



INSTITUT  
POLYTECHNIQUE  
DE PARIS



# *Quantum Many-Body Physics*

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<https://www.cpht.polytechnique.fr/cpht/uquantmat/>

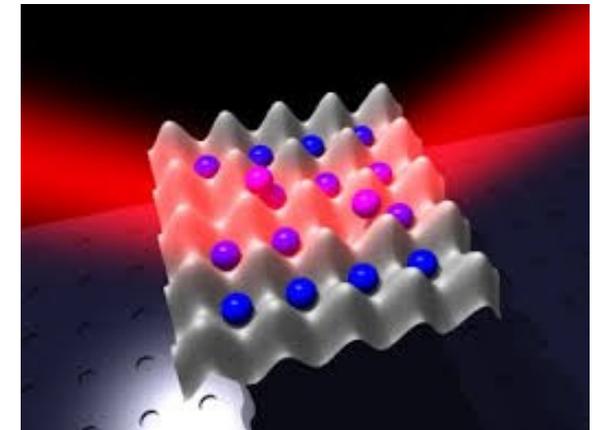


## Research field

Theory of quantum many-body systems

Quantum science and technologies ; Quantum simulation with ultracold atoms

Understand the behavior of correlated quantum matter and how to simulate them with controlled quantum systems



# Quantum Many-Body Physics

## 1. Introduction to many-body quantum physics

Few-body versus many-body problems

Quantum versus classical many-body systems

Quantum phases and quantum phase transitions

## 2. Models and methods

Second quantization approach

How to build a many-body model (a heuristic approach)

Case study of the Bose-Hubbard model

## 3. Strongly-correlated bosons in shallow quasiperiodic potentials

Physics of disorder

Bose-glass physics in shallow quasiperiodic potentials

Quantum quasicrystals : The Bose side

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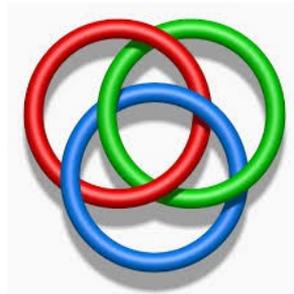
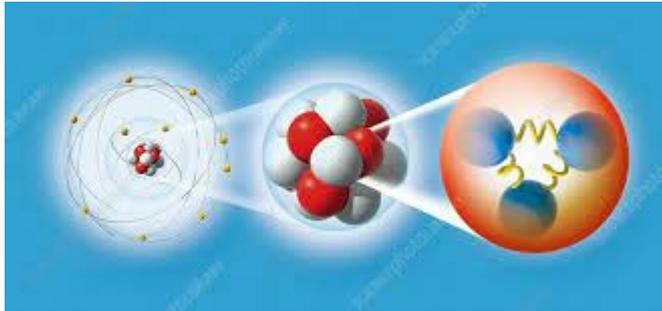
Quantum quasicrystals : The Bose side

# From Few-Body Physics ...

Understanding the physics of systems made of constituents that interact with each other

Many applications at all scales

A difficult problem already for 3 particles (eg dynamical instabilities, chaos, ...)



