

Quantum fluids of light

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Quantum fluids of light?

General motivations

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Why cavity polaritons?

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Band structure?

How do we perform experiments?

Why non-linearities?

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Where do we go??

Many body physics: fascinating macroscopic physics

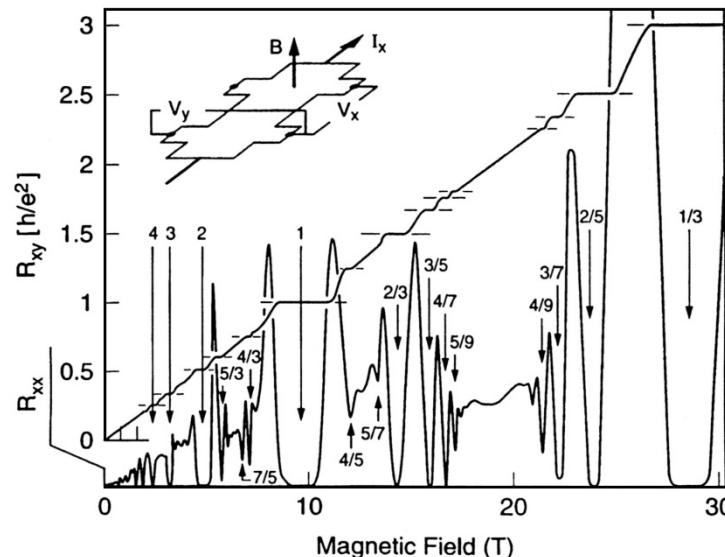
Superfluidity



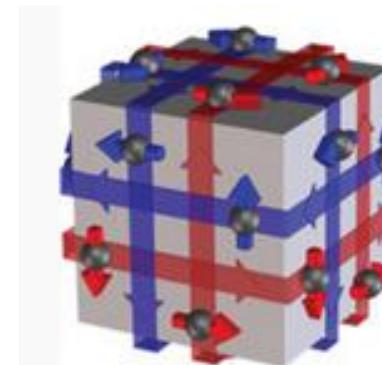
Graphene



Fractional Quantum Hall effect



Topological insulators



Quantum fluids of light as an analogous system

$$i\hbar \frac{\partial \psi}{\partial t} = H\psi$$

Same Hamiltonian => Same physical properties

Why use synthetic quantum material?

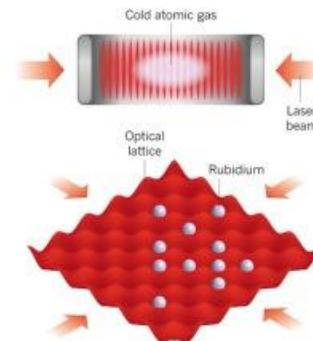
- System « easier » to probe, manipulate
- Create new geometries, new properties => Applications
- Realize experiments impossible to predict because of complexity of numerical simulations (many body physics): **quantum simulations**

QUANTUM BOARD GAMES

The set-ups of quantum simulators are different, but the concept is the same: first take atoms, ions or electrons, cool them to cryogenic temperatures and arrange them in an orderly grid. Then tune the interactions on the grid to mimic a more complex material.

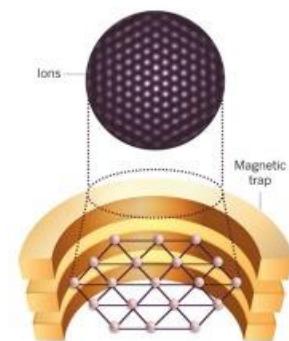
COLD ATOMS

Rubidium atoms are held in place by criss-crossed laser beams, which can also be used to tweak individual particles. A single pair of lasers holds the atoms in a one-dimensional column (top), whereas two pairs hold them in a grid (bottom). Some excitations in the grid system behave like the Higgs particle.



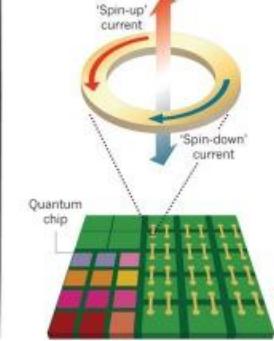
TRAPPED IONS

A combination of electric and magnetic fields trap charged, ionized atoms in an orderly grid. The ions wiggle and rotate in a way that mimics the interactions of quantum magnetism — a phenomenon that can't be simulated in classical systems.



SUPERCONDUCTING LOOPS

A quantized loop of current can flow clockwise, anticlockwise or in a superposition of both in a superconducting circuit (top). An array of such loops (bottom) can be manipulated to simulate various quantum systems — and perhaps even biological processes such as photosynthesis.



Nature 14 November 2012

Original idea: [R. Feynman., Int. J. Theor. Phys. 21, 467 (1982)]

Quantum fluids of light as an analogous system

- Photons have no mass !
- Photons have no charge !
- Photons do not interact (or so weakly) !
- Photons are bosons (not fermions like electrons !)



"Well, there's your problem - you were holding the plans wrong."



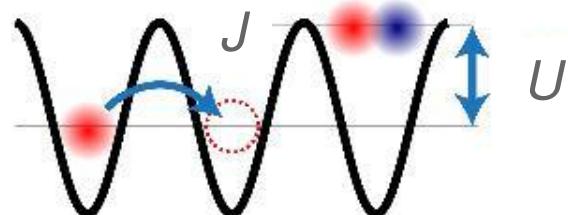
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Quantum fluids of light as an analogous system

- Photons have no mass ! **Effective mass in a cavity**
- Photons have no charge !! **Artificial gauge field**
- Photons do not interact (or so weakly) ! **Yes when coupled to electronic excitations**
- Photons are bosons (not fermions like electrons !) **Nobody is perfect!**

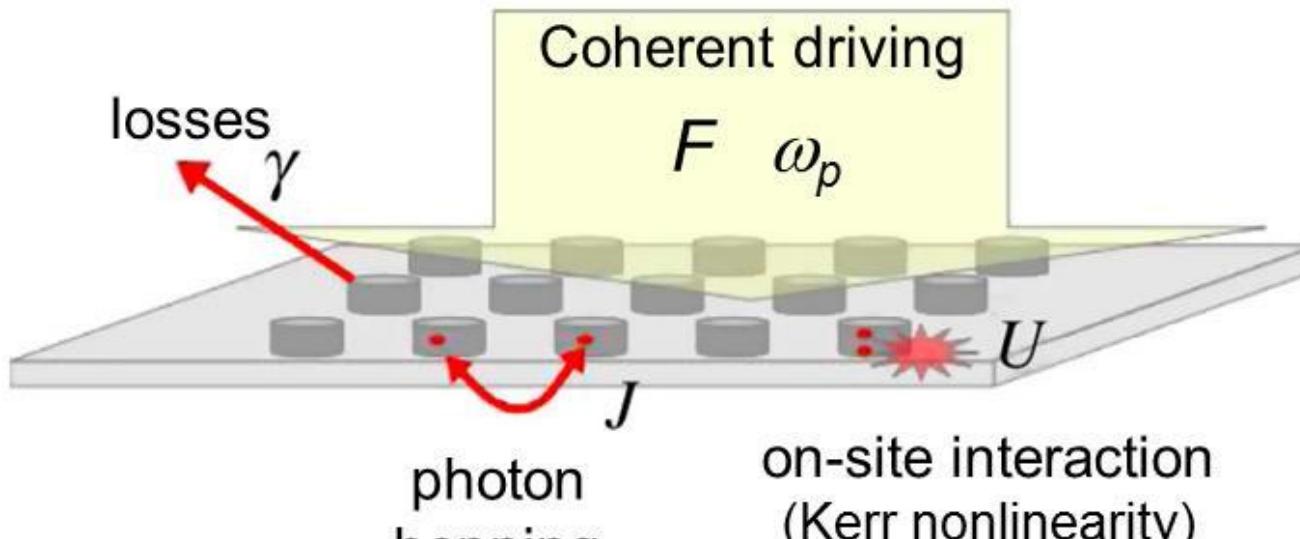


Bose-Hubbard $\hat{H} = -J \sum_{\langle m,n \rangle} \hat{a}_m^\dagger \hat{a}_n + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1)$



Driven-dissipative photonic Bose-Hubbard model

Out of equilibrium quantum physics



Ciuti & Carusotto, Rev. Mod. Phys. **85**, 299 (2013)

M.J. Hartman, Journal of Optics (2016)

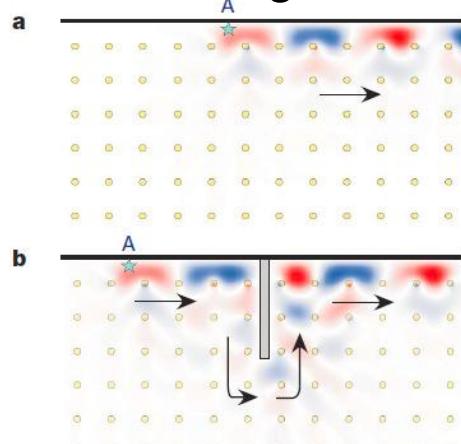
C.Noh and DG Angelakis, Report on progress in Physics (2016)

A. Biella et al., Phys. Rev. A 96, 023839 (2017)

F. Vincentini et al., Phys. Rev. A 97, 013853 (2018)

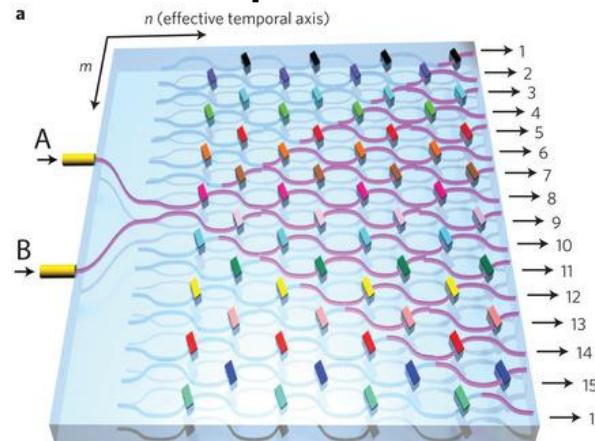
Emulation with light

Chiral edge states



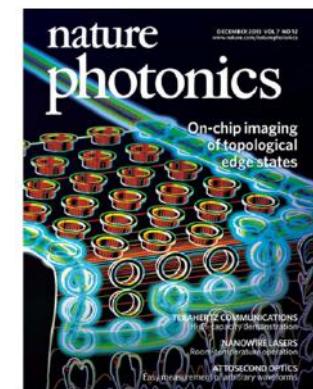
Zheng Wang et al.,
Nature **461** 772 (2009)

Random quantum walk



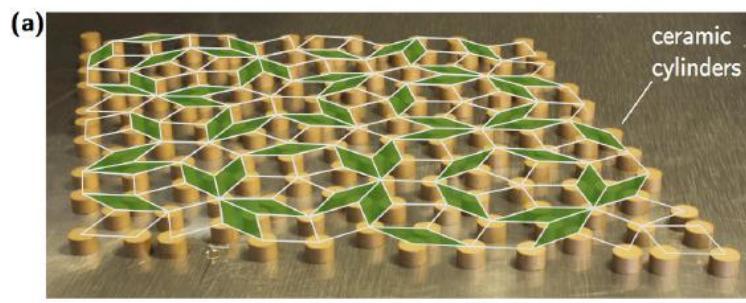
A. Crespi, Nature Photonics 7, 322 (2013)

Topological edge states Si



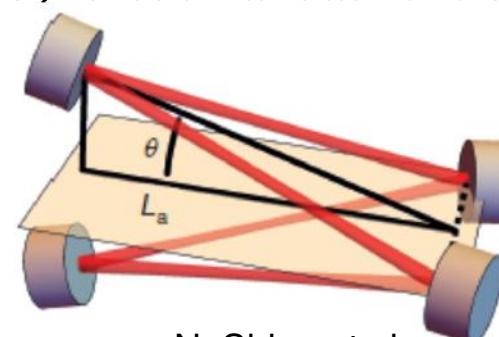
M. Hafezi,
Nat. Phot. **7** 1001 (2013)

Quasicrystal



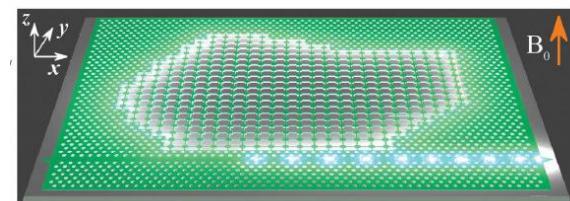
P. Vignolo et al., New J. Phys. 16 (2014) 043013 Nature 354, 671 (2016)

Synthetic Landau levels



N. Shine et al.

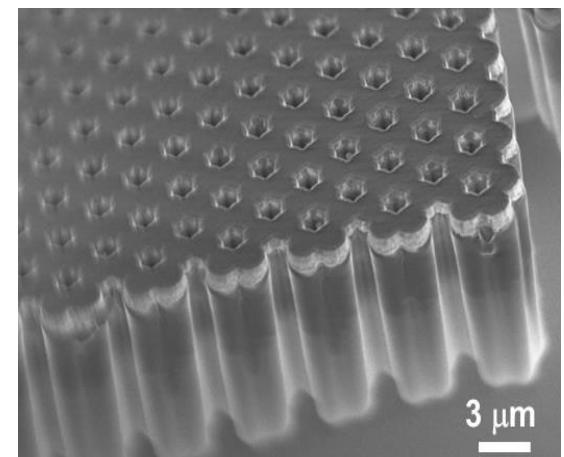
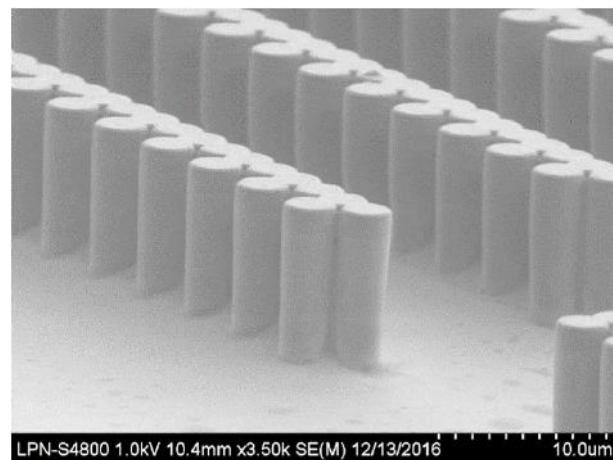
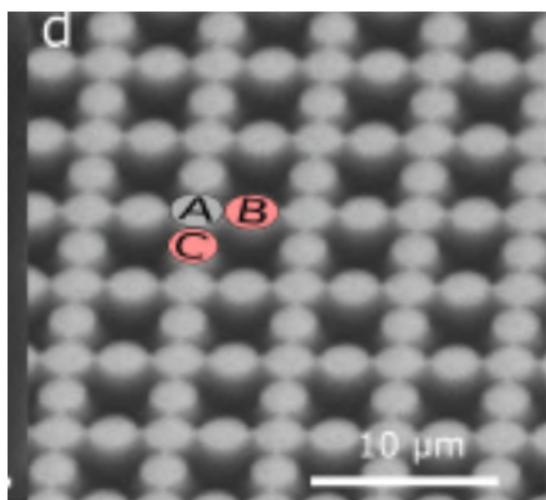
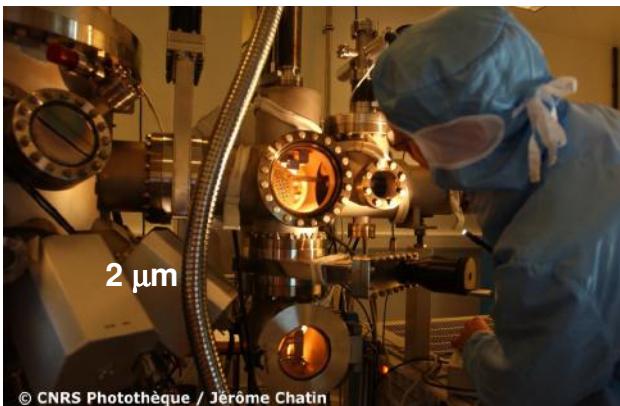
Non reciprocal lasing



B.Bahari et al., Science
10.1126/science.aao4551(2017)

M. A. Bandres et al. Science
10.1126/science.aar4005 (2018)

Semiconductor microcavities : cavity polaritons

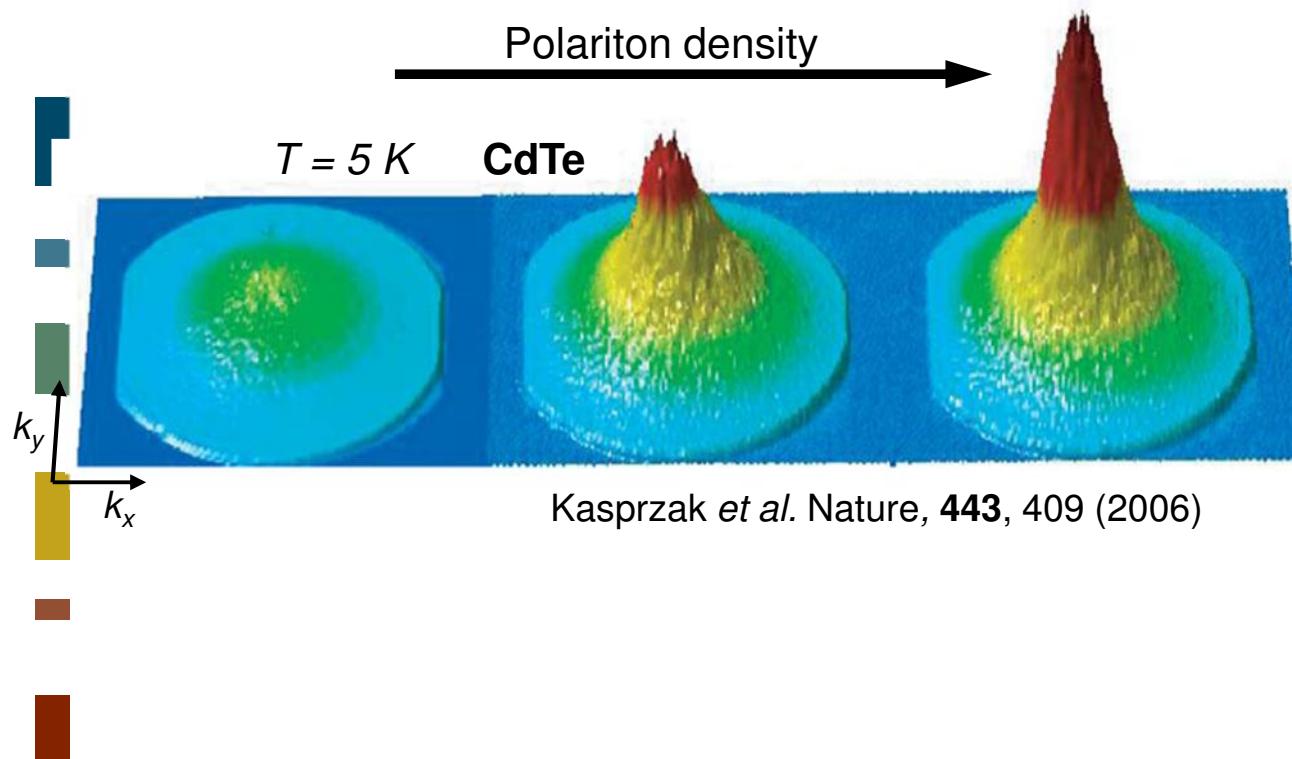


Bose-Einstein condensation of exciton polaritons

J. Kasprzak¹, M. Richard², S. Kundermann², A. Baas², P. Jeambrun², J. M. J. Keeling³, F. M. Marchetti⁴, M. H. Szymańska⁵, R. André¹, J. L. Staehli², V. Savona², P. B. Littlewood⁴, B. Deveaud² & Le Si Dang¹

T = 10 K

Nature 443, 409 (2006)



