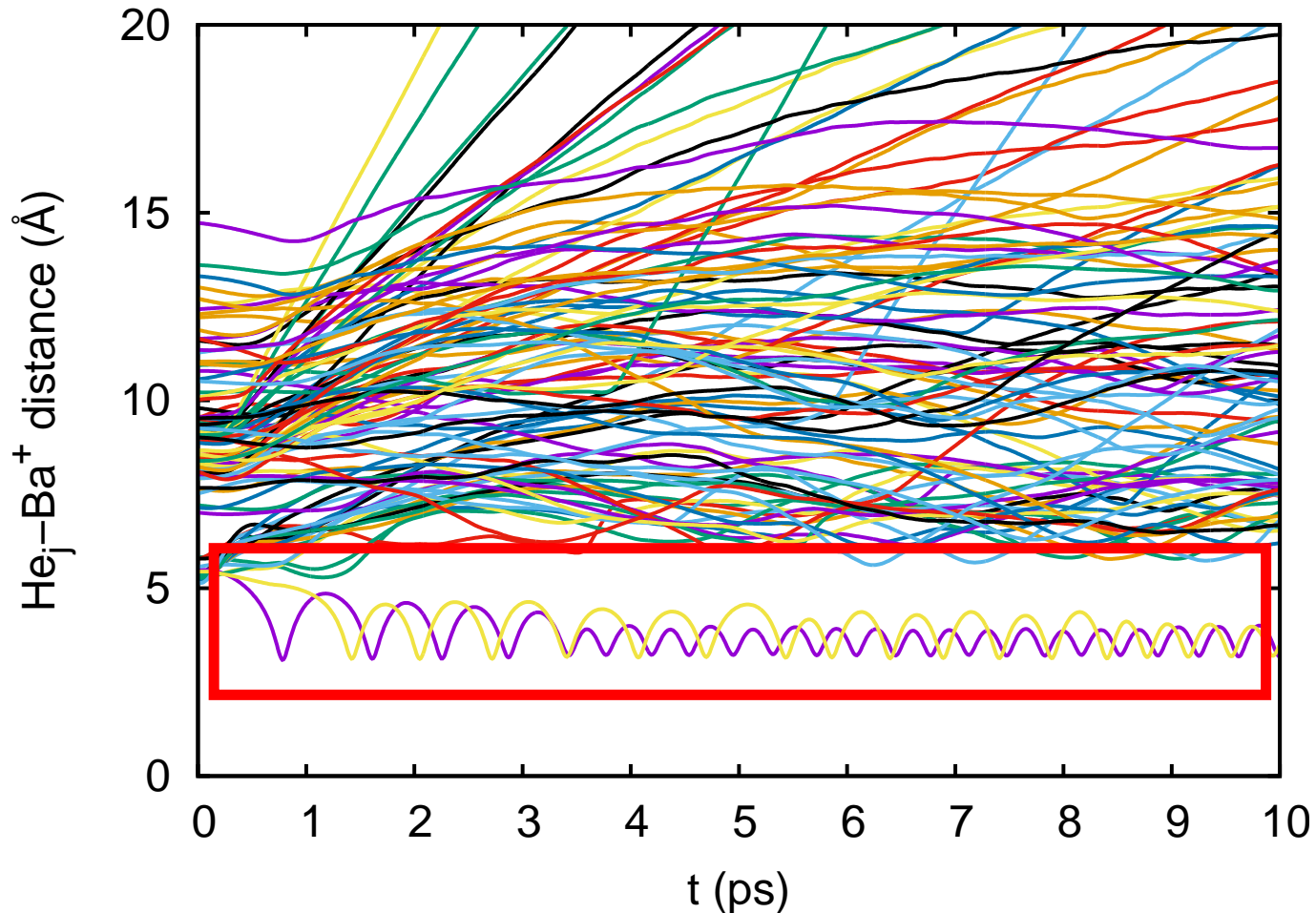
## *Dynamics for $Ba^+ \ ^2\Sigma_{1/2}$ excitation: typical trajectory*



As expected, many He dissociate, some of them fast and early.