# ... to Many-Body Physics

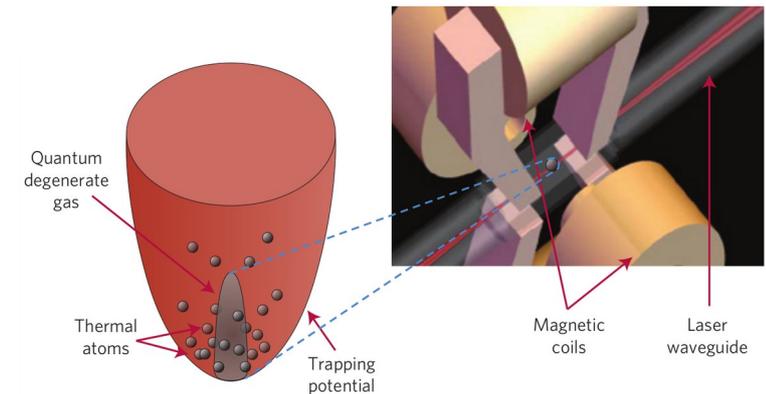
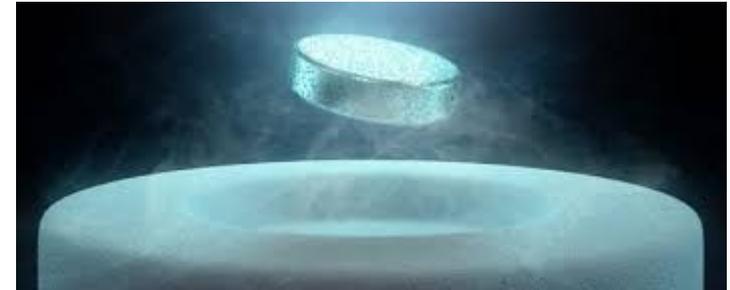
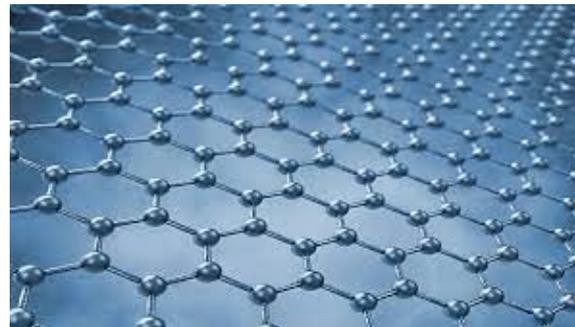
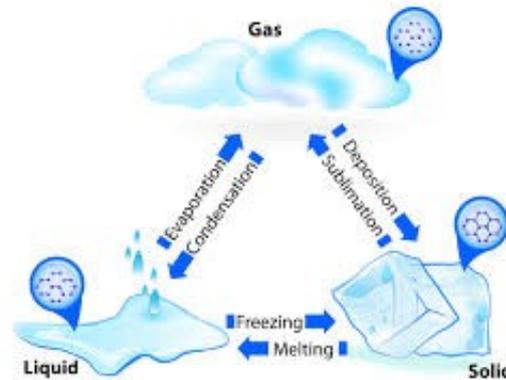
Systems made of huge number of constituents, which all interact with each other

$10^{4-6}$  (ultracold atoms),  $10^{23}$  (Avogadro),  $10^{27}$  (human body),  $10^{50}$  (Earth),  $10^{80}$  (Universe)

Chaos and statistical physics  
(benefits of complexity)

Linear response theory

Linearity of quantum physics  
and the problem of  
thermalization (ergodicity,  
eigenstate thermalization  
hypothesis, and many-body  
localization)



# Quantum Many-Body Systems

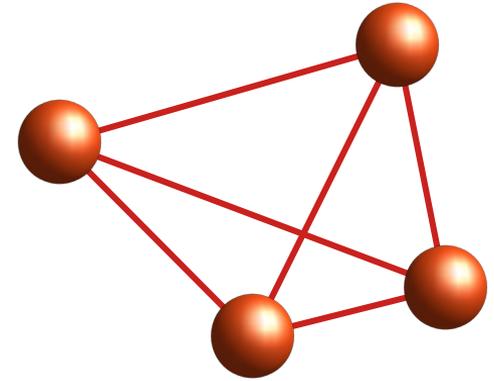
In spite of linearity, major issues arise

Quantum coherence (interference), non-commuting observables

Few-body interactions but collective effects, including non-separability of constituents

Entanglement, huge Hilbert space

Quantum indiscernability (bosons, fermions, symmetrization and anti-symmetrization) and spins