Benoit Deveaud



Le Si Dang

Polariton superfluidity



Iacopo Carusotto



Cristiano Ciuti



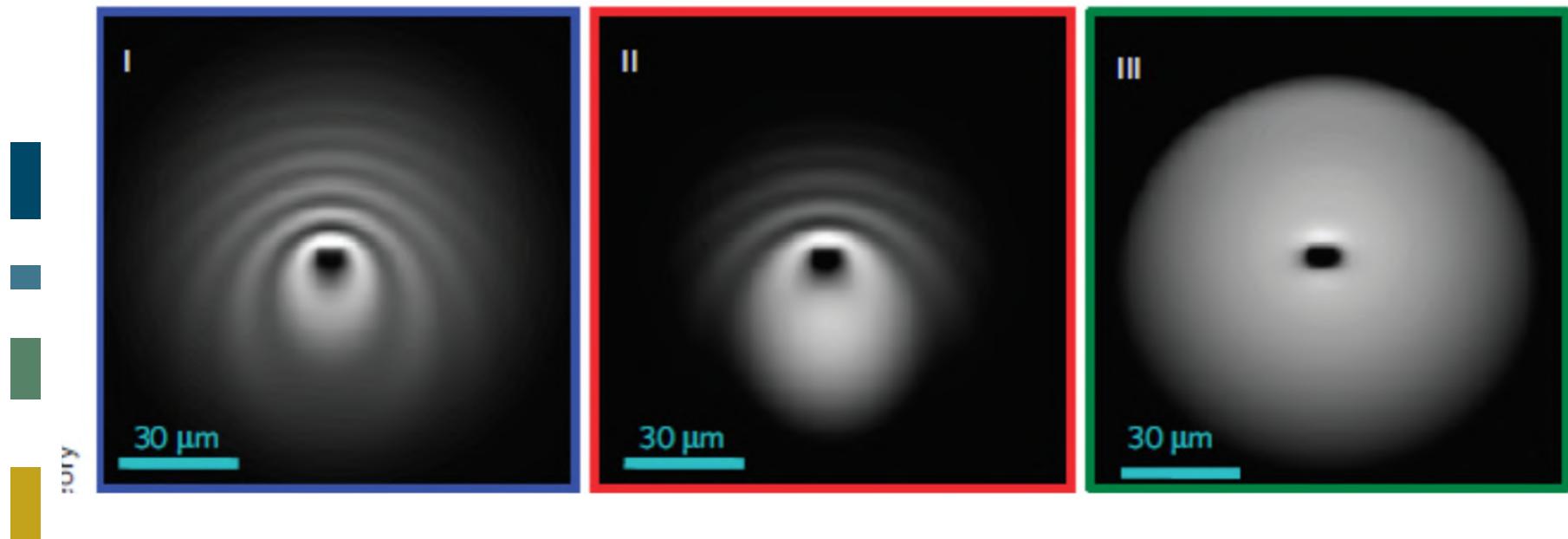
Alberto Amo



Alberto Bramati



Elisabeth Giacobino

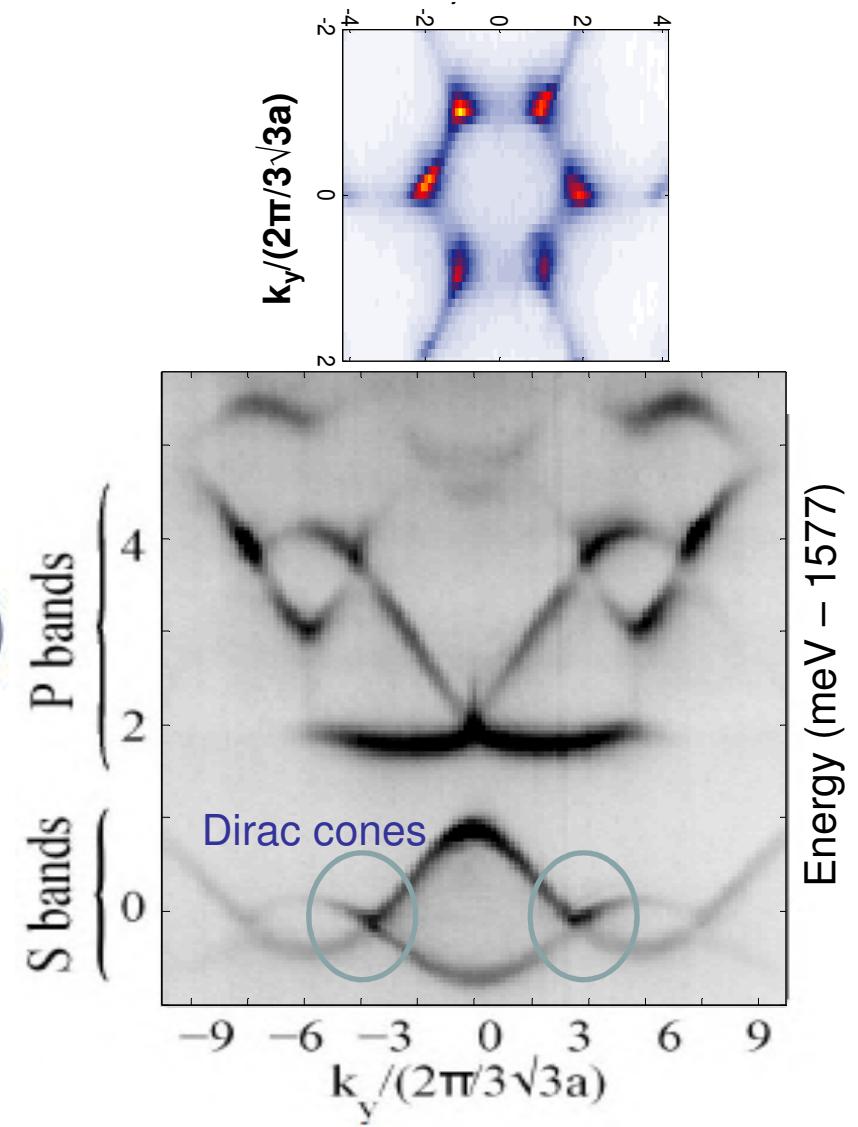
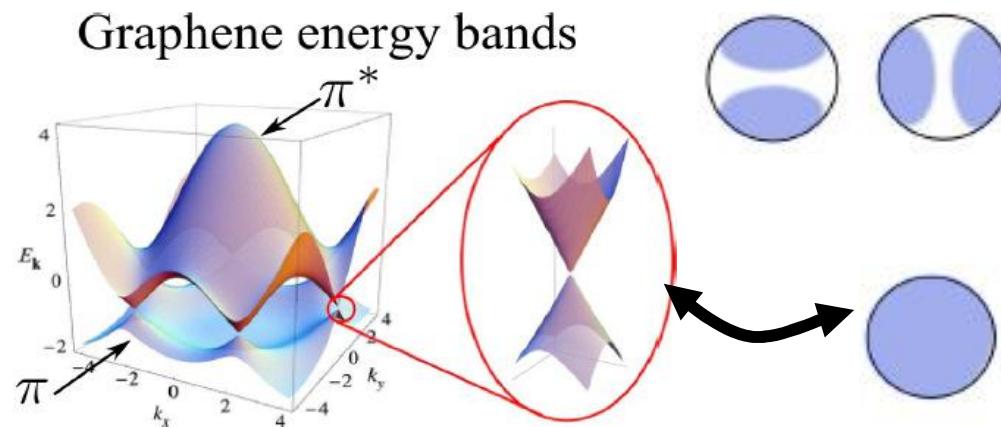
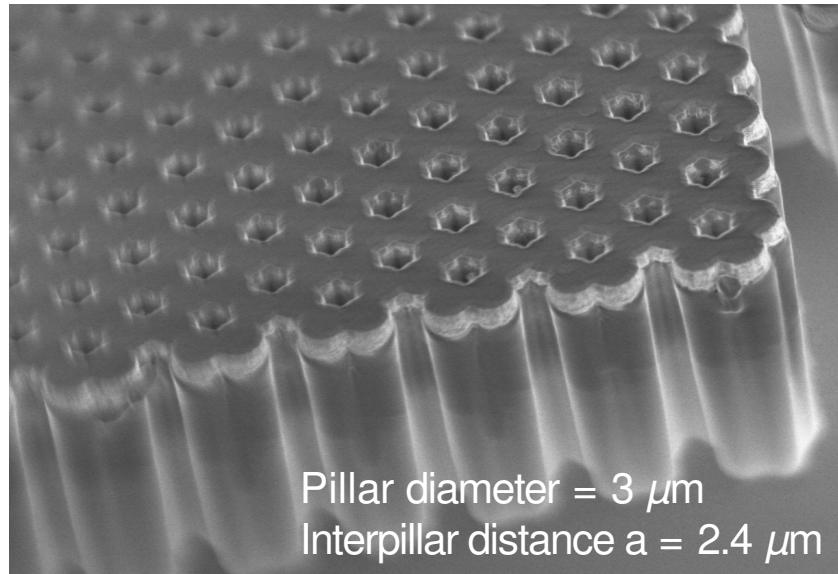


C. Ciuti and I. Carusotto PRL 242, 2224 (2005)

A. Amo et al. Nature Physics 5, 805 (2009)

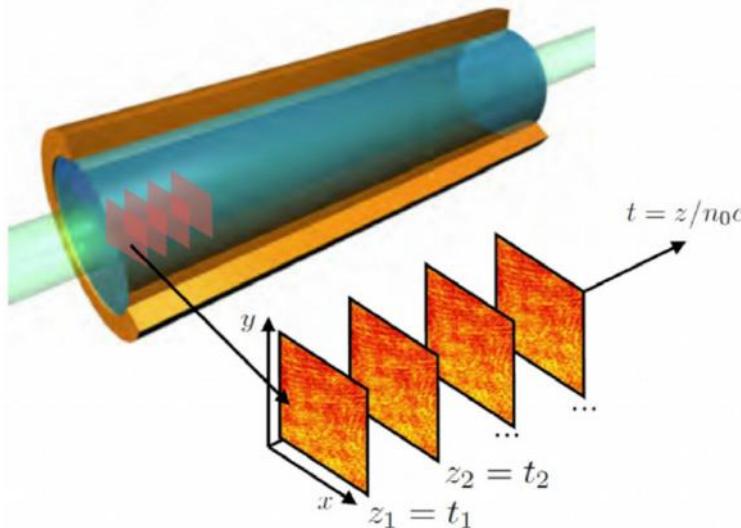
C. Ciuti & I. Carusotto, Rev. Mod. Phys. 85, 299 (2013)

Emulating Dirac Physics



Jacqmin et al., PRL 112, 116402 (2014)
See also Yamamoto (Stanford), Krizhanovskii (Sheffield)

Propagation of light in a non-linear medium



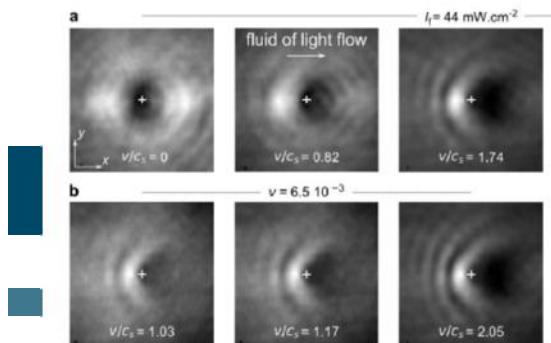
$$i \frac{\partial \mathcal{E}}{\partial z} = -\frac{1}{2k_0} \nabla_{\perp}^2 \mathcal{E} + U \mathcal{E} + \gamma |\mathcal{E}|^2 \mathcal{E}$$

kinetic energy external potential non-linear interaction

Propagation of light in a non-linear medium



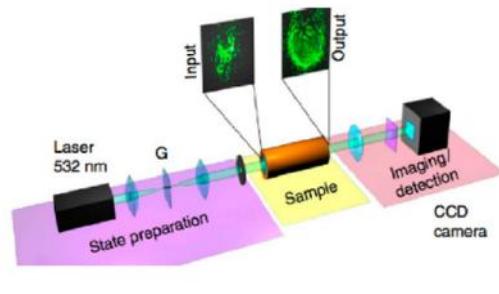
Photorefractive crystal



M. Bellec, C. Michel (INPHYNI)

C. Michel et al. Nature Com 2018

Thermo-optic liquid (methanol)



D. Faccio Edinburgh

Atomic vapor



Quentin
Glorieux

Q. Glorieux, A. Bramati (LKB)
PRL 120, 055301 (2018)

Outline

Lecture 1 : Introduction to cavity polaritons

Lecture 2: Polariton non-linearities : Superfluidity, BEC

December 10th

Lecture 3: Polariton lattices

December 12th

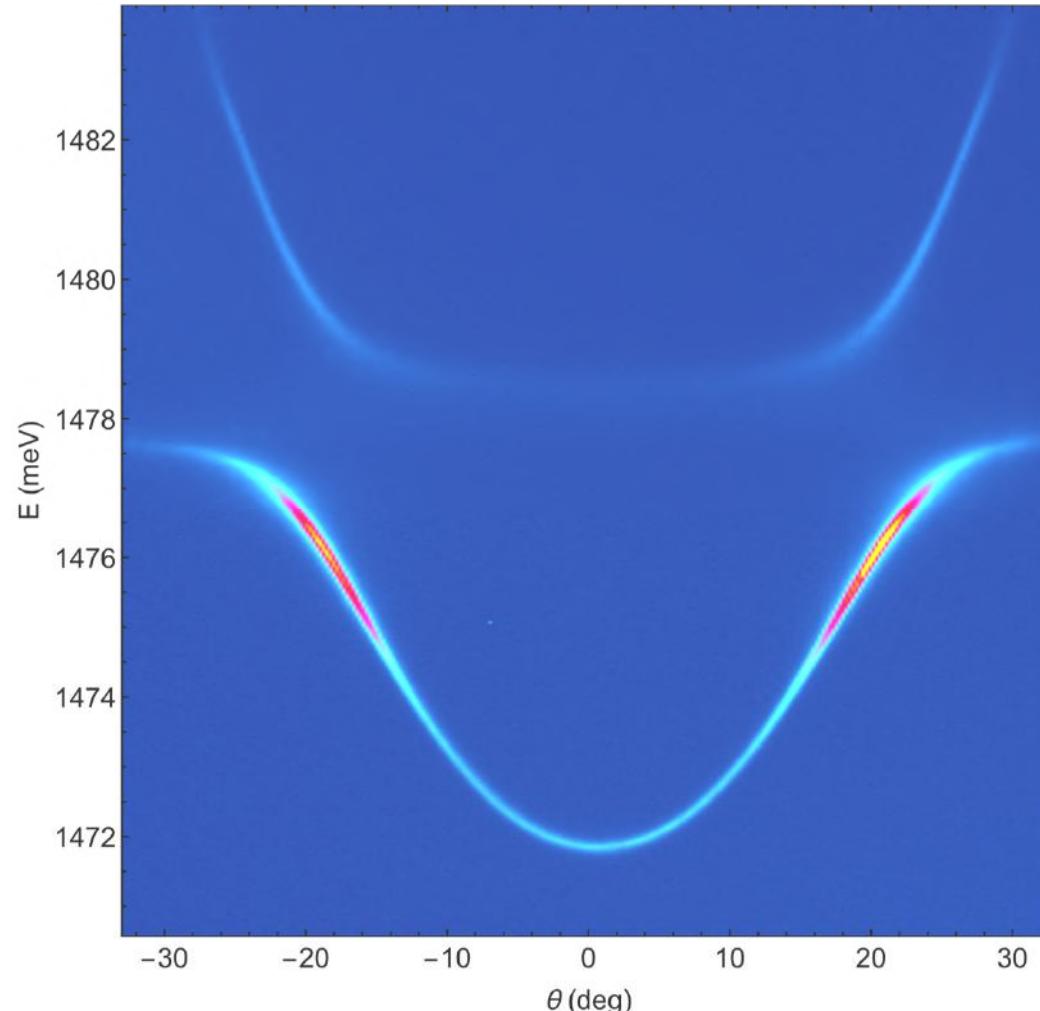
Lecture 4: Quantum fluids of light in propagating geometry

December 17th

Quentin Glorieux, LKB



Lecture 1: Introduction to cavity polaritons



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Microcavity polaritons

→ Excitonic polaritons : Mixed light-matter particles

Photons confined in an optical cavity

- **Very light ($m=0$ in vacuum)**
- Very fast
- No interactions

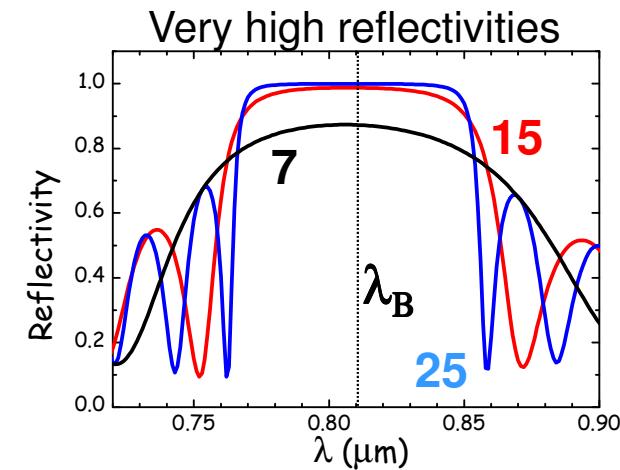
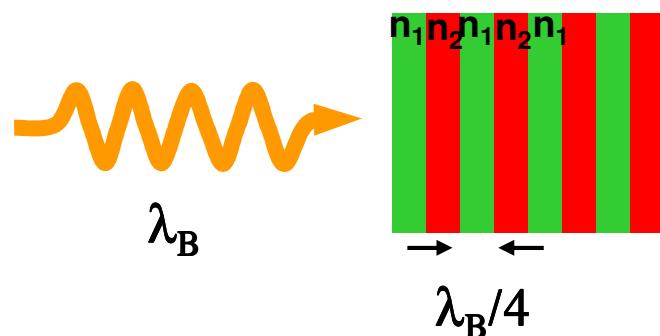
Excitons confined in a quantum well

- Very heavy
- Very slow
- **Interactions**



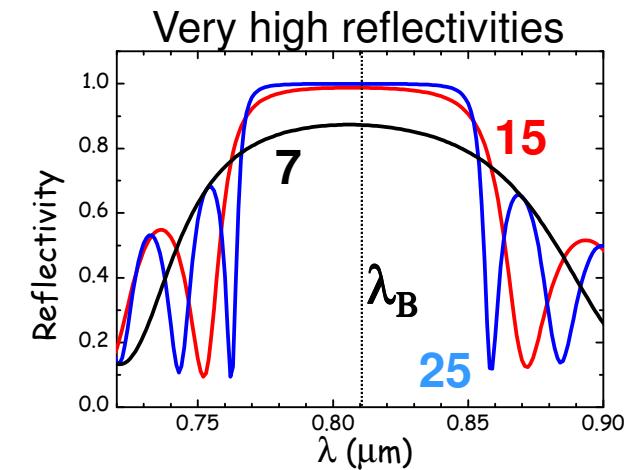
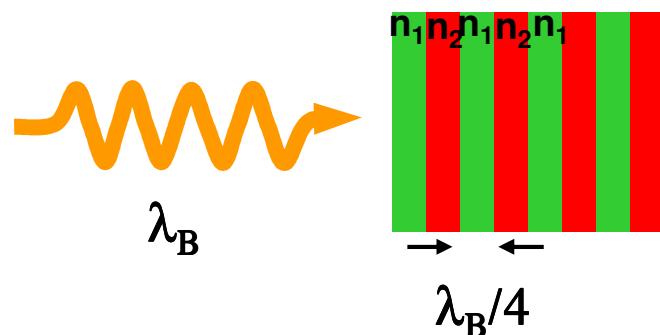
Photon confinement

→ Distributed Bragg reflector

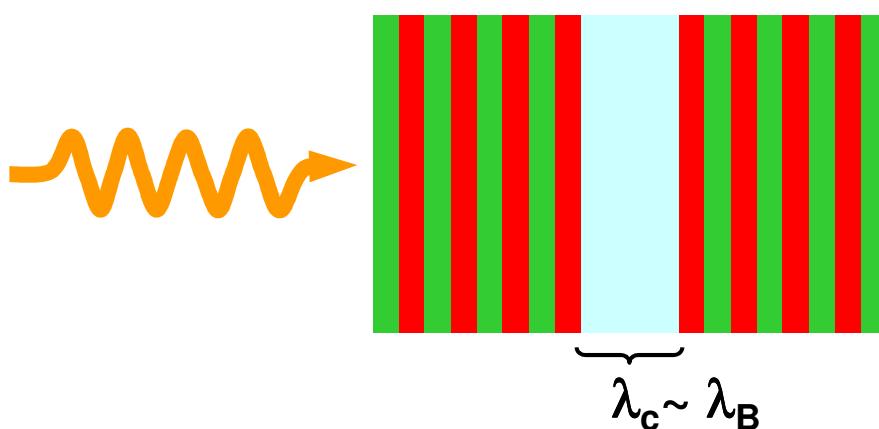


Photon confinement

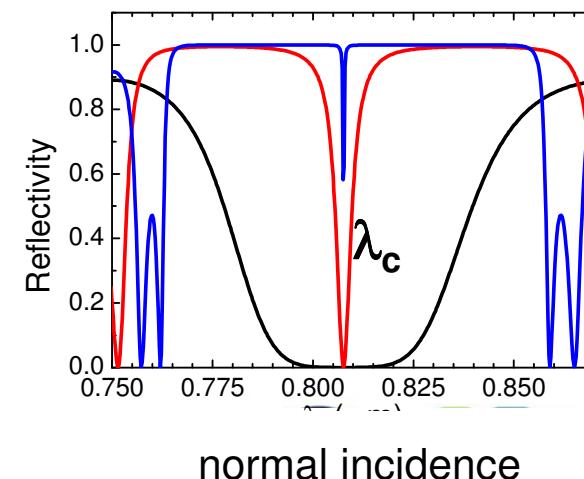
→ Distributed Bragg reflector



→ Fabry-Perot resonator (microcavity)

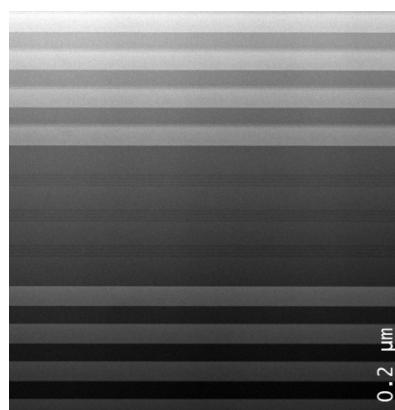
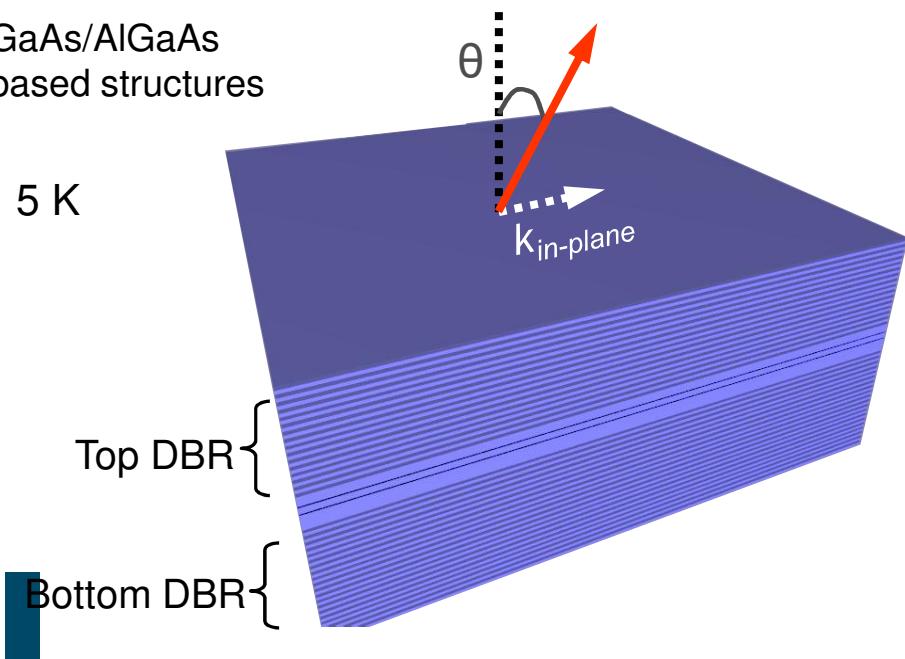