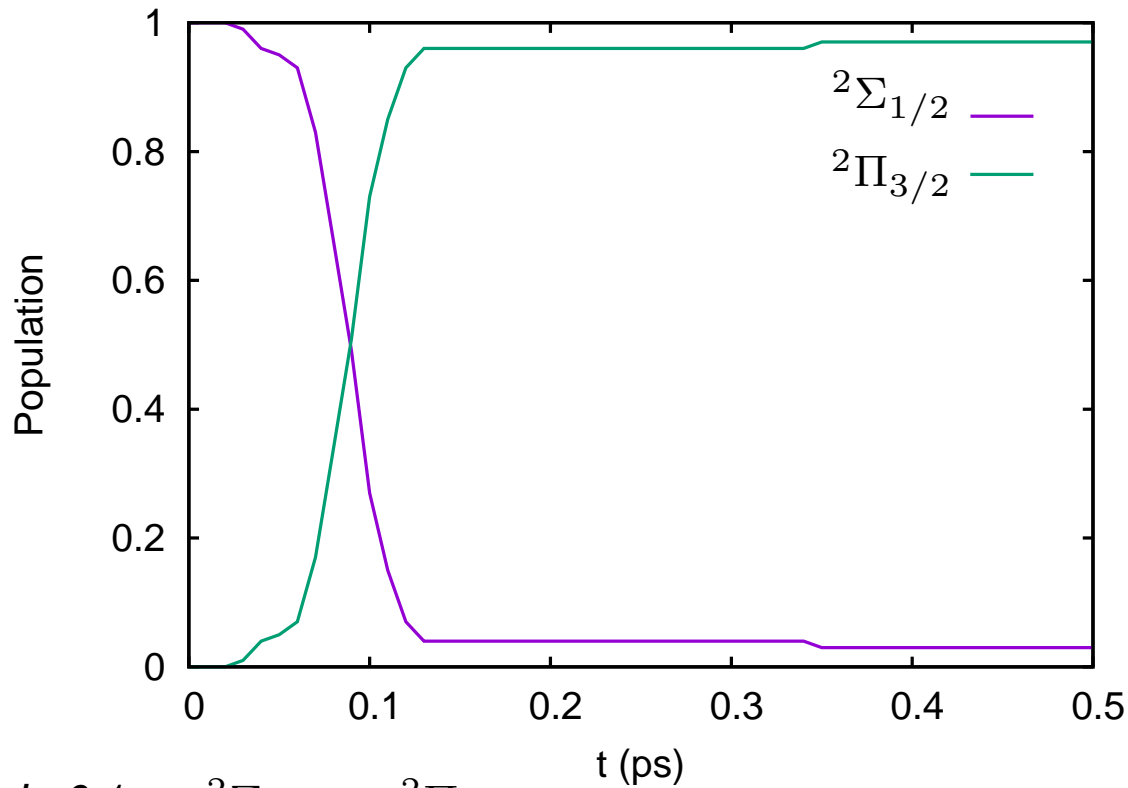
## Dynamics for $Ba^+ \ ^2\Sigma_{1/2}$ excitation: typical trajectory



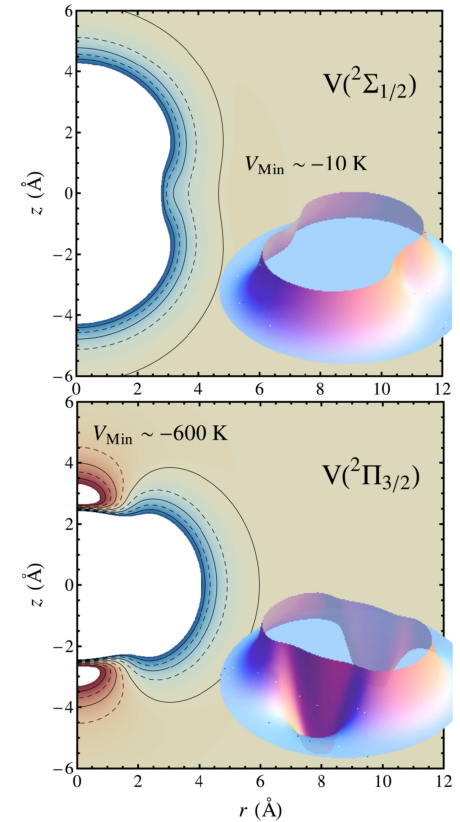
As expected, many He dissociate, some of them fast and early.