Theoretical physics has developed a number of tools, yet

Many-body physics remains a major issue in condensed matter physics and beyond

Each system requires specific attention and educated approximations

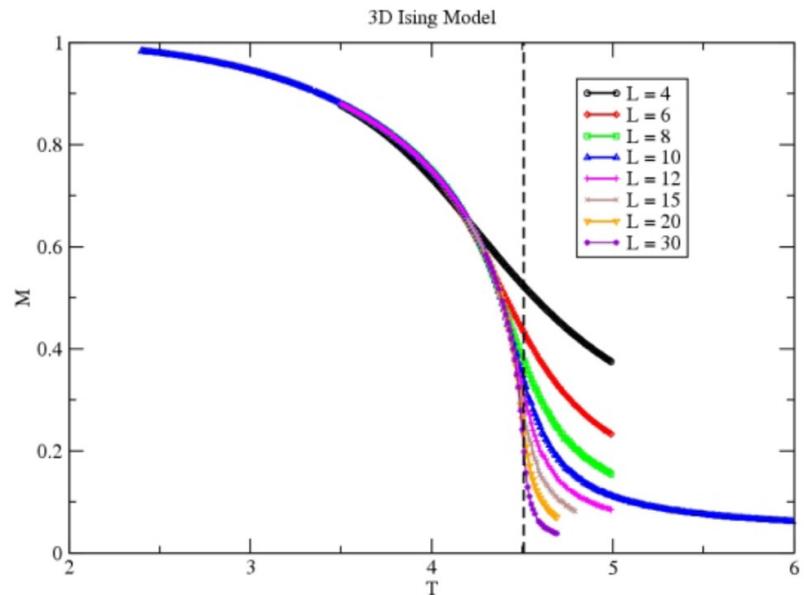
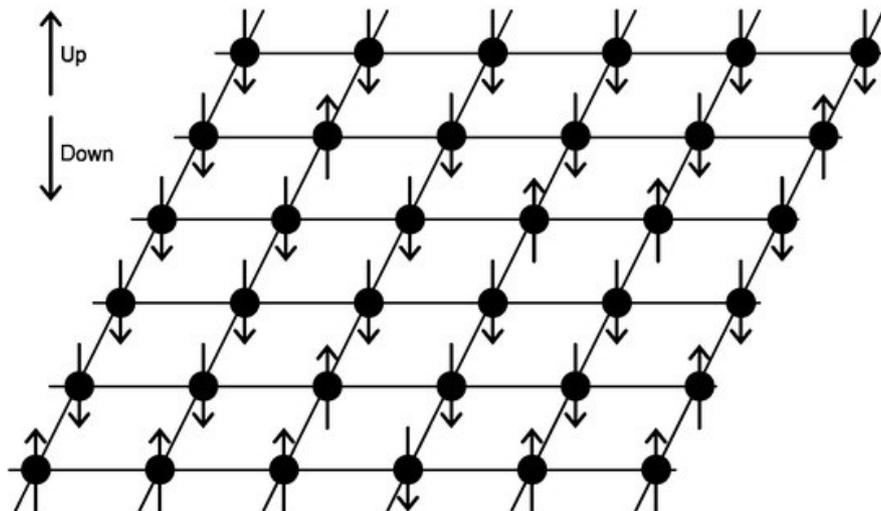
# Phase Transitions and Emerging Phenomena

## Phase transitions in classical many-body systems

Distinctive characteristics of a many-body system (« More is different », P. Anderson, 1972)

Solid-liquid-gas transitions ; Magnetism, etc ...

Classical phase transitions driven by thermal fluctuations



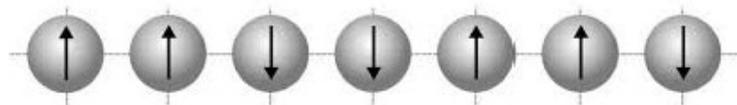
# Phase Transitions and Emerging Phenomena

## Phase transitions in quantum many-body systems

Phase transitions driven by thermal fluctuations (*classical phase transitions* ; superconductivity, superfluidity, magnetic transitions, ...)

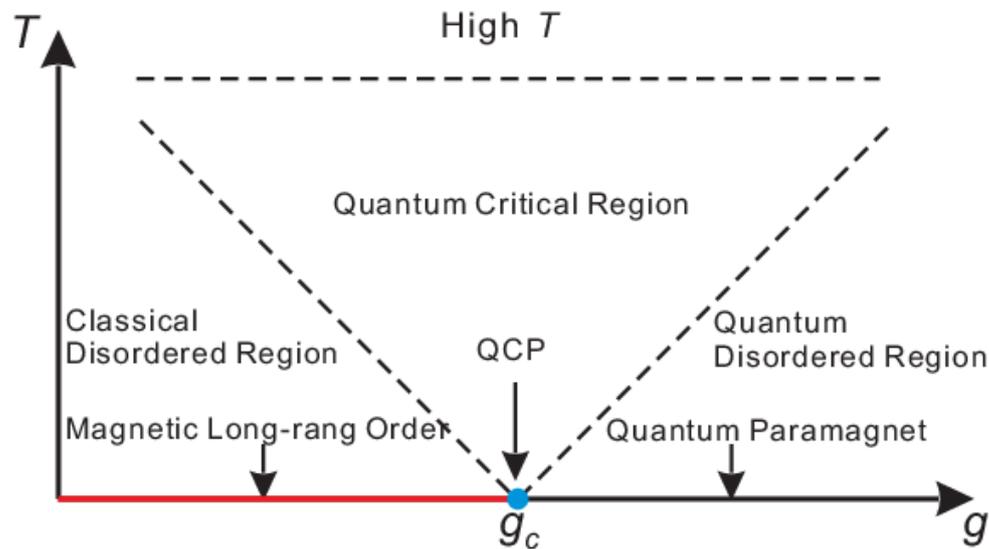
Phase transitions driven by quantum fluctuations (*quantum phase transitions* ; metal-insulator transitions, quantum magnetic transitions, ...)