Optical mode

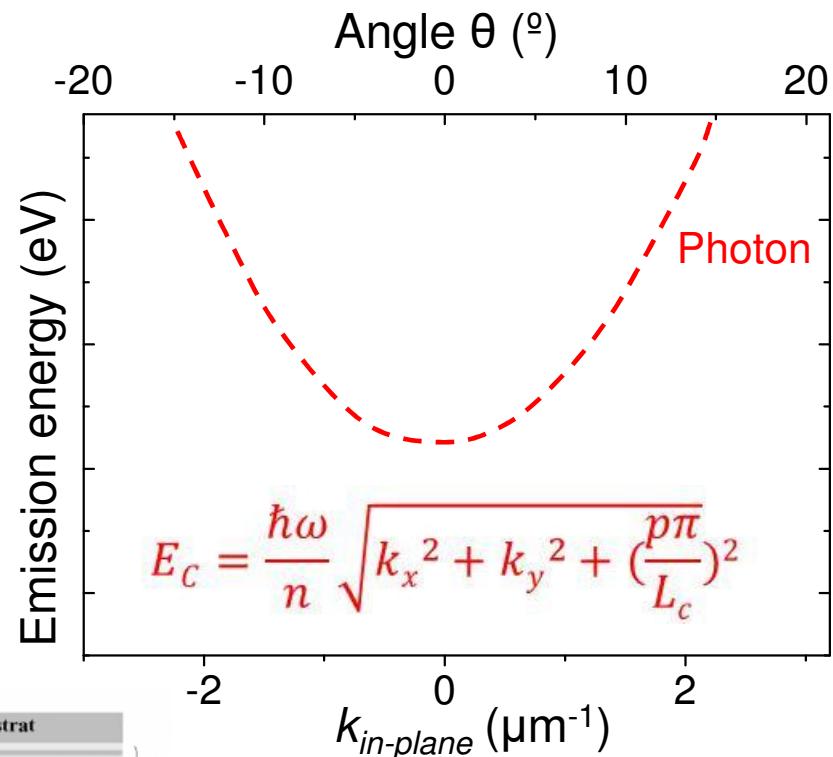


Microcavity polaritons

GaAs/AlGaAs
based structures



TEM, G. Patriarche, LPN



$$E_c(k) = E_c(k=0) + \frac{\hbar^2 k^2}{2 M p_{hot}}$$

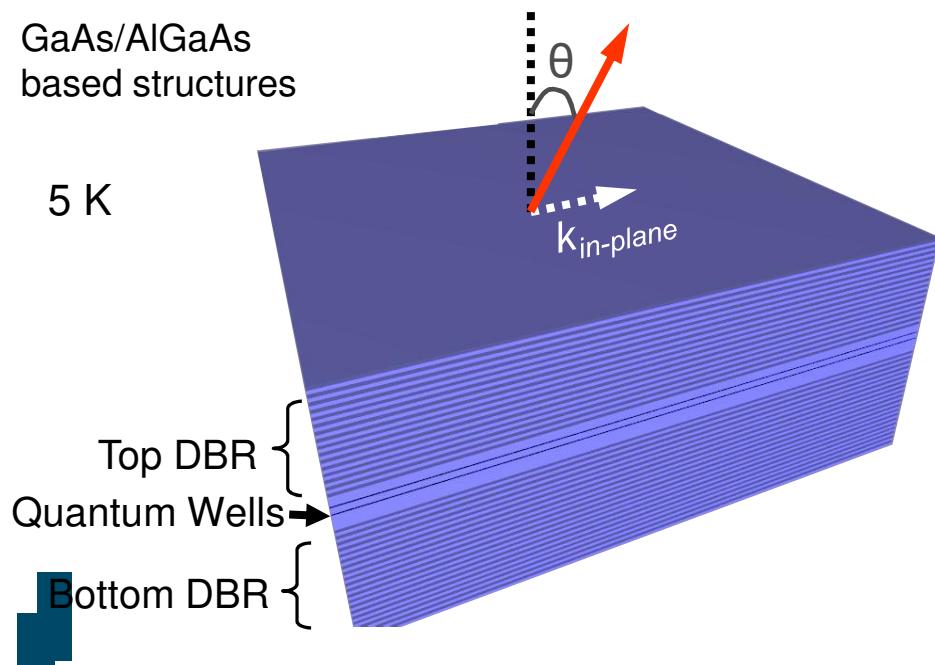
$$\text{with } M_{phot} = \frac{p^2 \pi^2 \hbar^2}{L_c^2 n^2}$$



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Microcavity polaritons

GaAs/AlGaAs
based structures

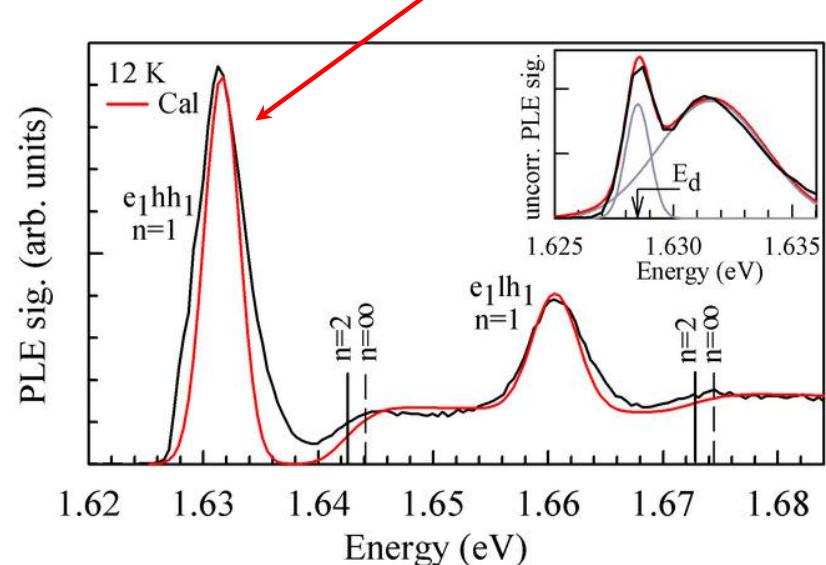


Top DBR
Quantum Wells
Bottom DBR

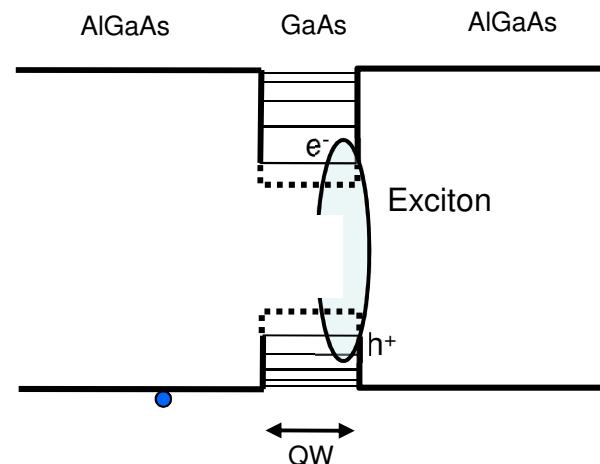
5 K



Excitonic resonance



Quantum well exciton



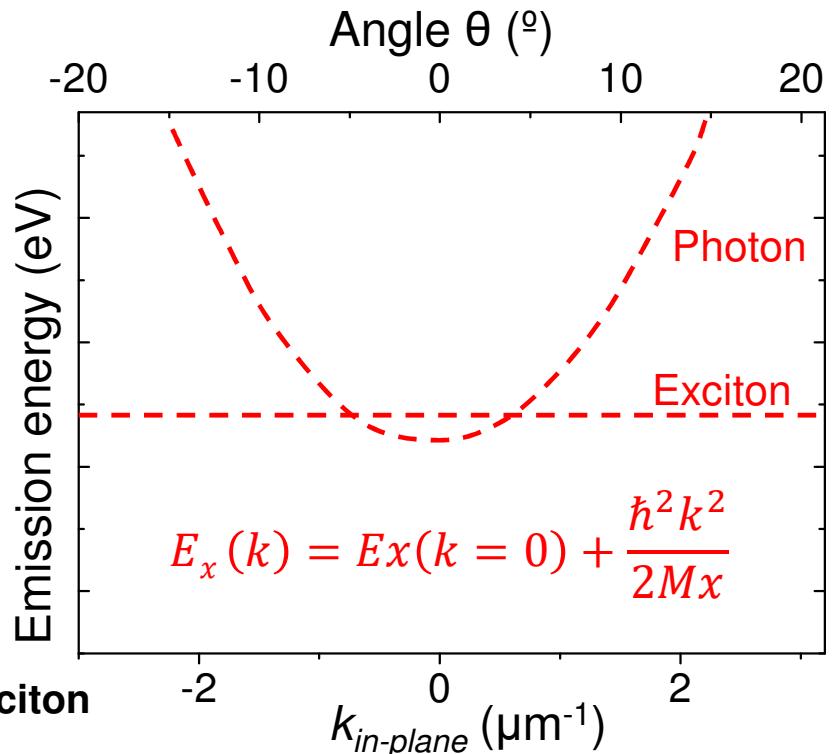
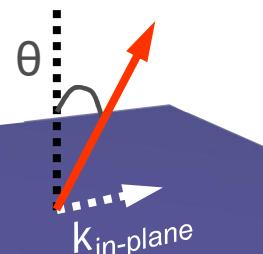
Microcavity polaritons

GaAs/AlGaAs
based structures

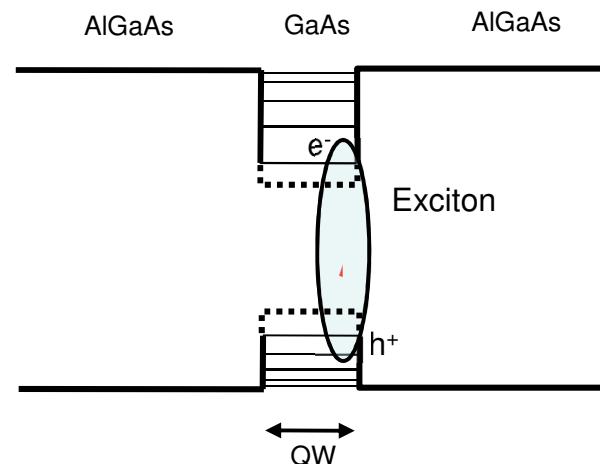
5 K

Top DBR {
Quantum Wells →

Bottom DBR {



Quantum well exciton



with $M_x = m_e + m_h$

Typically $\frac{M_x}{M_{phot}} = 10^4$



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Microcavity polaritons

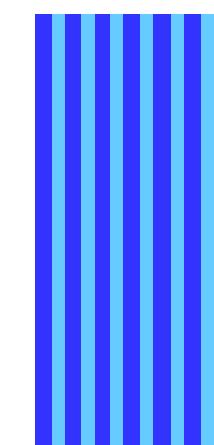
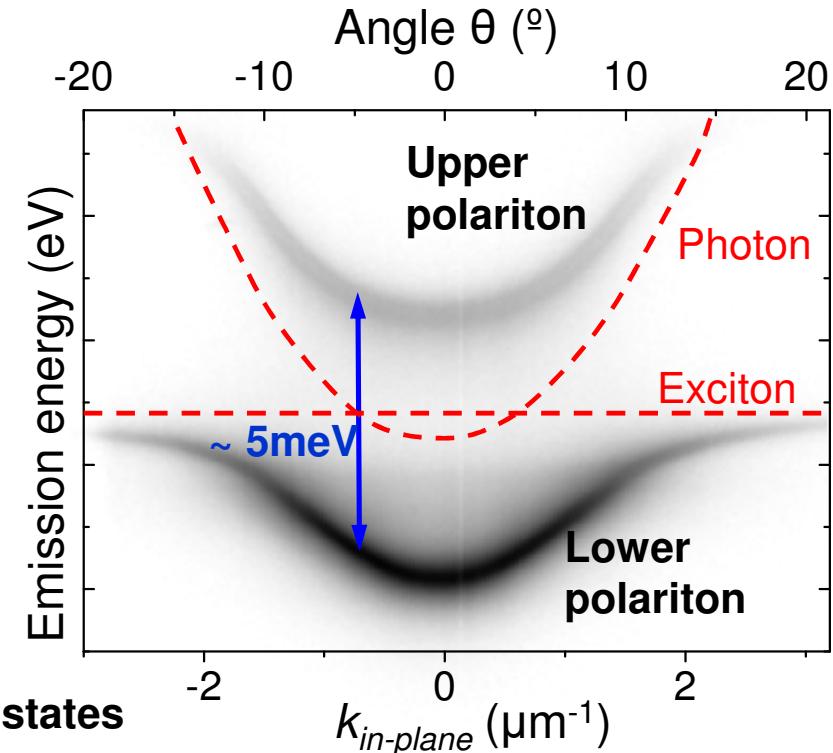
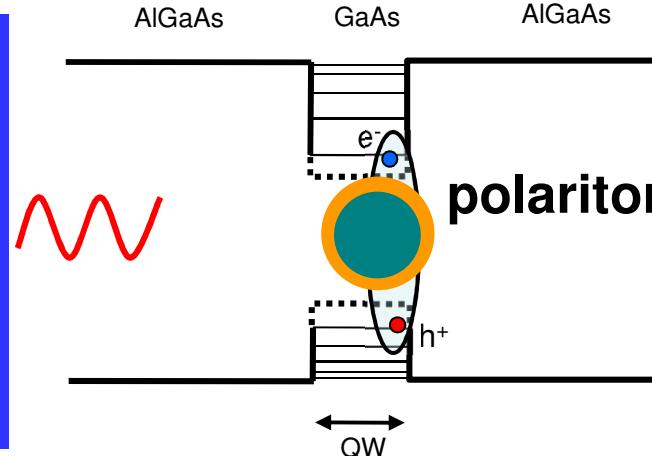
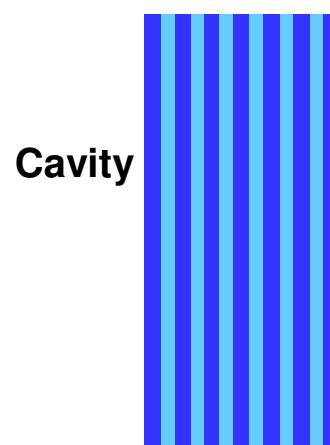
GaAs/AlGaAs
based structures

5 K

Top DBR
Quantum Wells

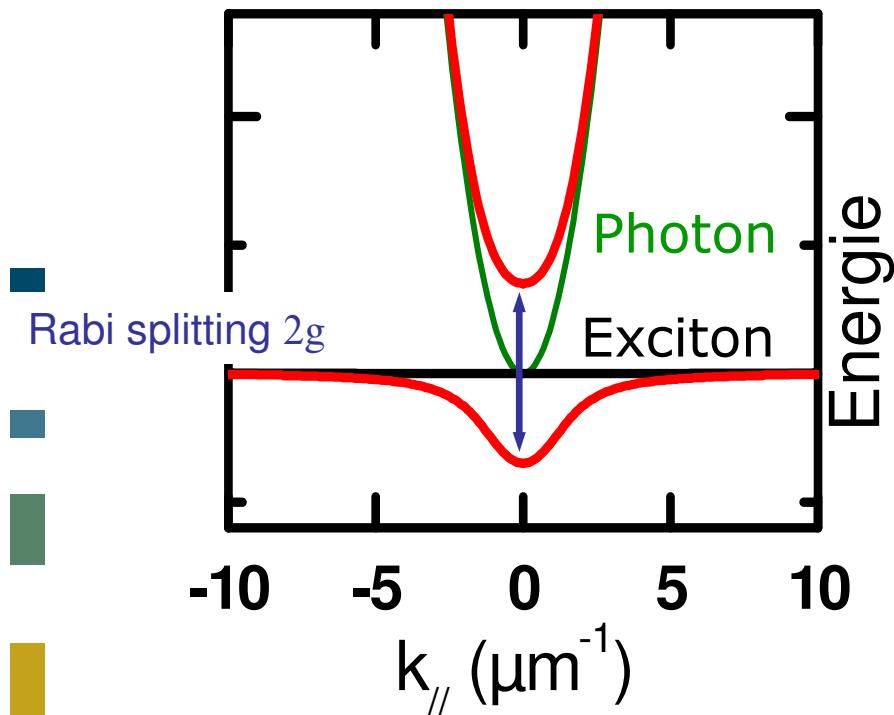
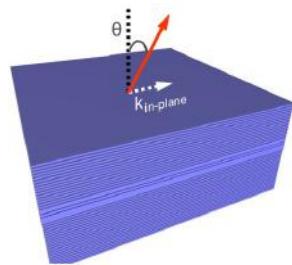
Bottom DBR

Microcavity polaritons : mixed exciton-photon states



Claude Weisbuch
PRL 69, 3314 (1992)

Microcavity polaritons



$$H_{k_{\parallel}} = \begin{pmatrix} E_X(k_{\parallel}) & g \\ g & E_C(k_{\parallel}) \end{pmatrix}$$

$$E_1 = \frac{E_X(k_{\parallel}) + E_C(k_{\parallel})}{2} - \frac{\Delta(k_{\parallel})}{2}$$

$$E_2 = \frac{E_X(k_{\parallel}) + E_C(k_{\parallel})}{2} + \frac{\Delta(k_{\parallel})}{2}$$

with

$$\Delta(k_{\parallel}) = \sqrt{(E_C(k_{\parallel}) - E_X(k_{\parallel}))^2 + 4g^2}$$



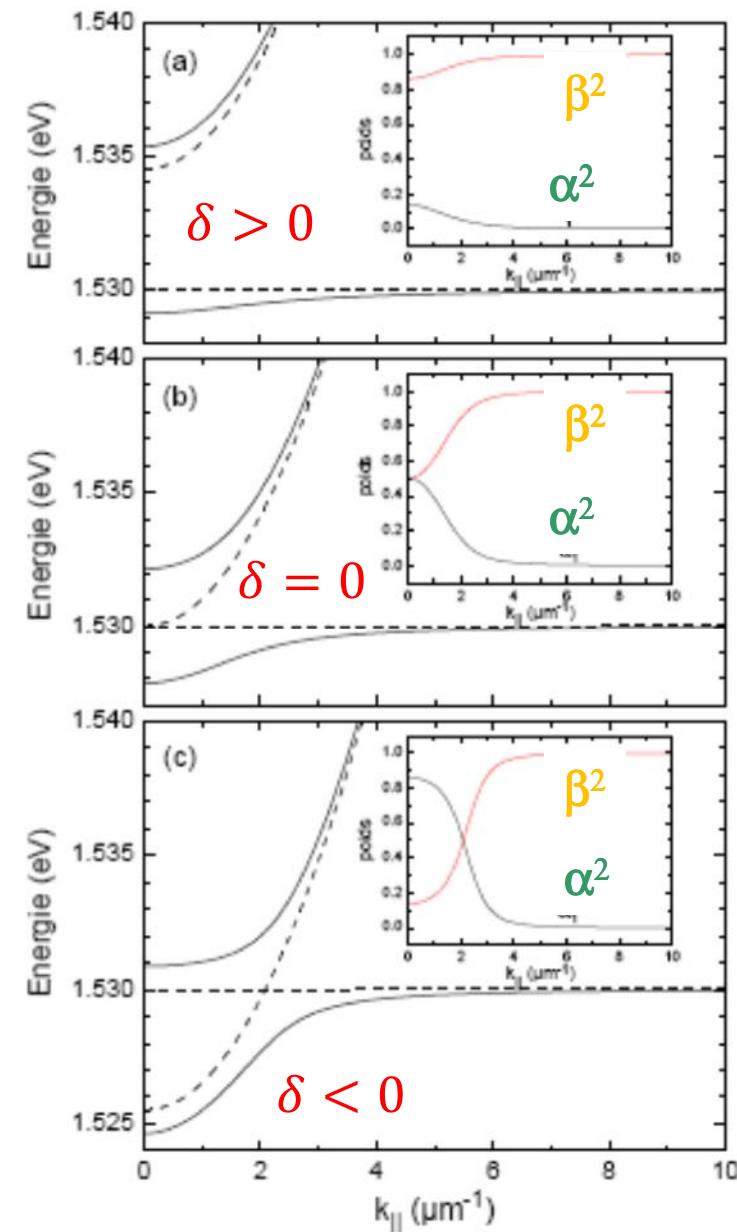
Microcavity polaritons

$$|\text{polariton}\rangle = \alpha |\text{photon}\rangle + \beta |\text{exciton}\rangle$$