**BUT:** a double excimer is formed??? This normally only happens for  $^2\Pi_{3/2}$

# Electronic transition at work



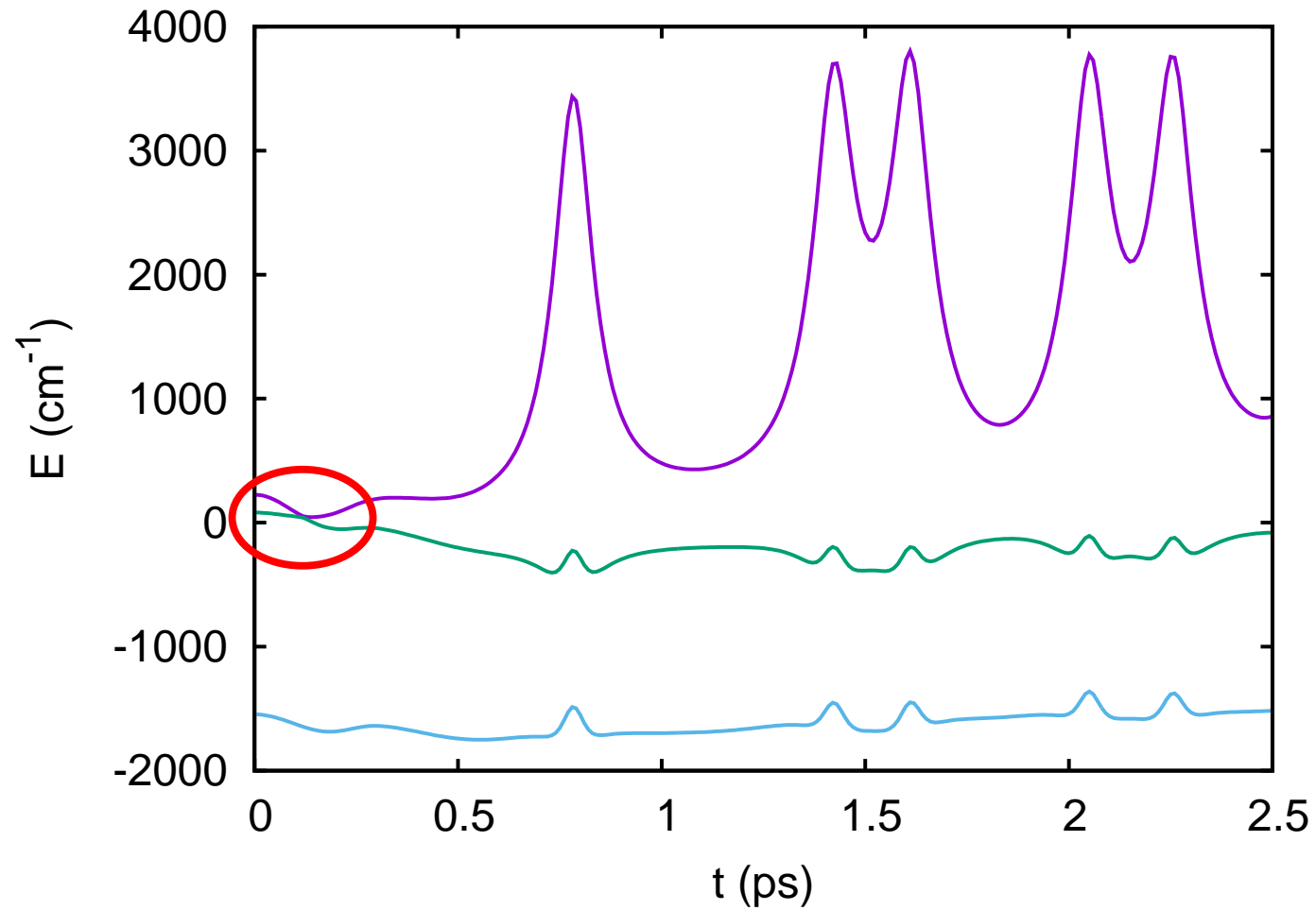
In 0.1 ps,  $2\Sigma_{1/2} \rightarrow 2\Pi_{3/2}$   
(here: averaged over 100 trajectories)



3D PES from Leal et al.,  
*JCP* 144, 094302 (2016)