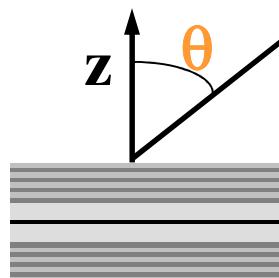
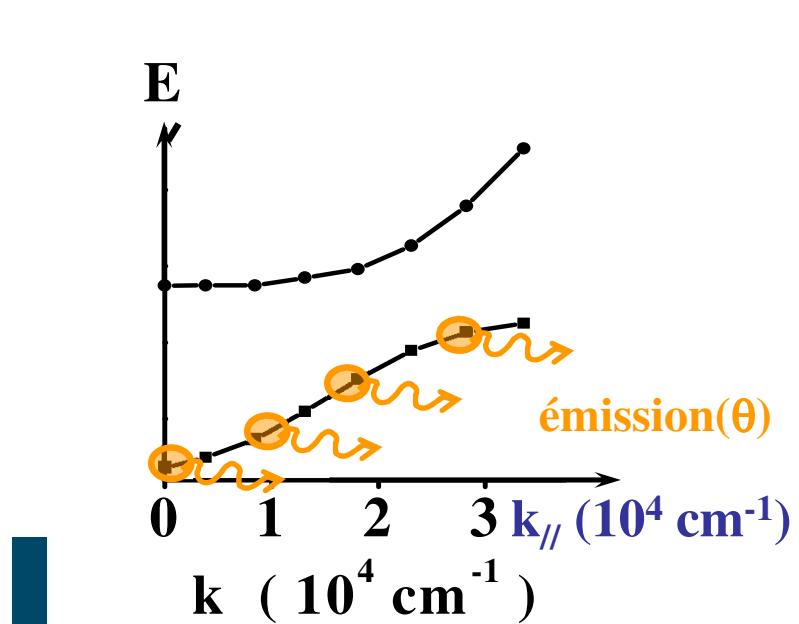
Exciton photon detuning:

$$\delta = E_c(k=0) - E_x(k=0)$$

s-shaped dispersion : inflexion point



Probing polariton states: Angle resolved experiments



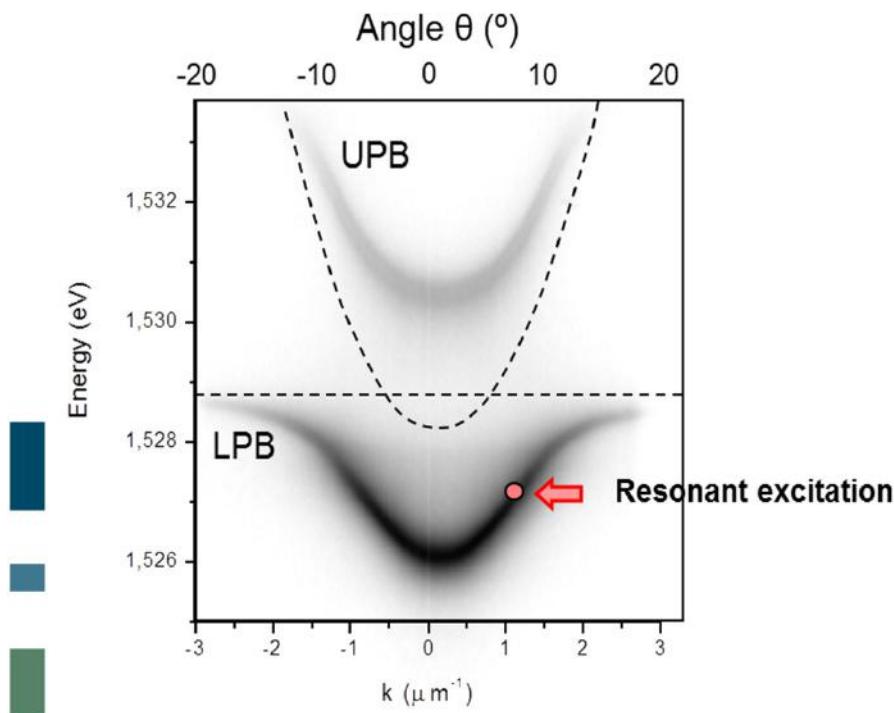
$$\theta \longleftrightarrow k_{\parallel}$$

$$k_{\parallel} = \omega/c \sin(\theta)$$

Selective excitation and probe of polariton states

How to generate microcavity polaritons?

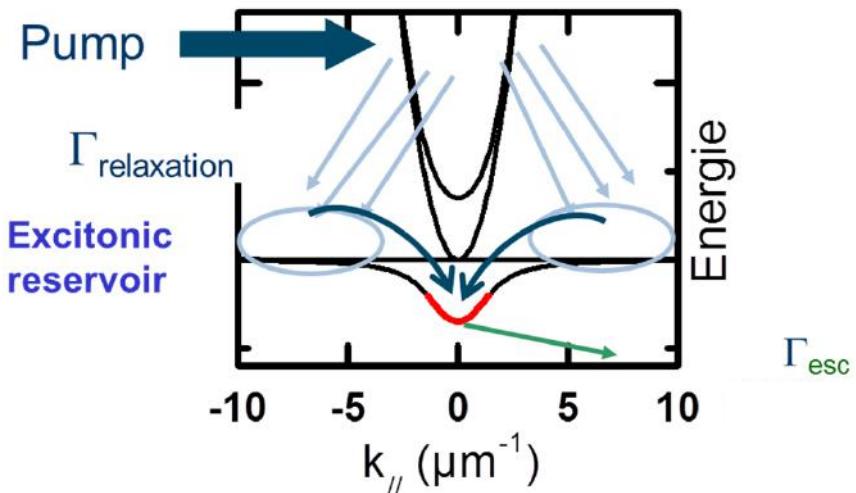
Resonant excitation



Control of :

- Energy
- k_{\parallel}
- Phase

Non resonant excitation



- Population of many low energy modes
- Creation of an excitonic reservoir
- Possibility to enter a stimulated regime: lasing, Bose E

Measurement of Cavity-Polariton Dispersion Curve from Angle-Resolved Photoluminescence Experiments

R. Houdré,¹ C. Weisbuch,^{1,2} R. P. Stanley,¹ U. Oesterle,¹ P. Pellandini,¹ and M. Illegems¹

¹*Institut de Micro- et Optoélectronique, Ecole Polytechnique Fédérale de Lausanne, CH 1015, Lausanne, Switzerland*

²*Laboratoire de Physique de la Matière Condensée, Ecole Polytechnique, F 91128 Palaiseau, France*

(Received 11 March 1994)

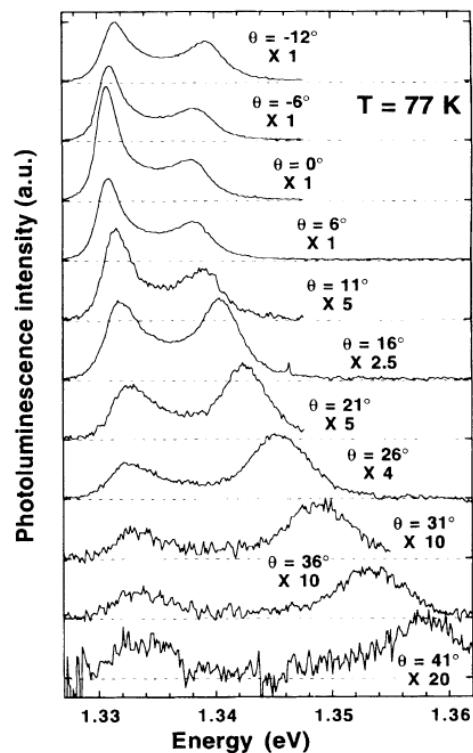


FIG. 2. Series of photoluminescence spectra at $T = 77$ K, for an emission angle from -12° to 41° . The Fabry-Pérot at normal incidence is resonant with the quantum well exciton.

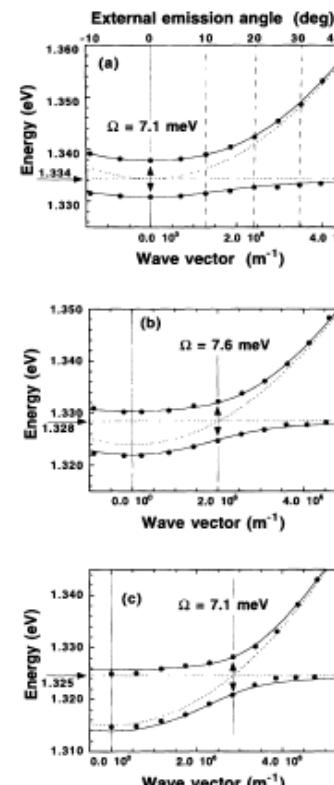


FIG. 3. Cavity-polariton dispersion curves, deduced from angle-resolved photoluminescence measurements, for different resonance conditions. (a) Resonance at $\theta = 0^\circ$ (case of Fig. 2), (b) resonance at $\theta = 29^\circ$, and (c) $\theta = 35^\circ$. The continuous lines are theoretical calculations and the dashed lines are the uncoupled exciton and cavity dispersion curves. The interaction energy Ω and exact resonance position are determined from the minimum splitting between both photoluminescence lines. An external emission angle grid is drawn on (a).

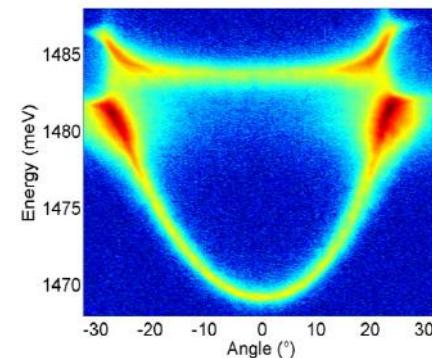
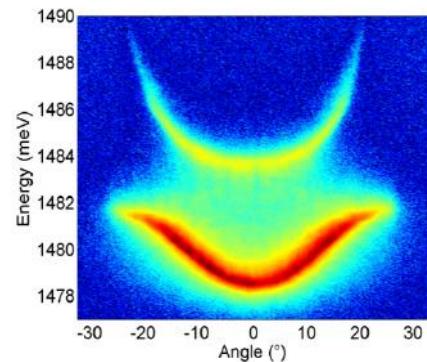
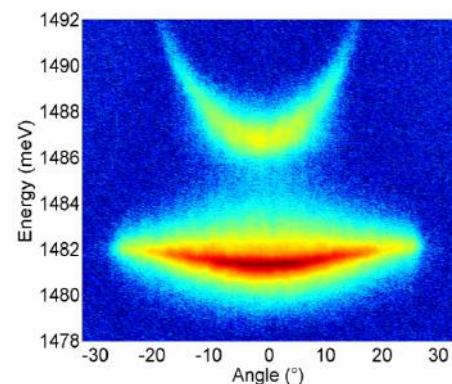
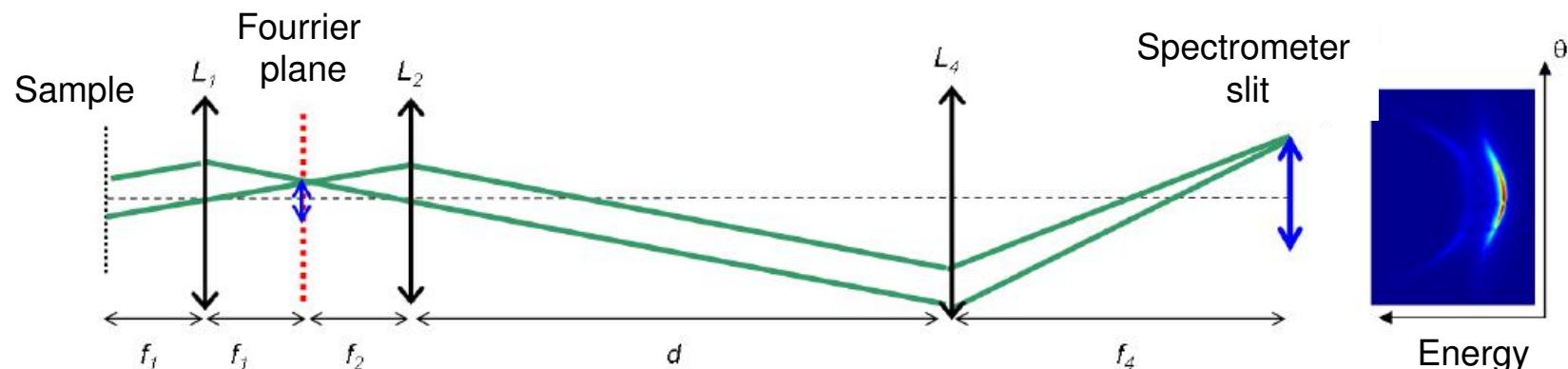
Phys. Rev. Lett. **73**, 2043 (1994)



Romuald Houdré

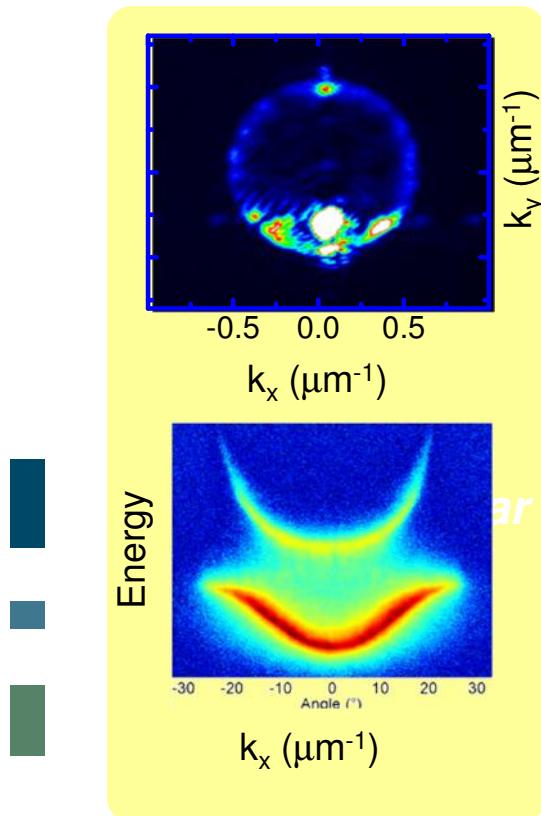
EPFL

Probing polariton states: Angle resolved experiments

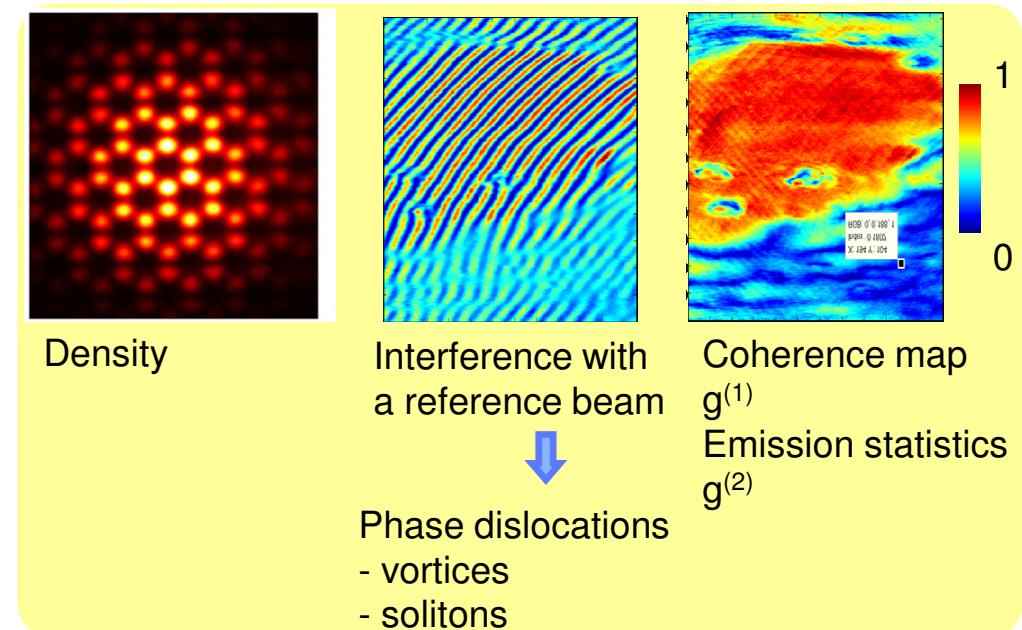


Typical experimental scheme

Far field imaging: k space

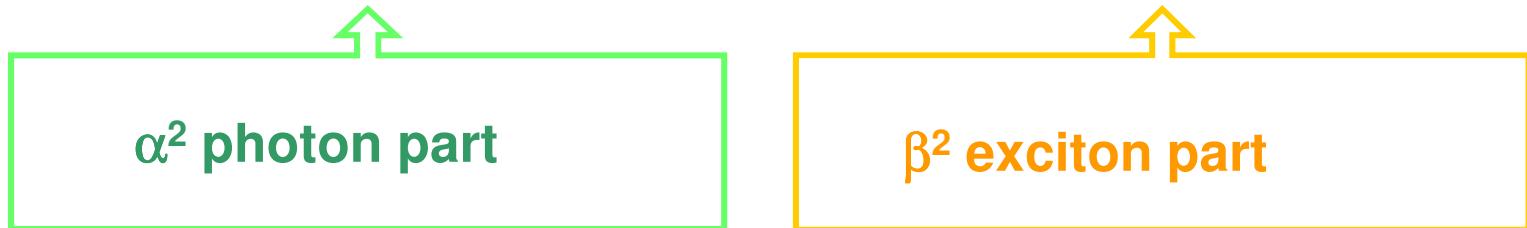


Real space imaging



Cavity polaritons : an exciton-photon mixed state

$$| \text{polariton} \rangle = \alpha | \text{photon} \rangle + \beta | \text{exciton} \rangle$$



Properties

- Photonic component \rightarrow low mass ($10^{-5} m_e$)

-
-
-
-

Microcavity polaritons

$$|\text{polariton}\rangle = \alpha |\text{photon}\rangle + \beta |\text{exciton}\rangle$$

Exciton photon detuning:

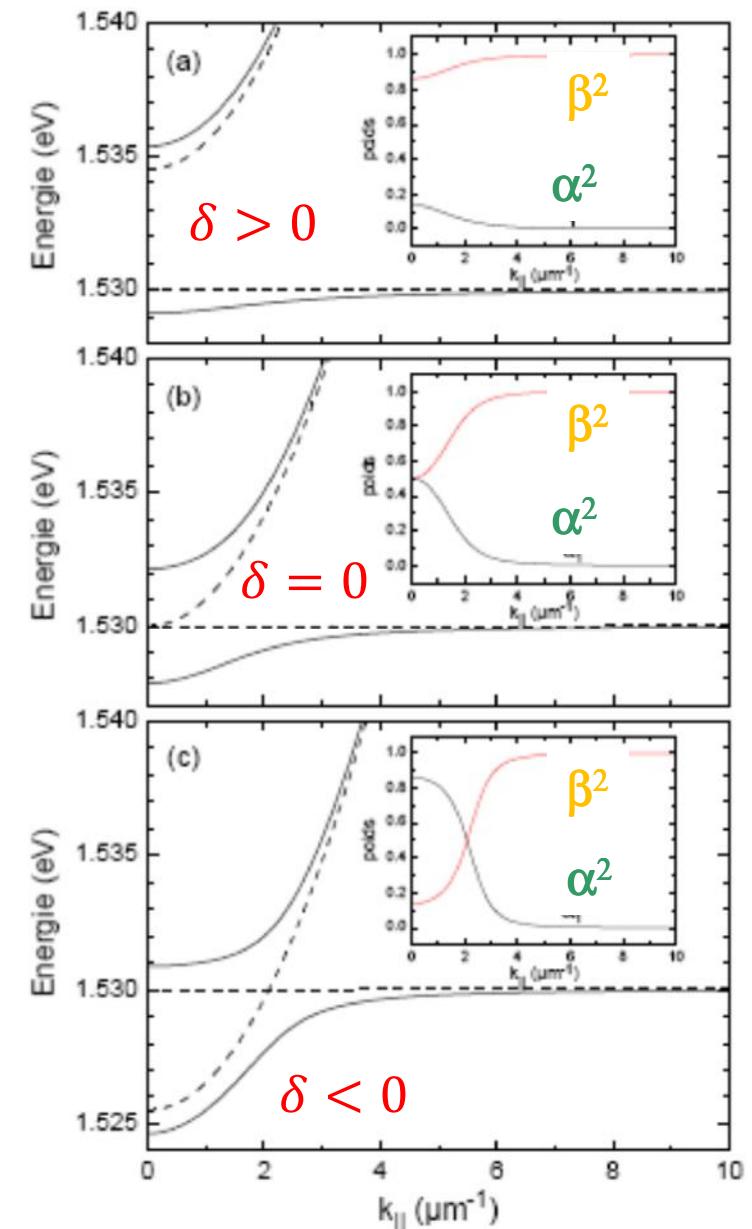
$$\delta = E_c(k=0) - E_x(k=0)$$

s-shaped dispersion : inflexion point

■ Effective mass: $\frac{1}{M_{pol}} = \frac{1}{\hbar^2} \frac{\partial^2 E}{\partial k^2}$

■ At $k=0$: $\frac{1}{M_{pol}} = \frac{\alpha^2}{M_{phot}} + \frac{\beta^2}{M_{exc}}$

■ Beyond inflexion point: negative effective mass

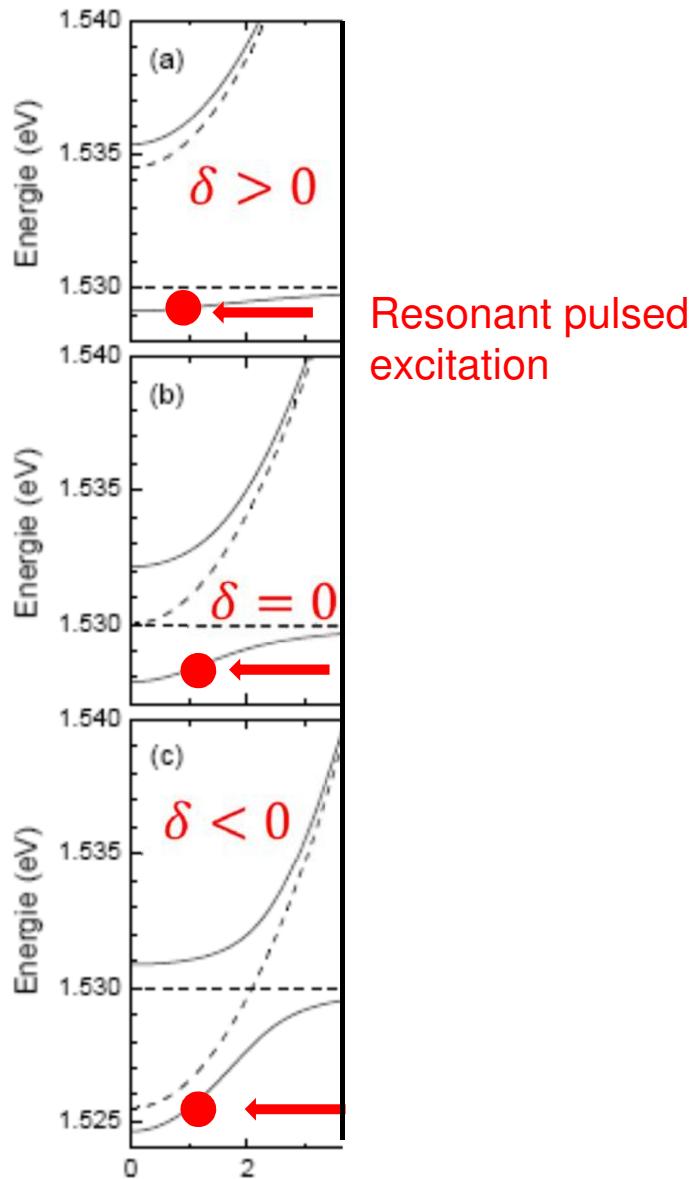


Probing polariton states: Real space propagation

PHYSICAL REVIEW B

VOLUME 61, NUMBER 11

15 MARCH 2000



Resonant pulsed excitation

In-plane propagation of excitonic cavity polaritons

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(Received 16 November 1999)

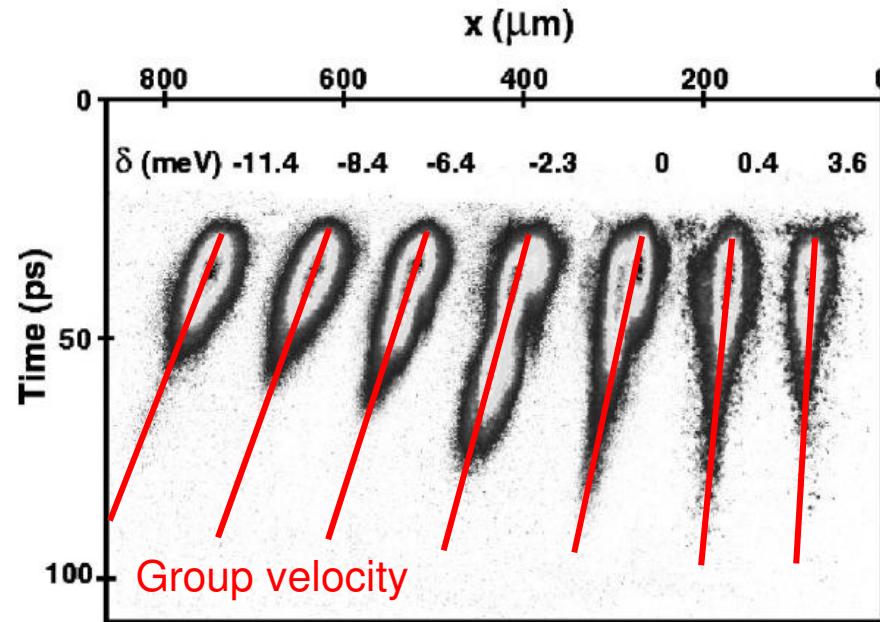
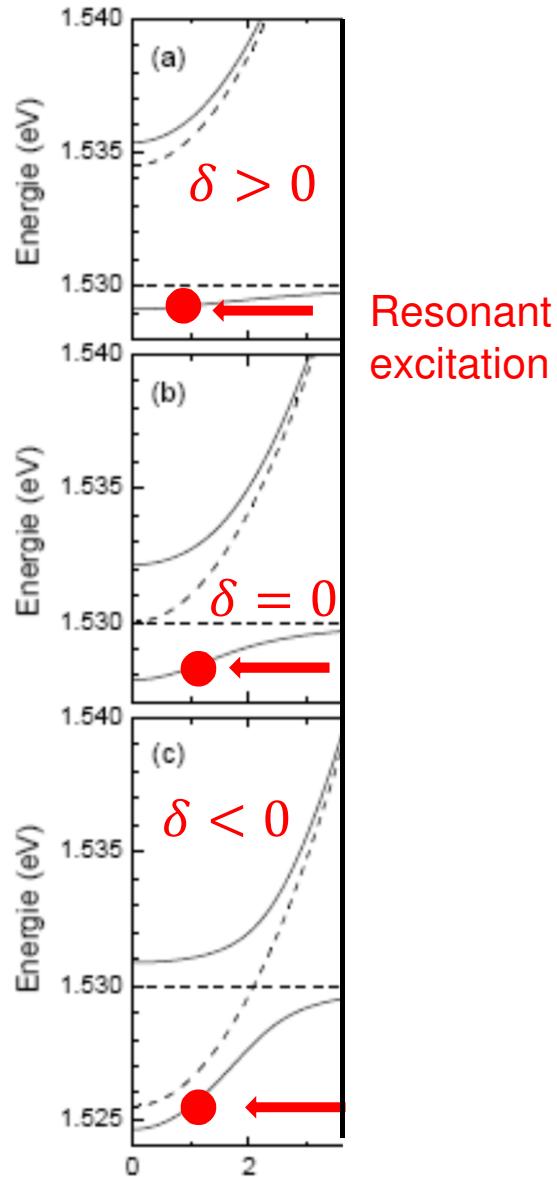


FIG. 2. Propagation in real space of the secondary emission source as observed on the streak camera for several detunings. The diaphragm is close to the reflection and $\theta_i = 7.8^\circ$.

Probing polariton states: Real space propagation



$$v_g = \frac{1}{\hbar} \frac{\partial E}{\partial k} \quad v_{g(k)} = \alpha^2 v_{phot}(k) + \beta^2 v_{ex}(k)$$

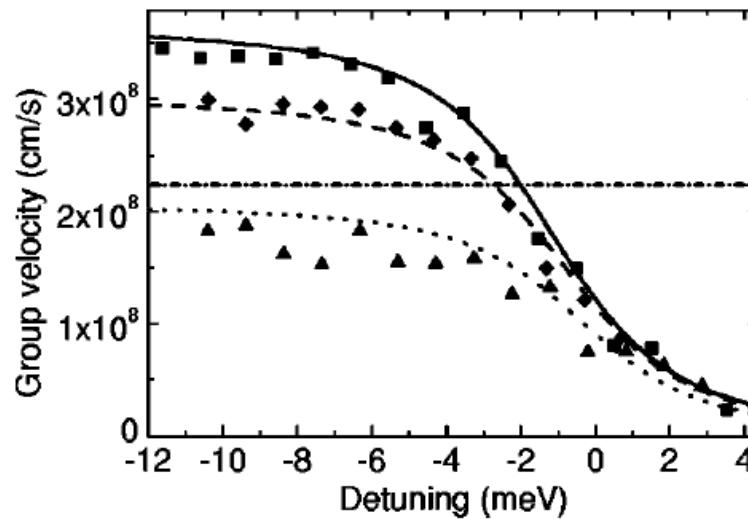
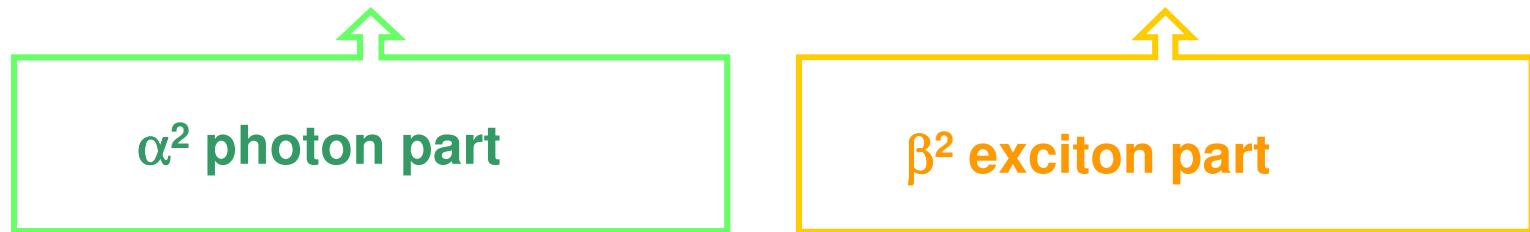


FIG. 3. Group velocity of the resonantly excited lower branch polaritons as a function of the detuning for different incident angles θ_i . The points are the measured values for $\theta_i=7.8^\circ$ (■), $\theta_i=6.2^\circ$ (◆), and $\theta_i=4.3^\circ$ (▲). The solid, dashed, and dotted lines

Cavity polaritons : an exciton-photon mixed state

$$| \text{polariton} \rangle = \alpha | \text{photon} \rangle + \beta | \text{exciton} \rangle$$



Properties

- Photonic component \rightarrow low mass ($10^{-5} m_e$)
- Short lifetime (1-100 ps) \rightarrow dissipative system
-
-

Polariton lifetime

PHYSICAL REVIEW B

VOLUME 53, NUMBER 24

15 JUNE 1996-II

Time-resolved spontaneous emission of excitons in a microcavity: Behavior of the individual exciton-photon mixed states

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FRANCE TELECOM, Centre National d'Etudes des Télecommunications, Paris B, Laboratoire de Bagneux, Boite Postale 107,
92225 Bagneux Cedex, France

J. Bloch, R. Planel, and V. Thierry-Mieg

Laboratoire de Microstructures et de Microélectronique, Centre National de la Recherche Scientifique, Boite Postale 107,
92225 Bagneux Cedex, France

(Received 1 May 1995; revised manuscript received 12 January 1996)

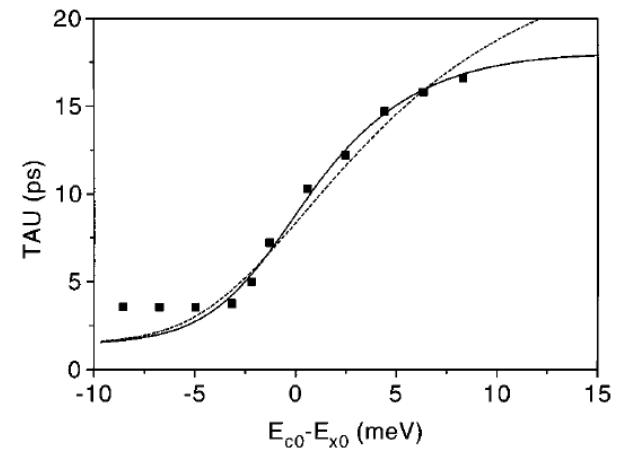
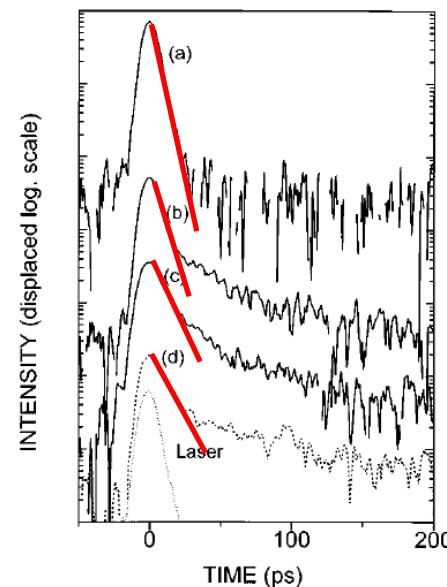
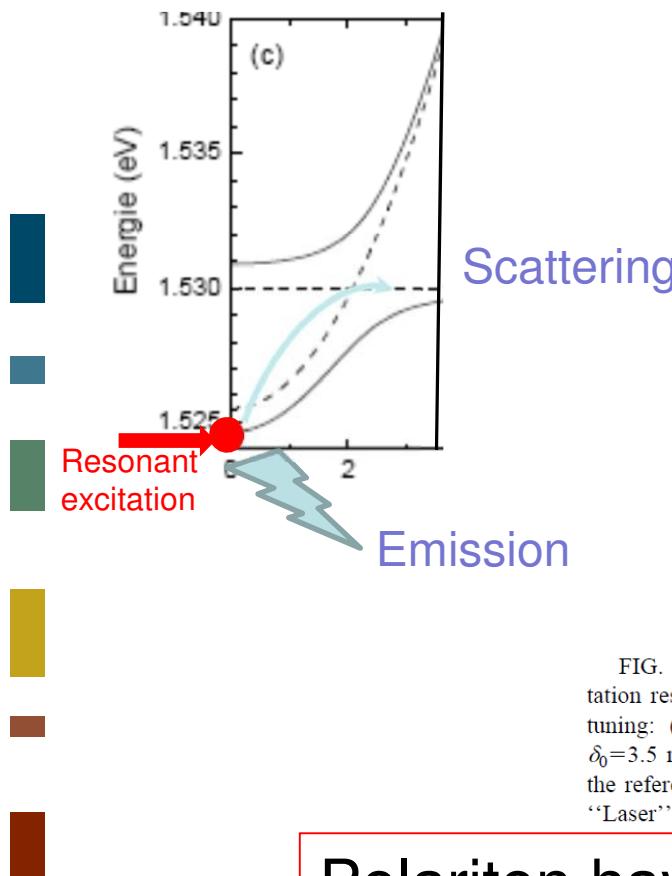


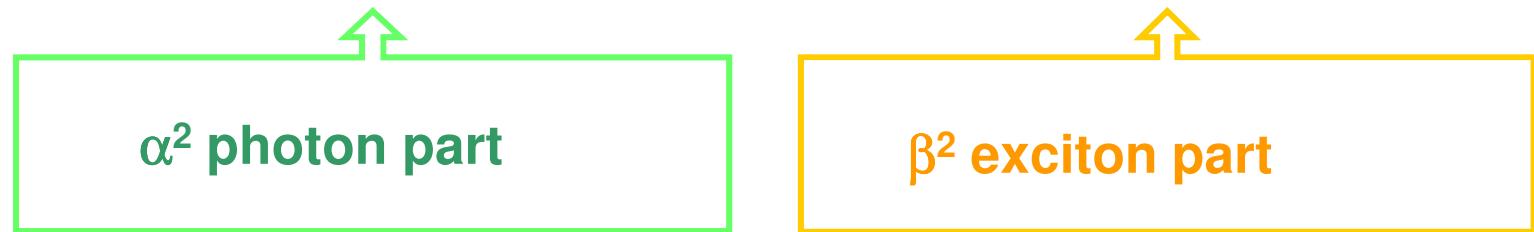
FIG. 1. Luminescence decay curves obtained with a laser excitation resonant with the lower state for different values of the detuning: (a) $\delta_0 = E_{c0} - E_{x0} = -8.5$ meV, (b) $\delta_0 = -1.5$, and (c) $\delta_0 = 3.5$ meV. The dashed curve (d) is the luminescence decay of the reference sample excited resonantly. The dotted curve (labeled "Laser") is the instrument response function.

$$\frac{1}{\tau_{pol}} = \frac{\alpha^2}{\tau_{phot}} + \frac{\beta^2}{\tau_{exc}}$$

Polariton have a finite (short) lifetime : 1-100 ps

Cavity polaritons : an exciton-photon mixed state

$$| \text{polariton} \rangle = \alpha | \text{photon} \rangle + \beta | \text{exciton} \rangle$$



Properties

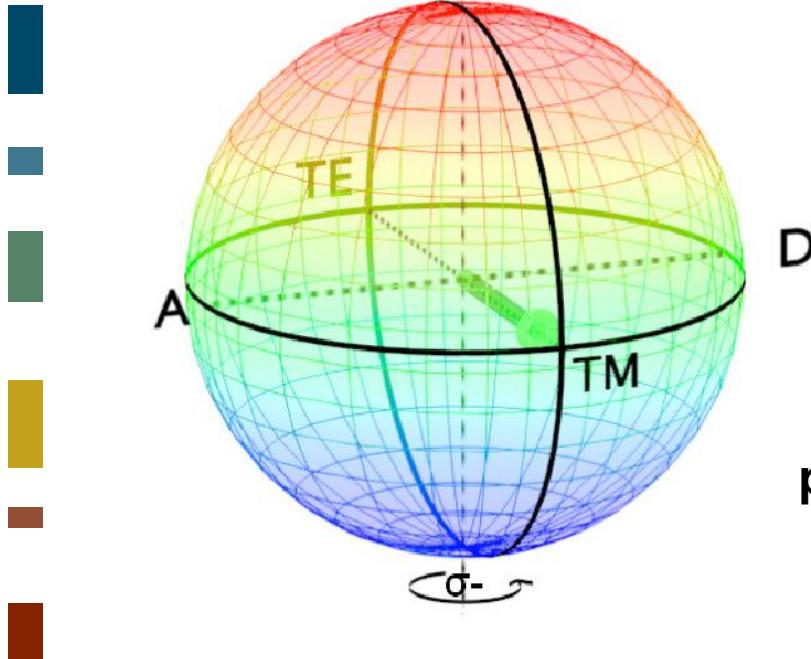
- Photonic component \rightarrow low mass ($10^{-5} m_e$)
- Short lifetime (1-100 ps) \rightarrow dissipative system
- Pseudo spin
-

Polariton spin

Spin : electron : $+- \frac{1}{2}$
heavy hole : $+- \frac{3}{2}$

Exciton : $J_z = +- 1$ $e \uparrow \downarrow h$ $e \downarrow \uparrow h$
 $J_z = +- 2$ $e \uparrow \uparrow h$ $e \downarrow \downarrow h$

Photon have an angular momentum : $+- 1$



Only $J=1$ excitons are coupled to light

Polaritons have two spin projections:

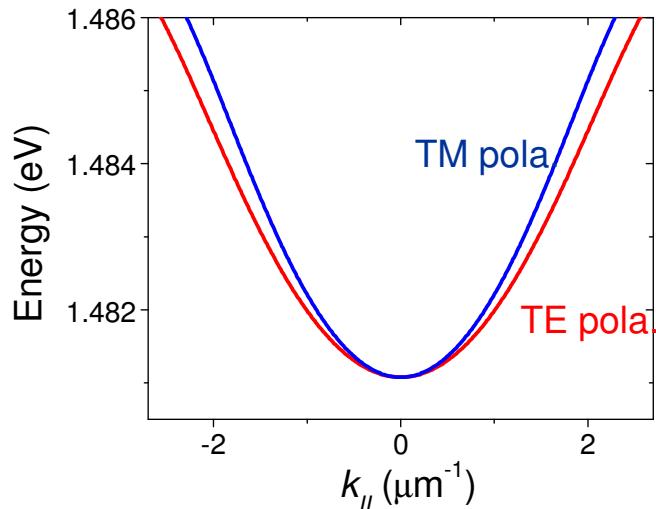
$j_z = +1$ σ^+
 $j_z = -1$ σ^-

} $\frac{1}{2}$ pseudospin

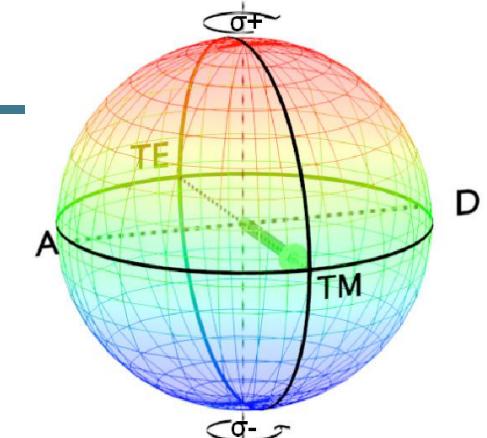
One-to-one relationship between
pseudospin state and polarisation degree

Polariton spin: Intrinsic magnetic field

Cavity modes linearly polarized: TE-TM splitting



boundary conditions for the electromagnetic field at the interface of the different dielectric layers

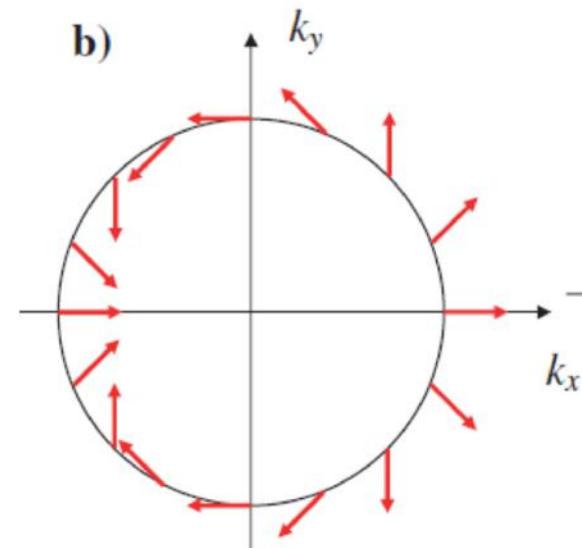


Effective magnetic field

$$\mathbf{H}_{\text{eff}} = \frac{\hbar}{\mu_B g} \boldsymbol{\Omega}_{\mathbf{k}}$$

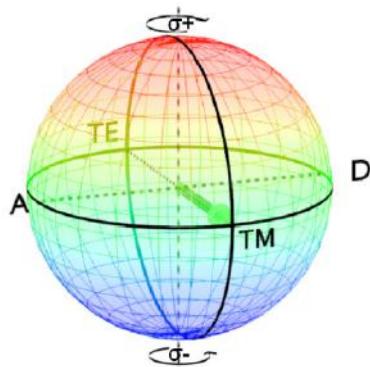
$$\Omega_x = \frac{\Omega}{k^2} (k_x^2 - k_y^2), \quad \Omega_y = 2 \frac{\Omega}{k^2} k_x k_y,$$

$$\boldsymbol{\Omega} = \frac{\Delta_{\text{LT}}}{\hbar}$$



Kavokin et al. PRL 95, 136601 (2005)