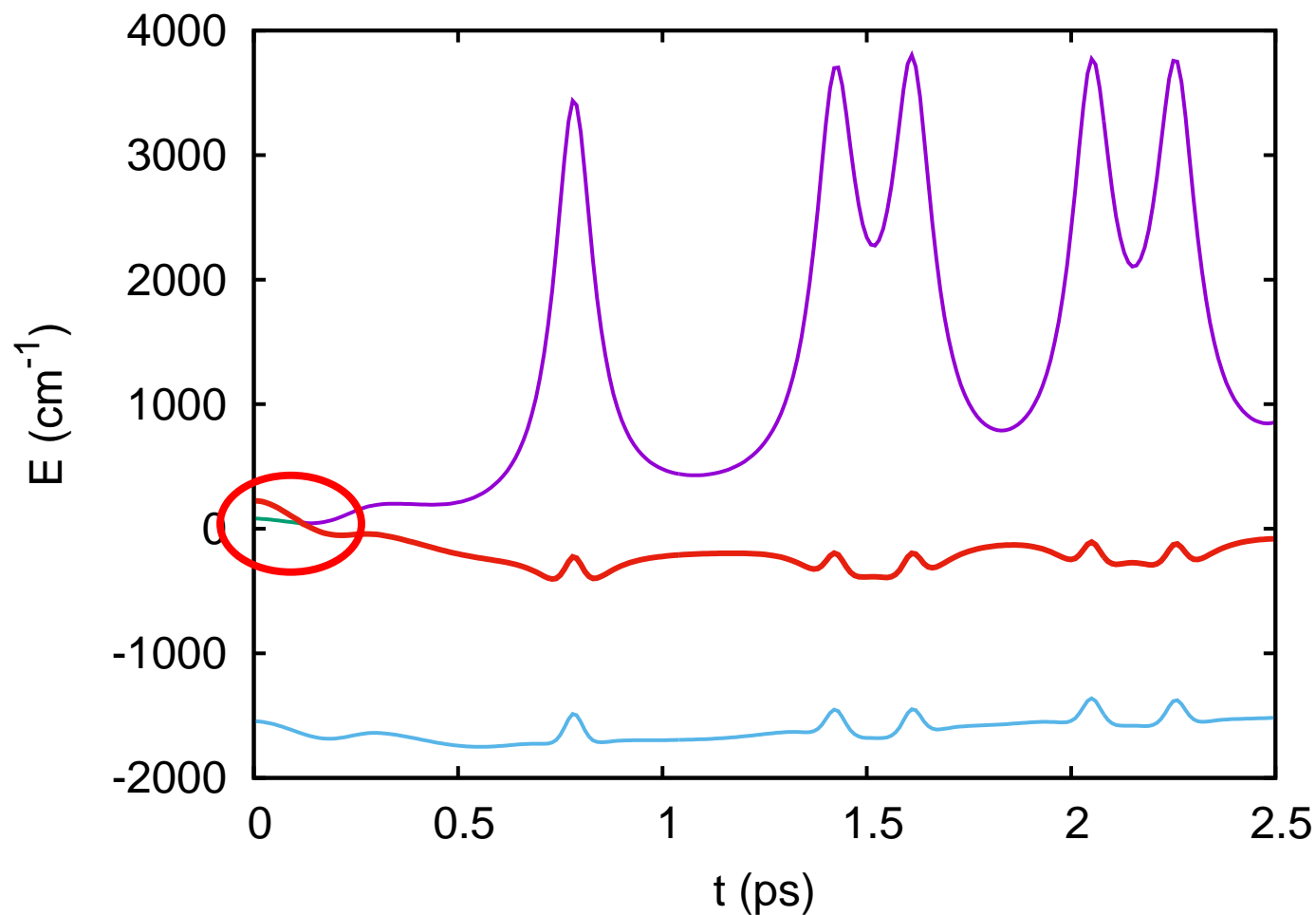
# How did that happen?