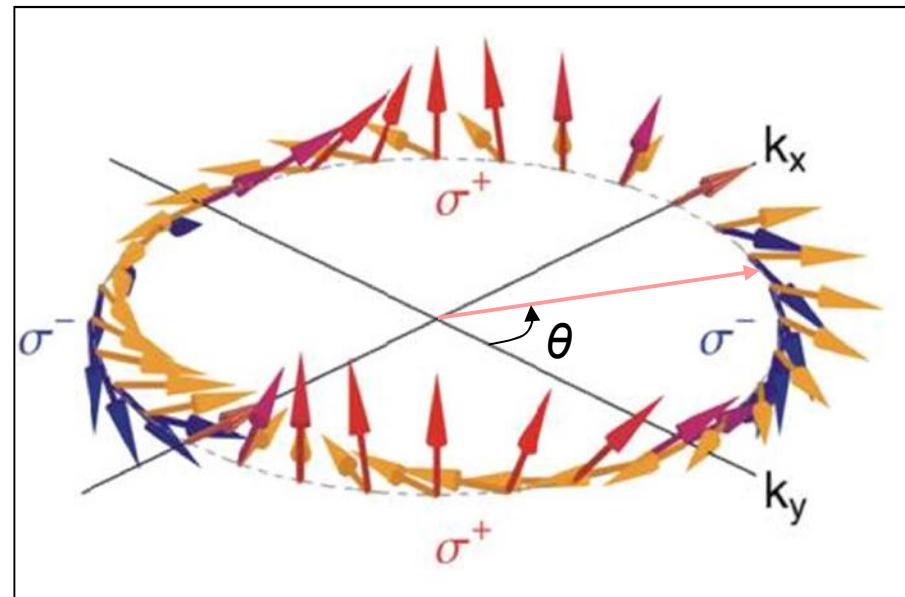
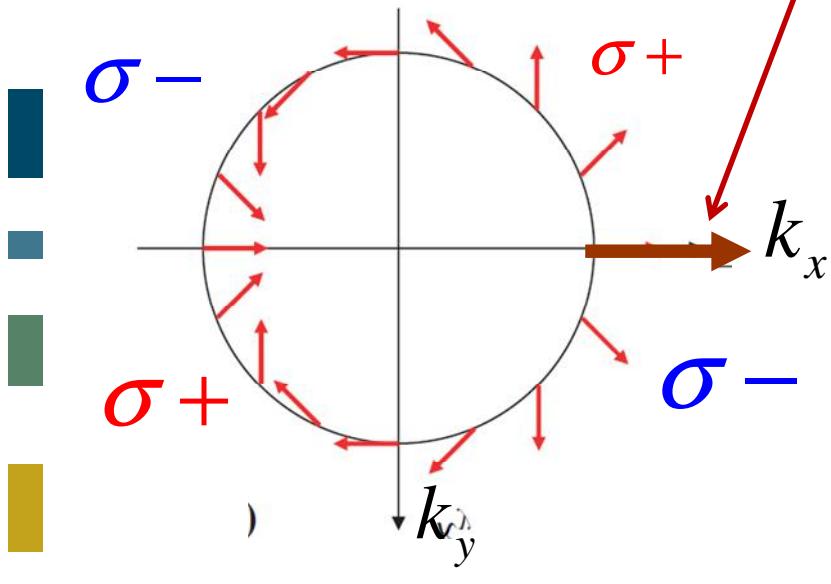
Optical spin Hall effect



Initial
Pseudospin

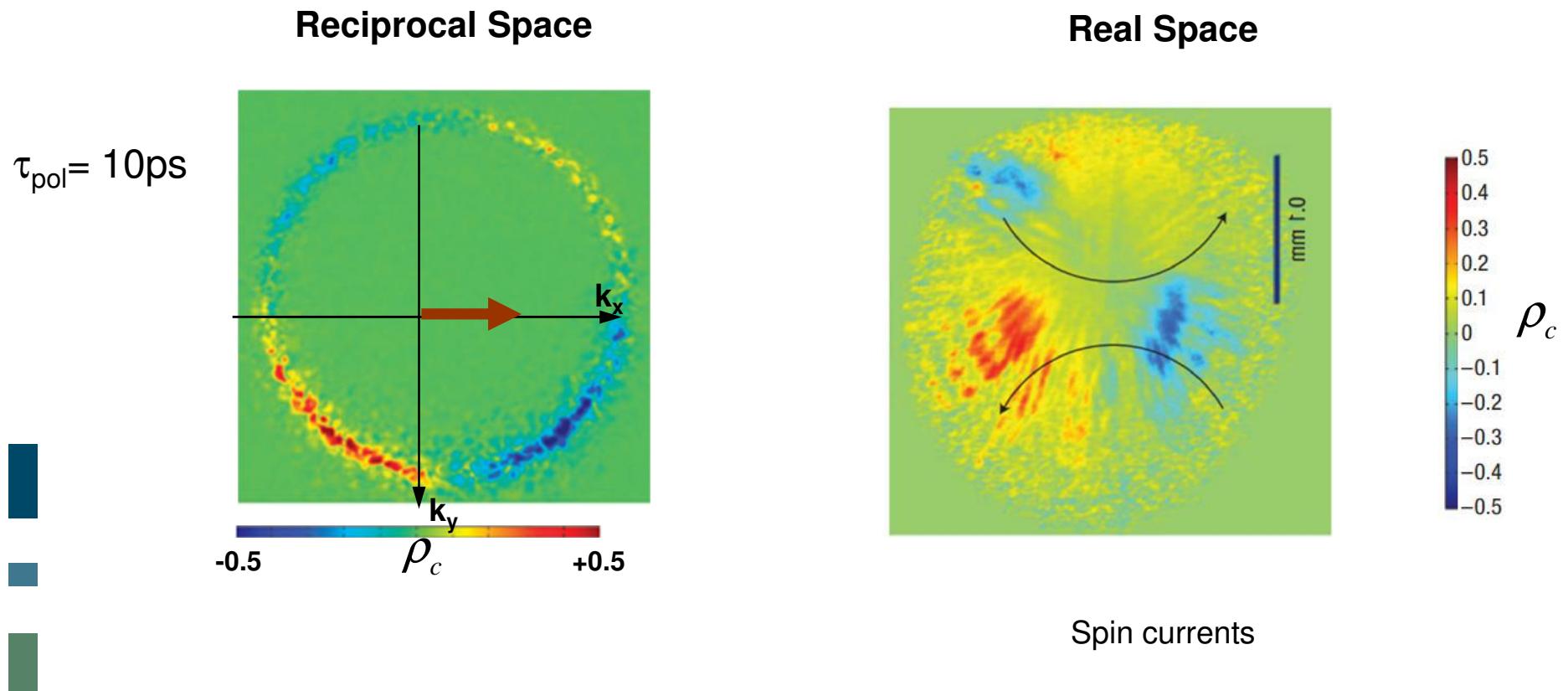
Linear polarisation along x

$$\frac{\partial \vec{S}}{\partial t} = \vec{S} \wedge \vec{\Omega}(\theta) + \frac{\vec{S}_0}{\tau_1} - \frac{\vec{S}}{\tau}$$



Kavokin et al. PRL 95, 136601 (2005)
Leyder et al. Nature Phys. 3, 628 (2007).

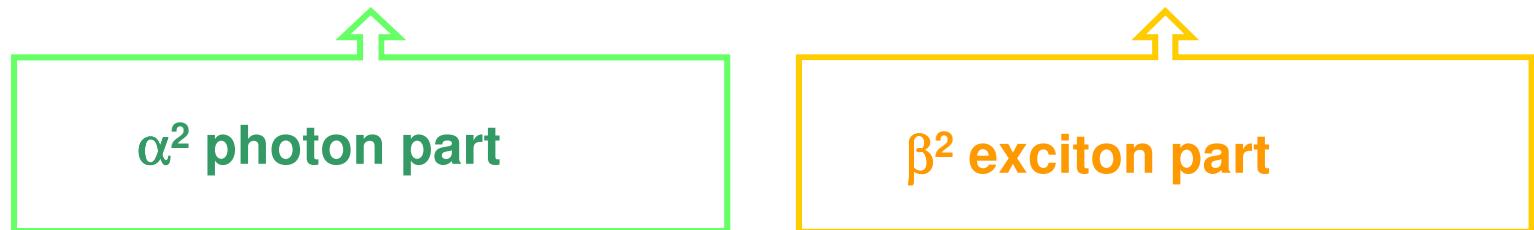
Optical spin Hall effect



Kavokin et al. PRL **95**, 136601 (2005)
Leyder et al. Nature Phys. **3**, 628 (2007).

Cavity polaritons : an exciton-photon mixed state

$$| \text{polariton} \rangle = \alpha | \text{photon} \rangle + \beta | \text{exciton} \rangle$$

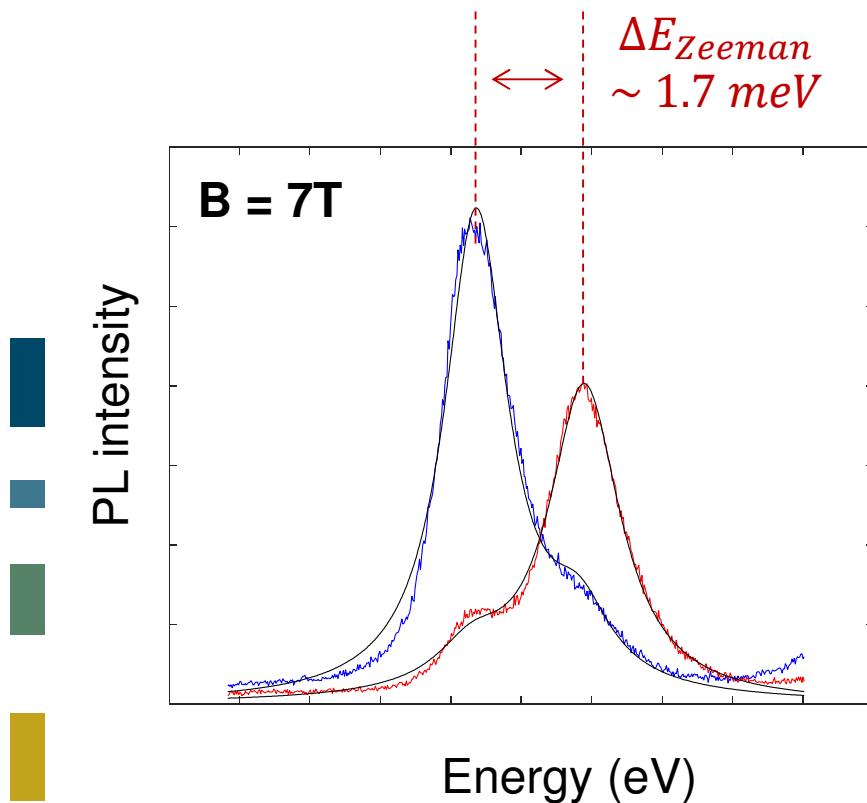


Properties

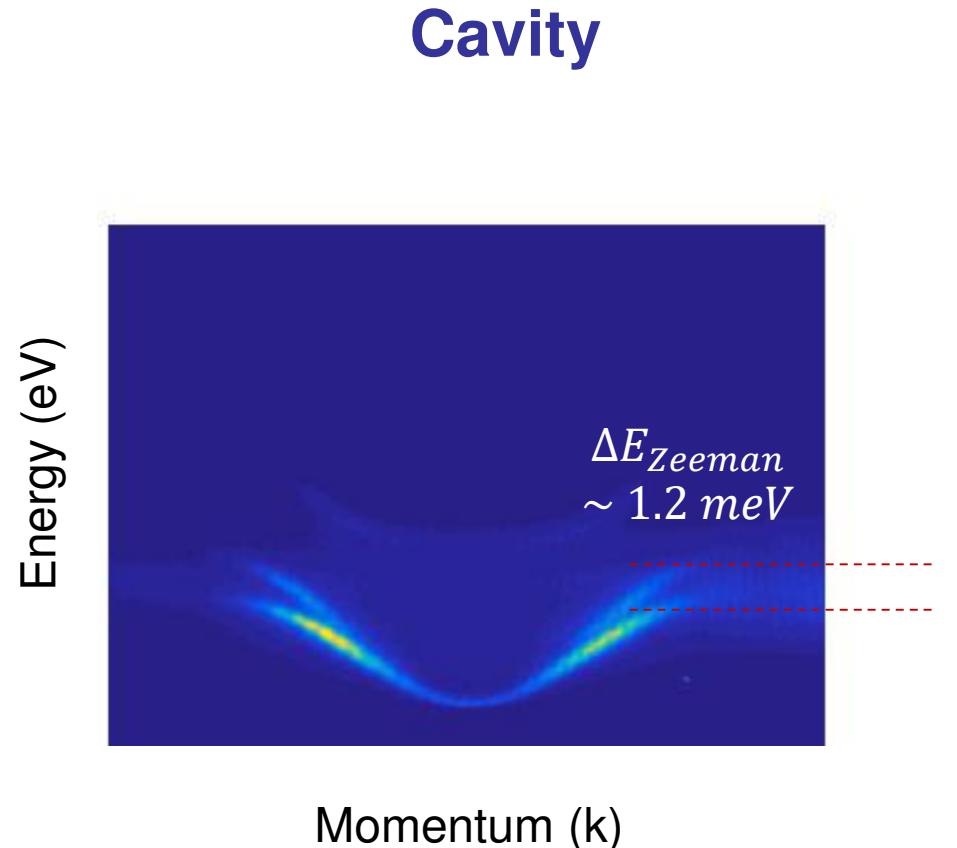
- Photonic component \rightarrow low mass ($10^{-5} m_e$)
- Short lifetime (1-100 ps) \rightarrow dissipative system
- Pseudo spin
- Sensitivity to magnetic field

Magneto-luminescence

Quantum well

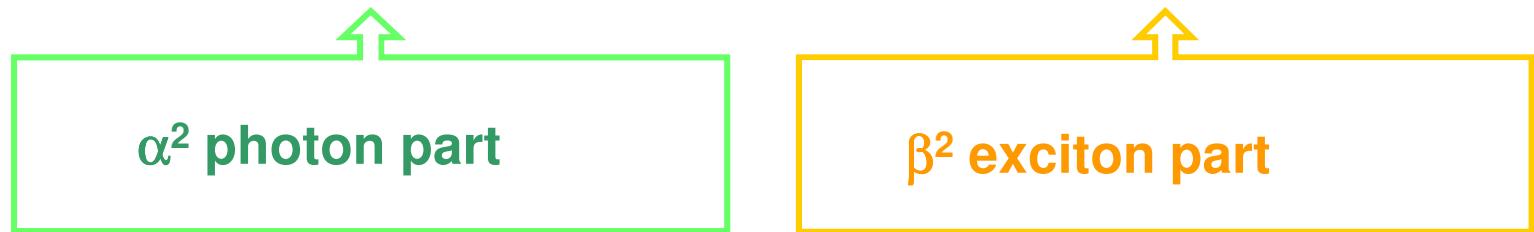


Cavity



Cavity polaritons : an exciton-photon mixed state

$$| \text{polariton} \rangle = \alpha | \text{photon} \rangle + \beta | \text{exciton} \rangle$$

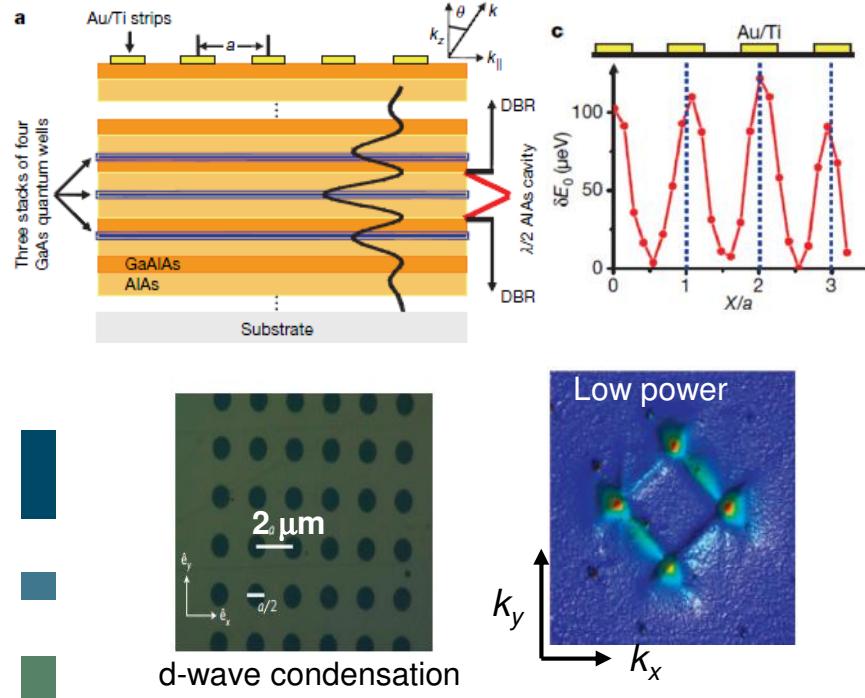


Properties

- Photonic component \rightarrow low mass ($10^{-5} m_e$)
- Short lifetime (1-100 ps) \rightarrow dissipative system
- Pseudo spin
- Sensitivity to magnetic field
- Polariton Interaction \rightarrow Highly non linear system (Kerr like)
- Polariton lattices

Polariton lattices

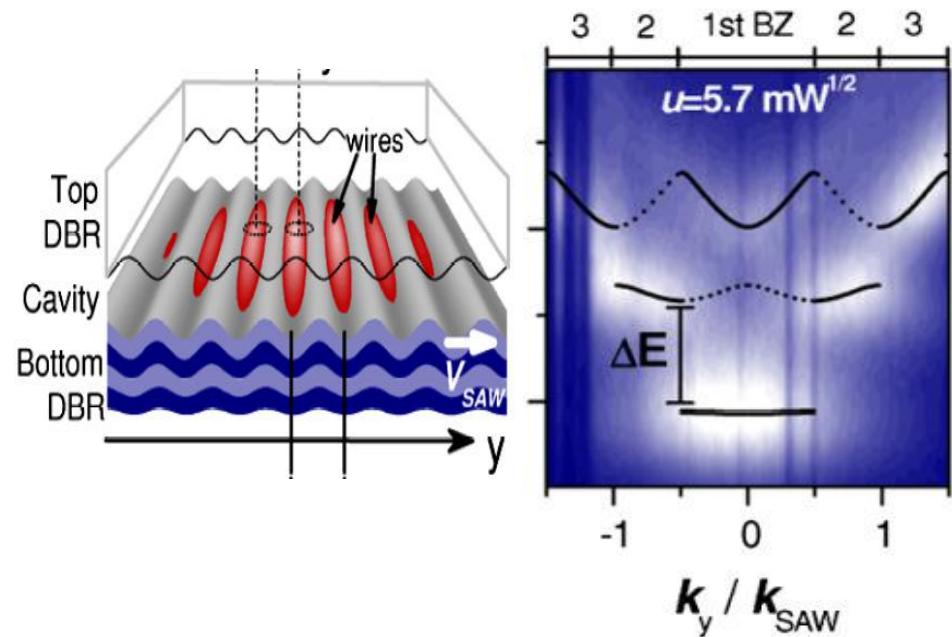
Gold deposition on top mirror



Lai et al., Nature 450, 529 (2007)
Kim et al., Nature Phys 7, 681 (2011)

Stanford

Surface acoustic waves



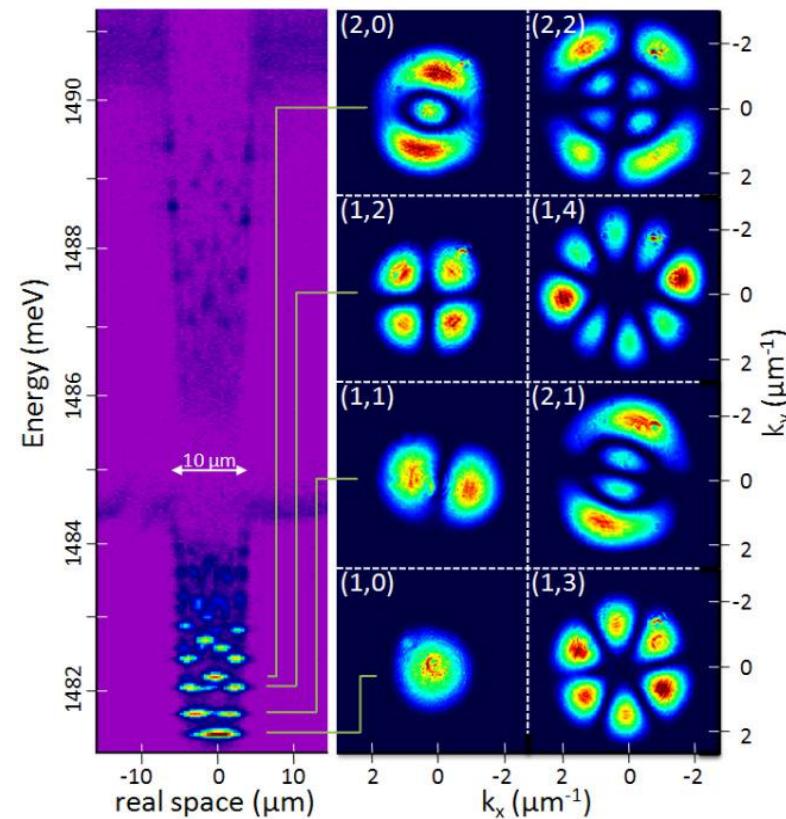
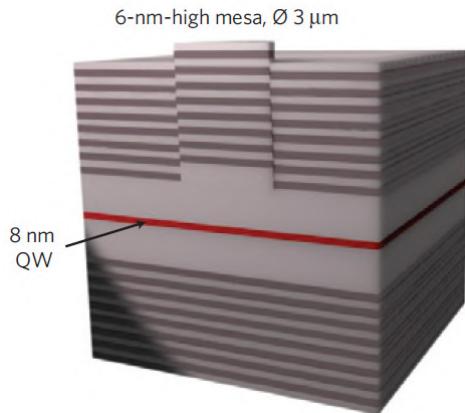
Cerdá-Méndez et al., PRL 105, 116402 (2010)
de Lima et al., PRL 97, 045501 (2006)

Cerdá-Méndez et al. PRB 86, 100301(R) (2012)

Berlin

During-growth photonic trap

Deveaud's group at EPFL



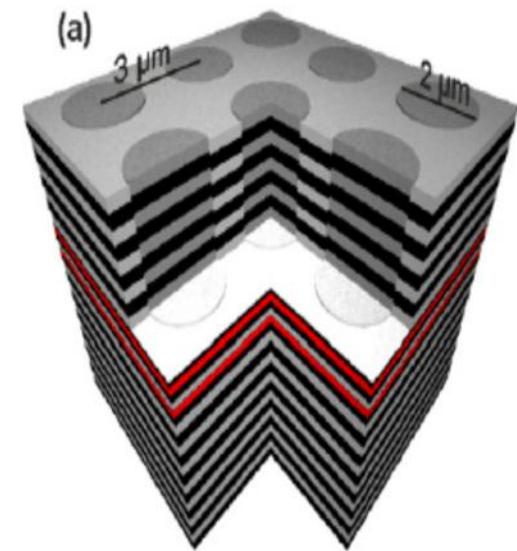
El Daïf et al., APL **88**, 061105 (2006)

Idrissi Kaitouni et al., PRB **74**, 155311 (2006)

Cerna et al., PRB **80**, 121309 (2009)

Würzburg

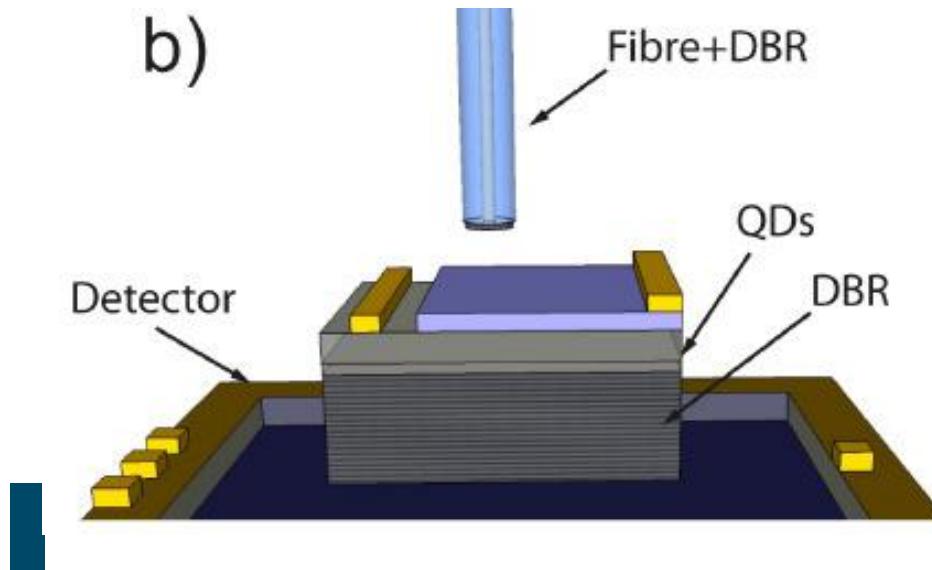
Coupled mesas



Winkler et al., New J. Phys. **17**, 23001 (2015)

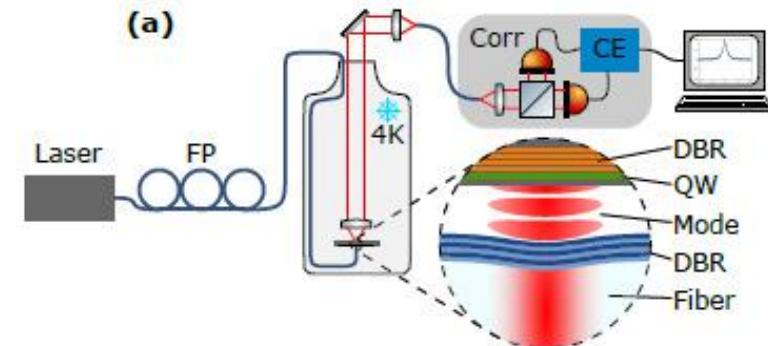
Hybrid cavities

b)



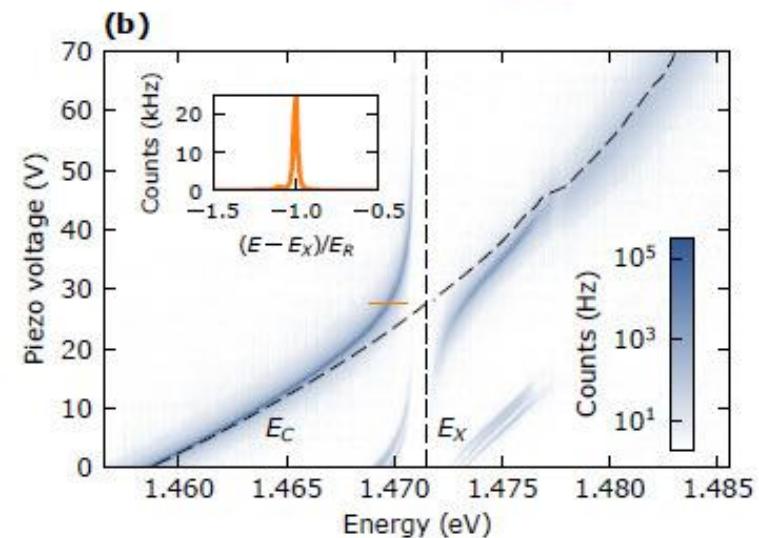
Fiber-closed cavity

Imamoglu, Reichel,



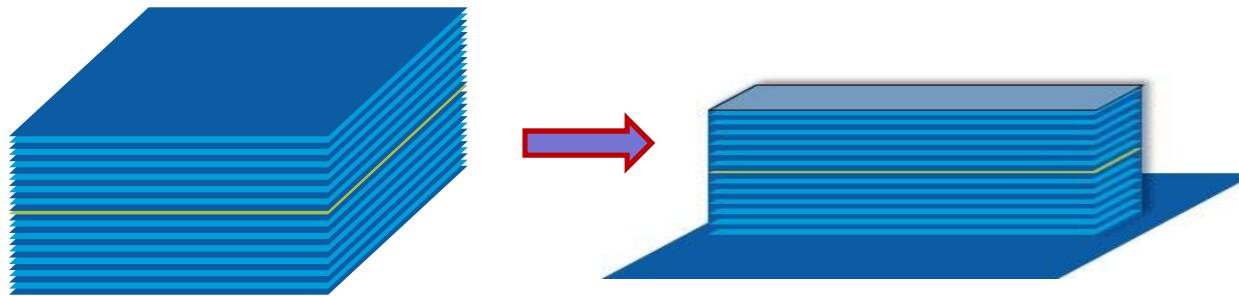
Besga et al., Light: S&A 3, e135 (2014)
S. Dufferwiel et al. Appl. Phys. Lett. 104, 192107 (2014)

See also T. Fink et al., Nature Physics 14, 365 (2018)
G. Muñoz-Matutano et al, Nature Materials 18, 213 (2019)
A. Delteil et al., Nature Materials 18, 219 (2019)

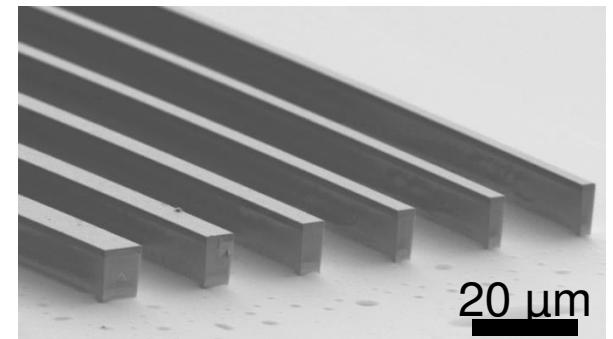


Polariton in 1D cavities

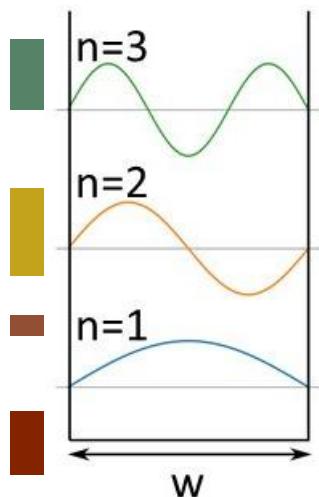
Electron beam lithography
+ ion dry etching



Polariton lateral confinement

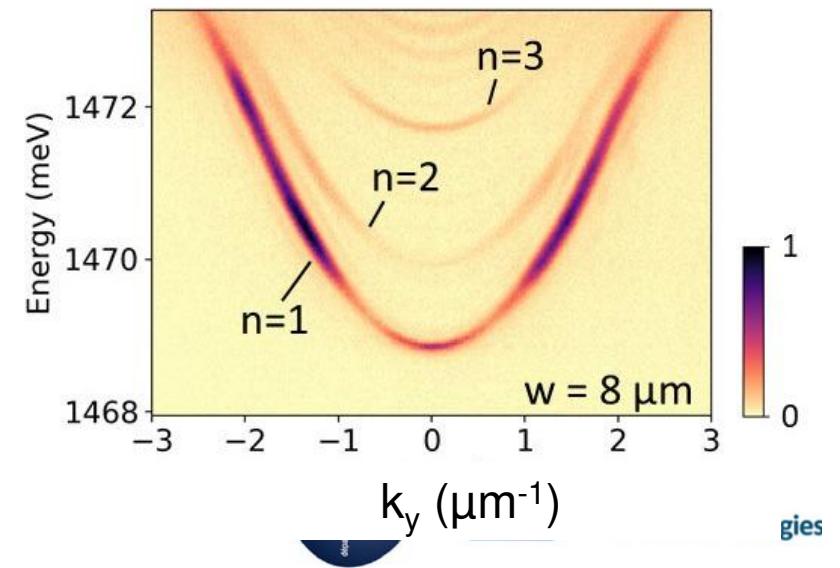


Quantization of k_x : $k_x = n\pi/w$
 $n=1,2,3,\dots$

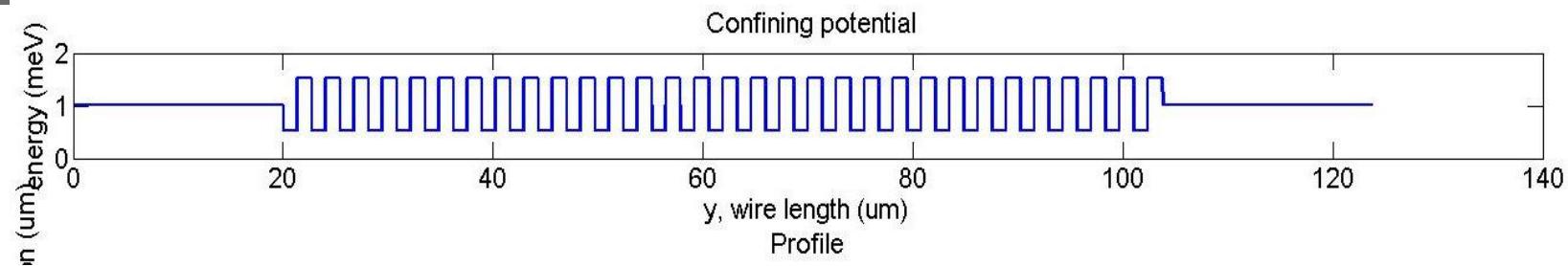
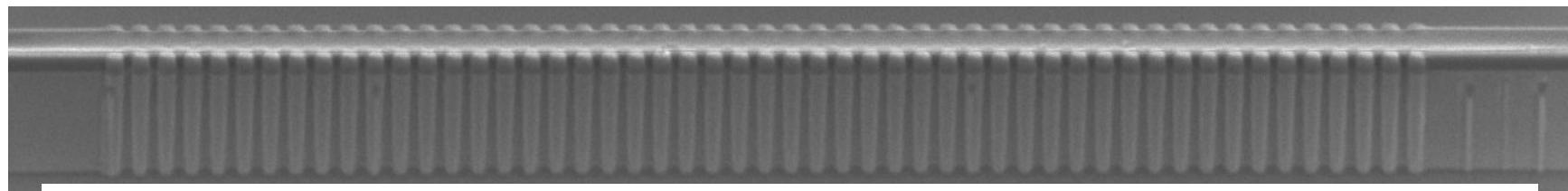


$$E_n(k_y) = E_0 + \frac{\hbar^2}{2m} \left(\frac{n^2\pi^2}{w^2} + k_y^2 \right)$$

1D polariton sub-bands

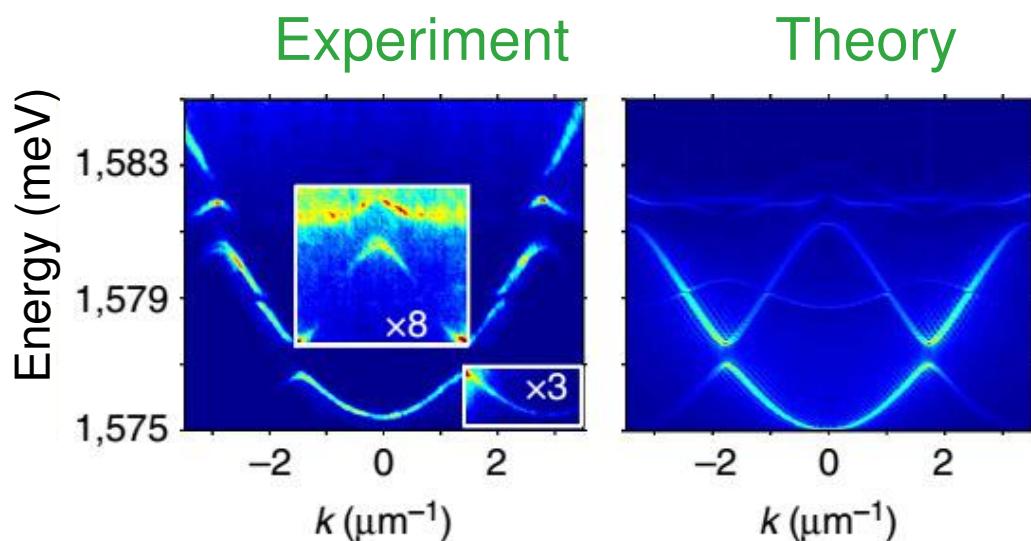


1D periodic potential: free particle approach



$$E\psi(x) = -\frac{\hbar^2}{2m} \frac{\partial\psi(x)}{\partial x} + V(x)\psi(x)$$

$$V(x) = \frac{\hbar^2}{2m} \frac{n^2 \pi^2}{w(x)^2}$$



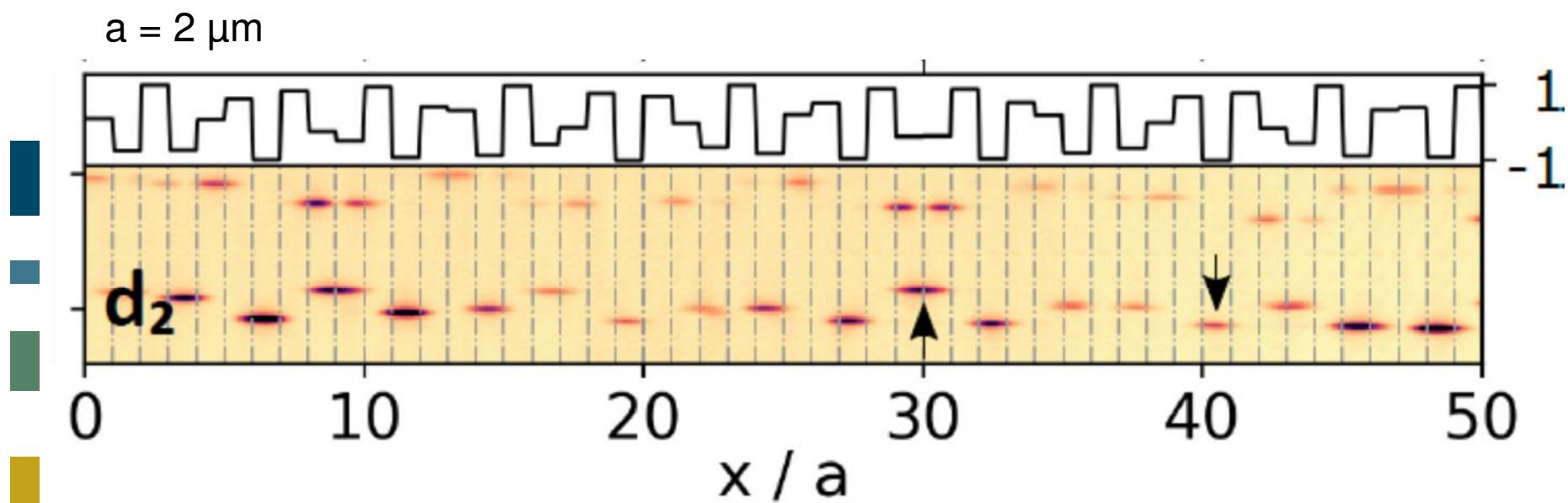
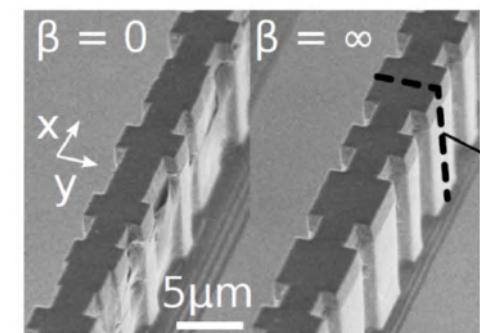
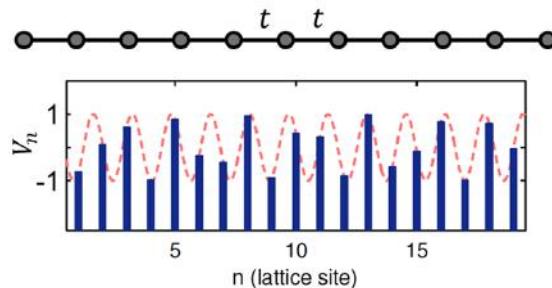
D. Tanese et al., Nature Com. 4, 1749 (2013)



Centre
de Nanosciences
et de Nanotechnologies

Aubry André (Harper) quasi-periodic potential

On-site potential incommensurate with the lattice period

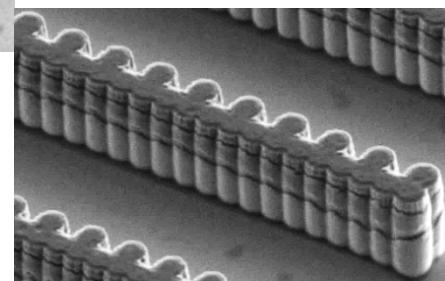
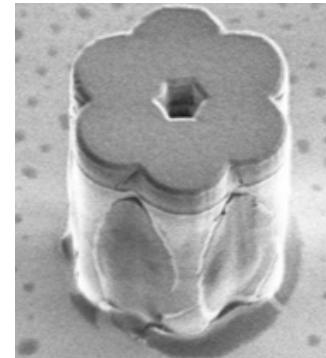
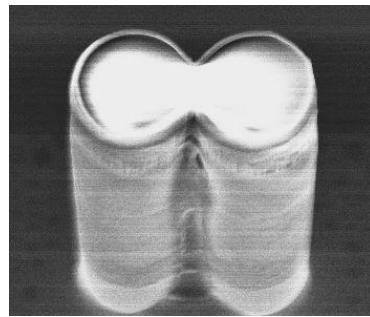
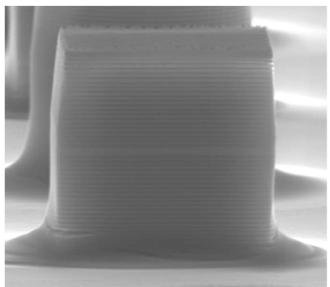


V. Goblot et al., Arxiv1911.07809

- Roati et al., Nature **453**, 895 (2008)
- Lahini et al., Phys. Rev. Lett. **103**, 013901 (2009)