Potential energy curves during the “typical” trajectory



# How did that happen?

Potential energy curves during the “typical” trajectory



At  $\sim 0.1$  ps,  $^2\Sigma_{1/2}$  and  $^2\Pi_{3/2}$  come very close:  $\Delta \simeq 20 \text{ cm}^{-1}$  (quasi-spherical symmetry)  $\rightarrow$  strong coupling  $\rightarrow$  “hop” (electronic transition)

# Conclusion/Perspectives

- Yes, there are electronic transitions induced by He at 0.4 K!
  - He-(TD)DFT is a very accurate and realistic tool to describe doped He droplet structures and dynamics
  - It can describe the dynamics following an electronic transition
  - Molecular model is complementary since it can describe potentials and couplings, and predict transitions
  - Approximate (ZPAD) quantum molecular dynamics for  $\text{Ba}^+$  at  $\text{He}_N$  (not finished): include quantum fluctuations in the initial conditions, ...
- couple TDDFT to other dynamics

*Many thanks for your attention!*



UNIVERSITÉ  
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PAUL SABATIER

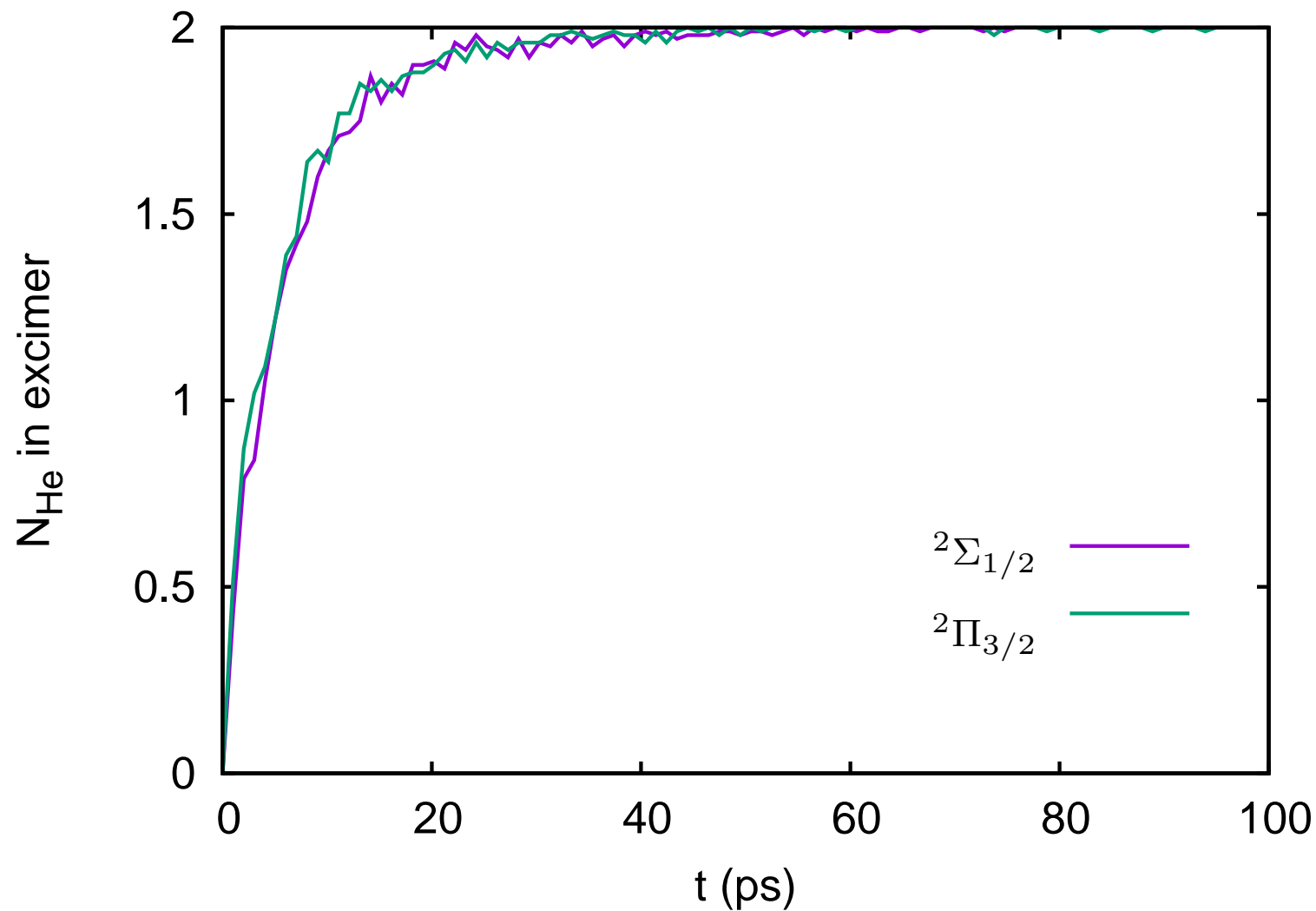


Université  
de Toulouse



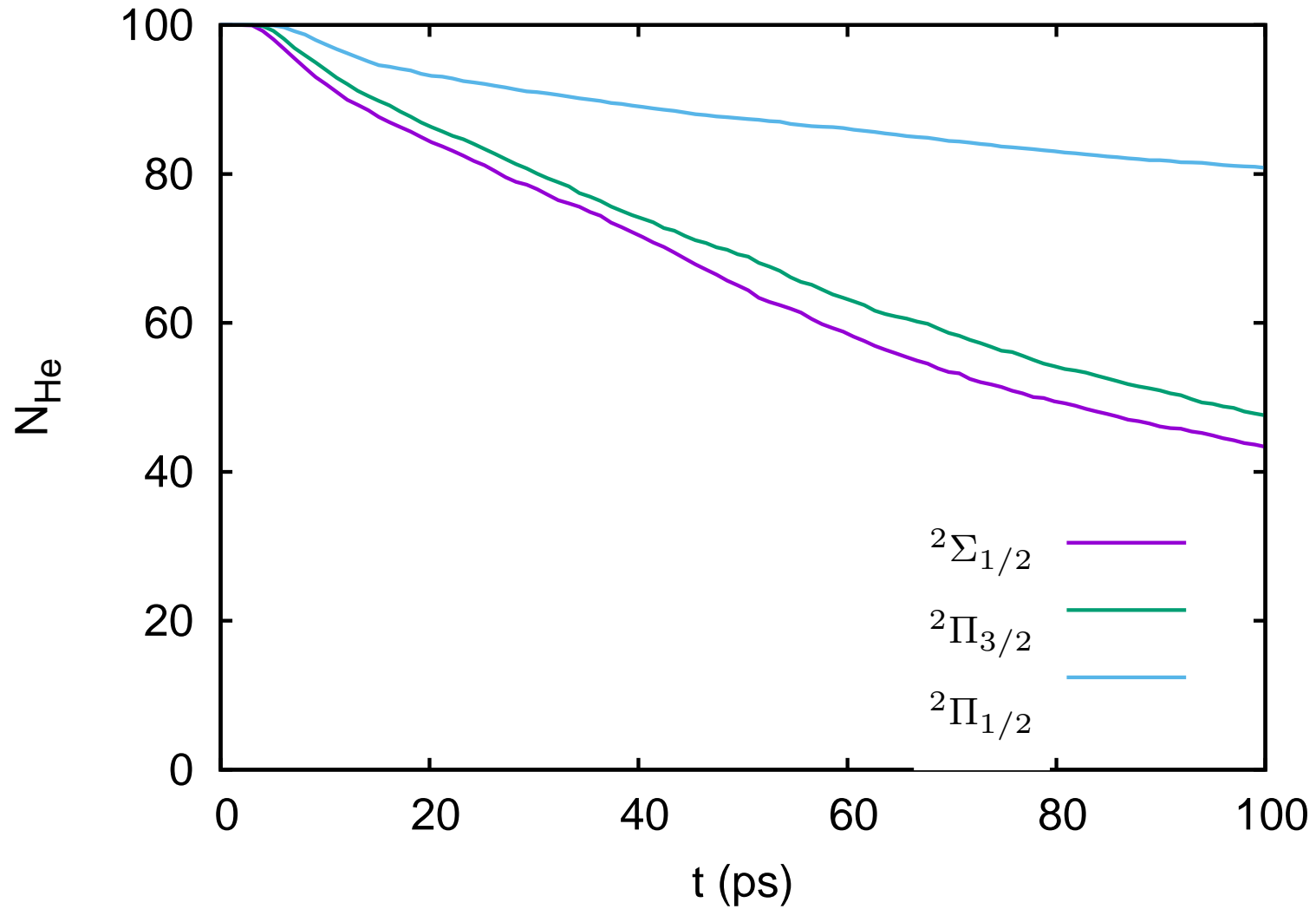
Programme Chaires d'Attractivité  
Manuel Barranco

## Excimer formation dynamics



*excimer formation: same kinetics for  $^2\Pi_{3/2}$  and  $^2\Sigma_{1/2}$*

# He evaporation dynamics



Slower dynamics for  $2\Pi_{1/2}$