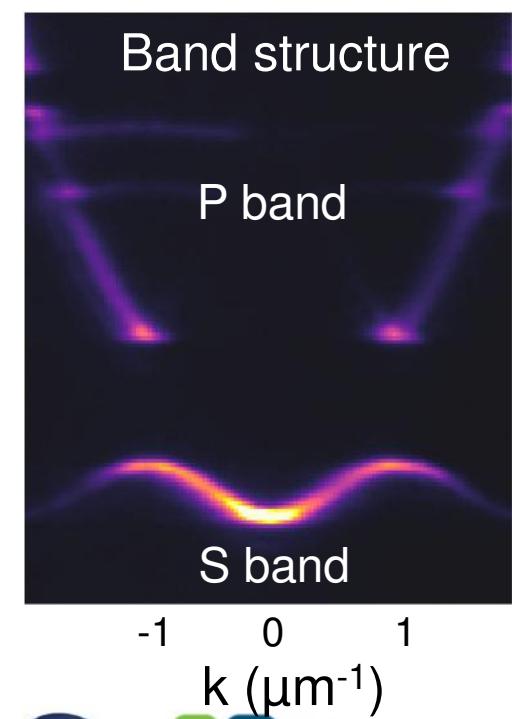
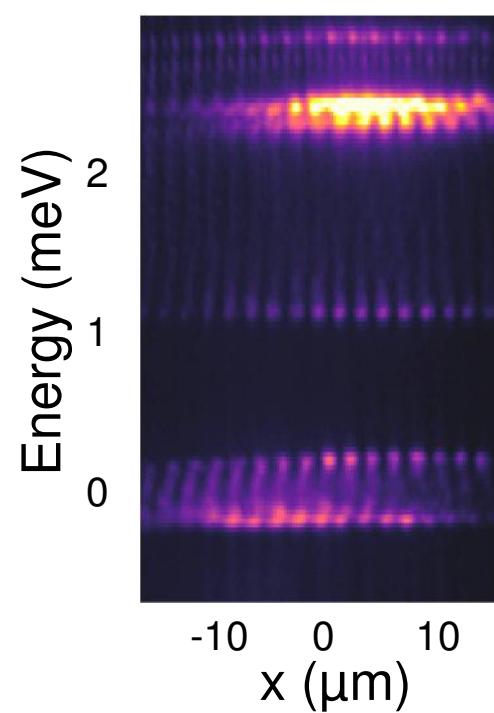
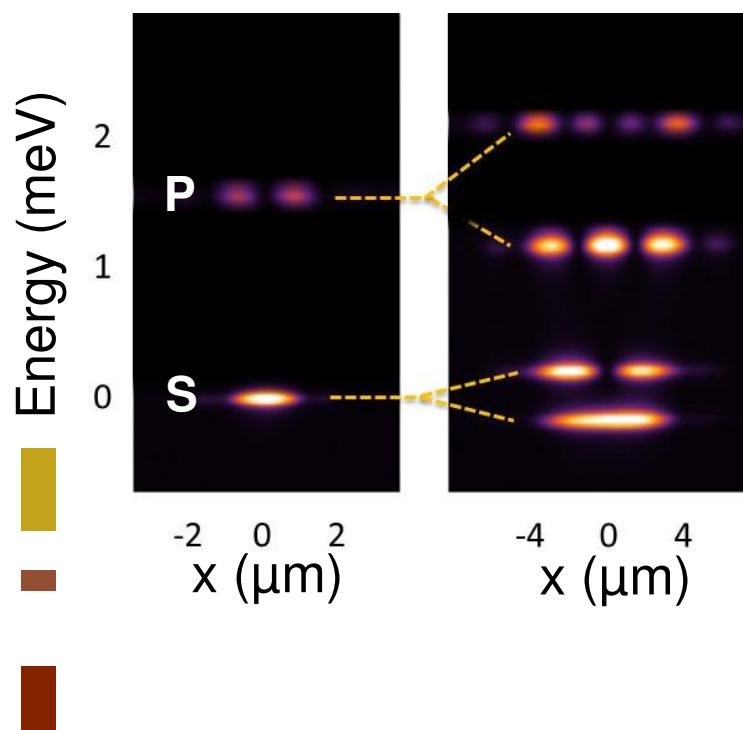
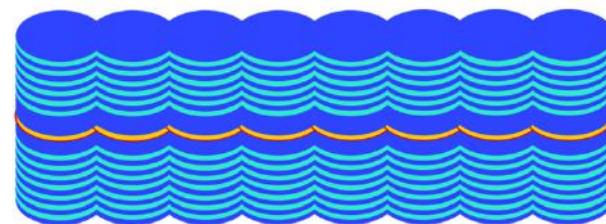
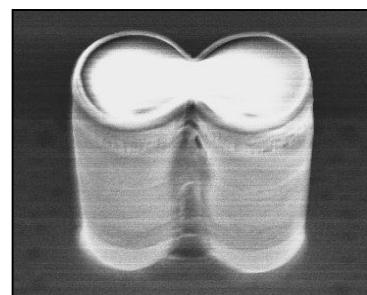
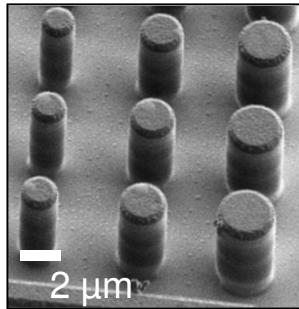


Polariton lattices:Tight binding approach

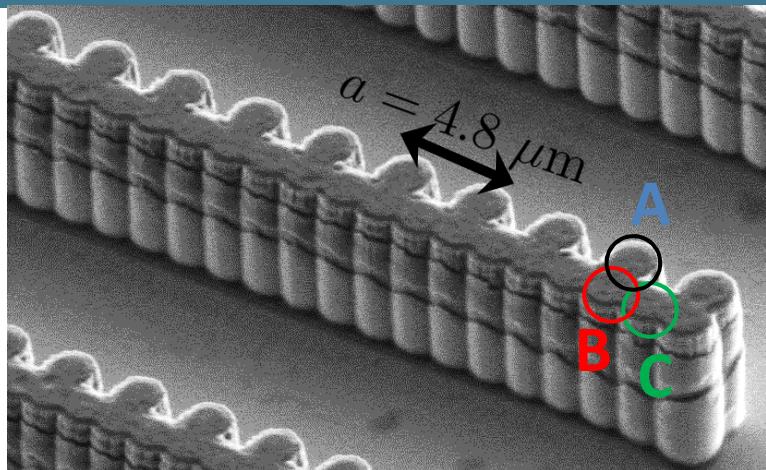


Review: C. Schneider et al., Rep. Prog. Phys. 80, 16503 (2017)

Polariton lattices:Tight binding approach



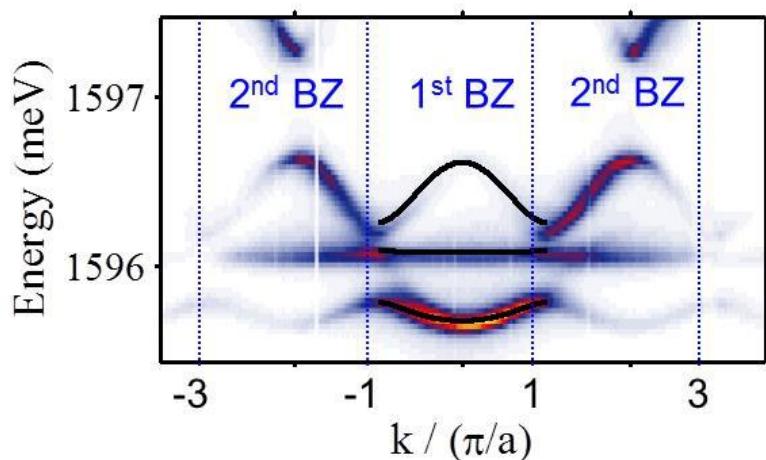
Engineering of a 1D flatband : “comb” lattice



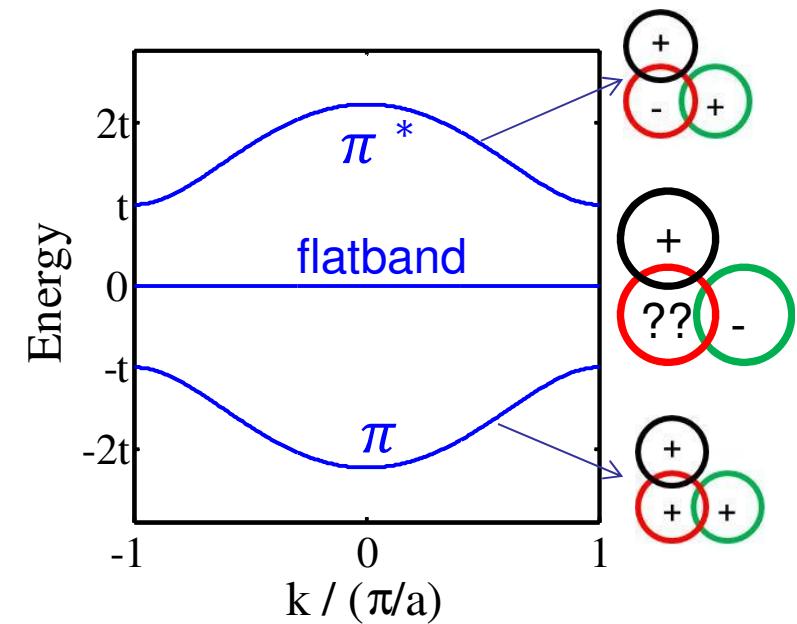
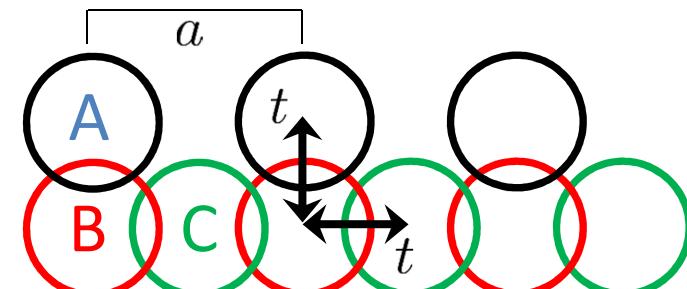
Pillar diameter = 3 μm
Interpillar distance = 2,4 μm



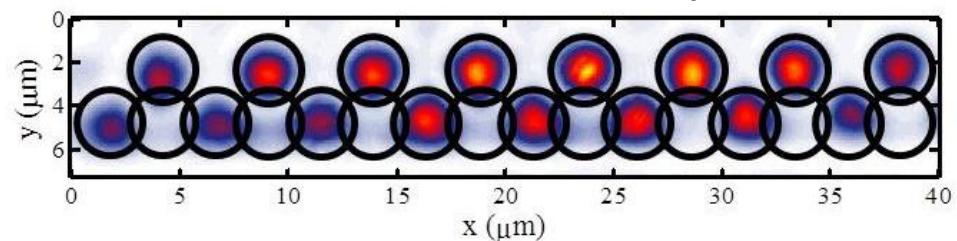
Far field emission



F. Baboux et al. PRL116, 066402 (2016)

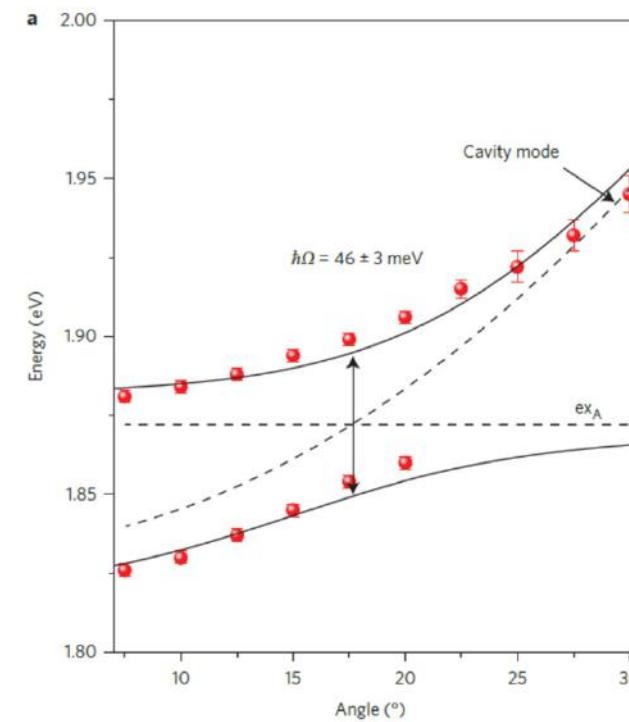
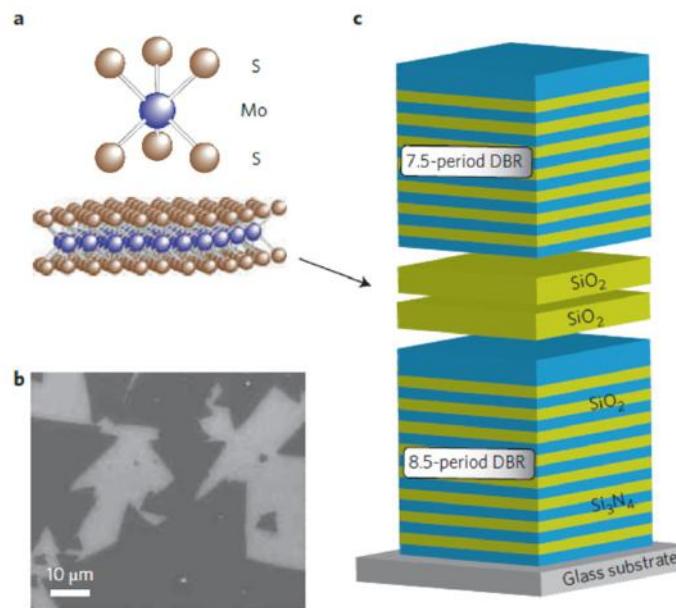


Emission of flat band in real space



Strong light-matter coupling in two-dimensional atomic crystals

Xiaoze Liu^{1,2}, Tal Galfsky^{1,2}, Zheng Sun^{1,2}, Fengnian Xia³, Erh-chen Lin⁴, Yi-Hsien Lee⁴, Stéphane Kéna-Cohen⁵ and Vinod M. Menon^{1,2*}





Ground-State Chemical Reactivity under Vibrational Coupling to the Vacuum Electromagnetic Field

Anoop Thomas[†], Jino George[†], Atef Shalabney, Marian Dryzhakov, Sreejith J. Varma, Joseph Moran, Thibault Chervy, Xiaolan Zhong, Eloïse Devaux, Cyriaque Genet, James A. Hutchison, and Thomas W. Ebbesen*

Thomas Ebbesen
2019 CNRS Gold medal

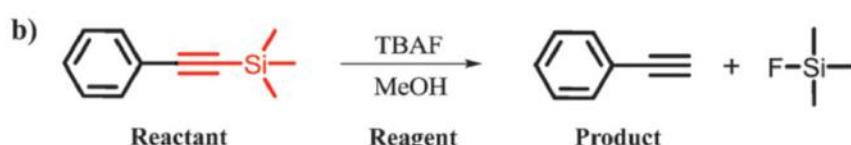
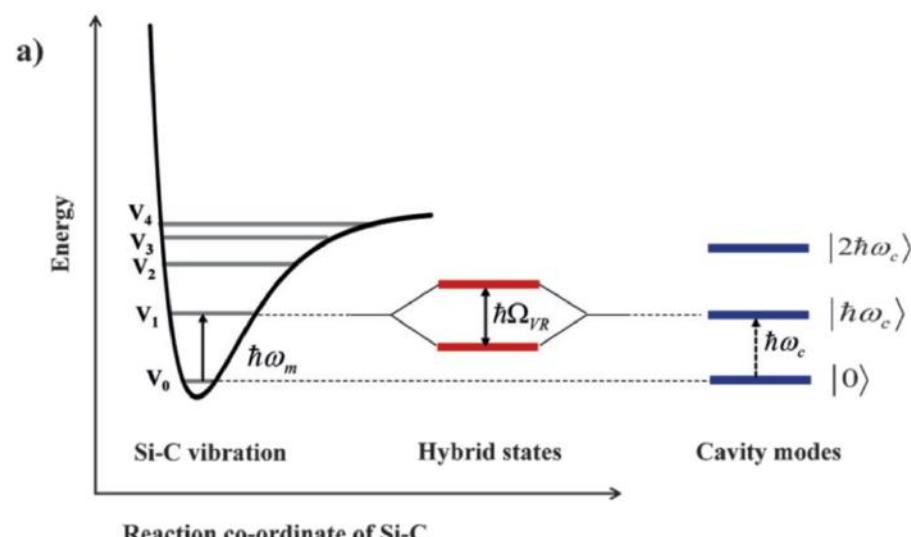


Figure 1. a) The light-matter strong coupling between the Si-C stretching vibrational transition and a cavity mode that results in the Rabi splitting. b) The silane deprotection reaction of 1-phenyl-2-trimethylsilylacetylene used in the present study.

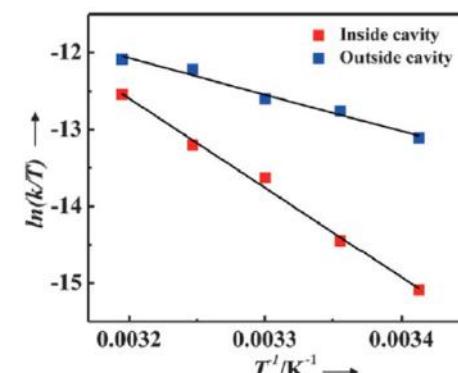
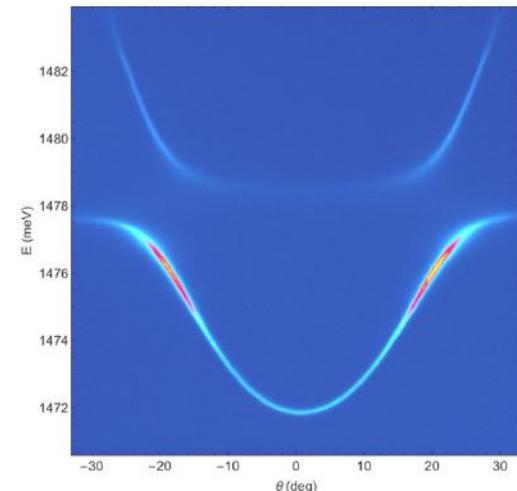


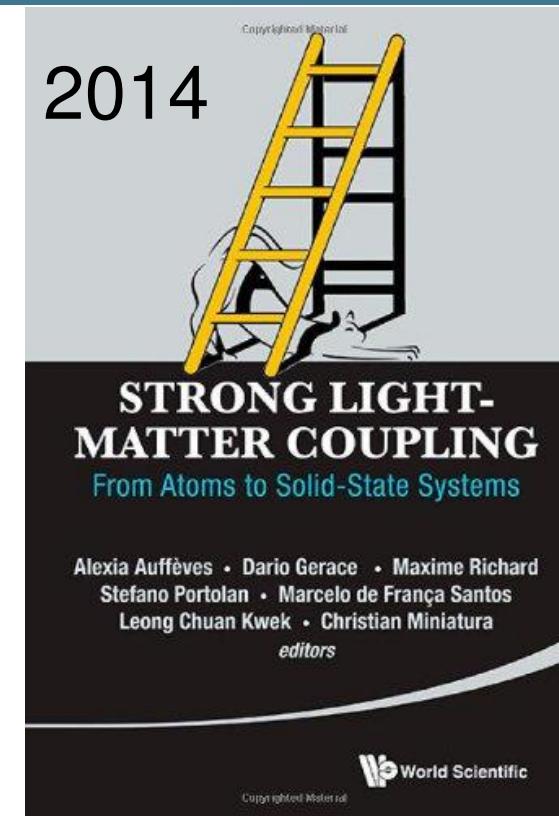
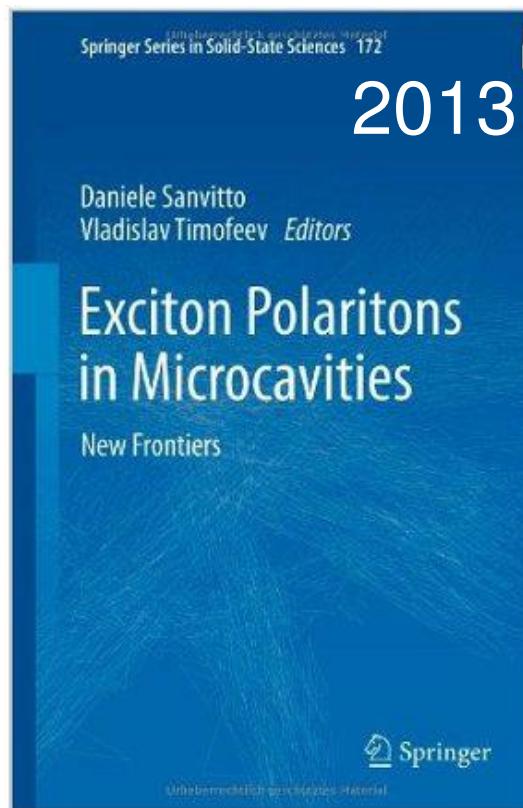
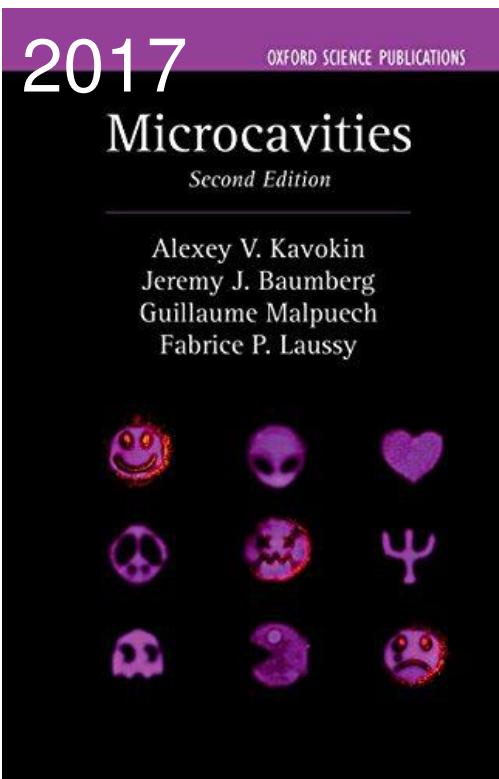
Figure 4. The reaction rate ($[PTA]=2.53\text{ M}$, $\hbar\Omega_{VR}=98\text{ cm}^{-1}$) as a function of temperature (Eyring equation plot) for reactions inside the ON resonance cavity (red squares) and outside the cavity (blue squares).

Summary

- Hybrid exciton-photon quasi-particles
- Tunable properties : effective mass, group velocity
- Lateral confinement : engineering of band structure
- ➤ Spin dependent interactions
- ➤ Optical access to all physical properties



Reviews



C. Ciuti & I. Carusotto, *Quantum fluids of light*: Rev. Mod. Phys. **85**, 299 (2013)

A. Amo and J. Bloch, “*Exciton-polaritons in lattices: A non-linear photonic simulator*”, Comptes Rendus de l’Académie des Sciences 8, 805 (2016)