### Transverse-field Ising chain



$$\hat{H} = -J \sum_{\langle j, \ell \rangle} \hat{\sigma}_j^x \cdot \hat{\sigma}_\ell^x + h \sum_j \hat{\sigma}_j^z$$

$$g = h/J$$



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# Second Quantization Formalism

A new formalism adapted to indistinguishable quantum particles

**Bosons** ( $\eta = +1; n_\lambda \in \mathbb{N}$ )

$$\hat{a}_\lambda |n_1 \dots n_\lambda \dots\rangle \equiv \sqrt{n_\lambda} |n_1 \dots n_\lambda - 1 \dots\rangle$$

$$\hat{a}_\lambda^\dagger |n_1 \dots n_\lambda \dots\rangle \equiv \sqrt{n_\lambda + 1} |n_1 \dots n_\lambda + 1 \dots\rangle$$

$$\hat{n}_\lambda = \hat{a}_\lambda^\dagger \hat{a}_\lambda$$

$$[\hat{a}_\lambda, \hat{a}_{\lambda'}] = [\hat{a}_\lambda^\dagger, \hat{a}_{\lambda'}^\dagger] = 0$$

$$[\hat{a}_\lambda, \hat{a}_{\lambda'}^\dagger] = \delta_{\lambda, \lambda'}$$

**Fermions** ( $\eta = -1; n_\lambda \in \{0, 1\}$ )

$$\hat{a}_\lambda |n_1 \dots n_\lambda \dots\rangle \equiv \prod_{\lambda' < \lambda} (-1)^{n_{\lambda'}} |n_1 \dots n_\lambda - 1 \dots\rangle$$

$$\hat{a}_\lambda^\dagger |n_1 \dots n_\lambda \dots\rangle \equiv \prod_{\lambda' < \lambda} (-1)^{n_{\lambda'}} |n_1 \dots n_\lambda + 1 \dots\rangle$$

$$\hat{n}_\lambda = \hat{a}_\lambda^\dagger \hat{a}_\lambda$$

$$\{\hat{a}_\lambda, \hat{a}_{\lambda'}\} = \{\hat{a}_\lambda^\dagger, \hat{a}_{\lambda'}^\dagger\} = 0$$

$$\{\hat{a}_\lambda, \hat{a}_{\lambda'}^\dagger\} = \delta_{\lambda, \lambda'}$$

One- and two-body observables

First quantization  $\longrightarrow$  Second quantization

$$\hat{A} = \sum_{1 \leq j \leq N} \hat{A}^{(j)}$$

$$\hat{A} = \sum_{\lambda, \lambda'} \langle \lambda' | \hat{A} | \lambda \rangle \hat{a}_{\lambda'}^\dagger \hat{a}_\lambda$$

$$\hat{B} = \frac{1}{2} \sum_{j \neq \ell} \hat{B}^{(j, \ell)}$$

$$\hat{B} = \frac{1}{2} \sum_{\lambda_1, \lambda_2, \lambda'_1, \lambda'_2} \langle \lambda'_1, \lambda'_2 | \hat{B} | \lambda_1, \lambda_2 \rangle \hat{a}_{\lambda'_1}^\dagger \hat{a}_{\lambda'_2}^\dagger \hat{a}_{\lambda_2} \hat{a}_{\lambda_1}$$

H. Bruus & K. Flensberg, *Many-body Quantum Theory in Condensed Matter Physics* (Oxford Univ Press, UK, 2004)

J. W. Negele & H. Orland, *Quantum many-particle systems* (Perseus books, USA, 1998)

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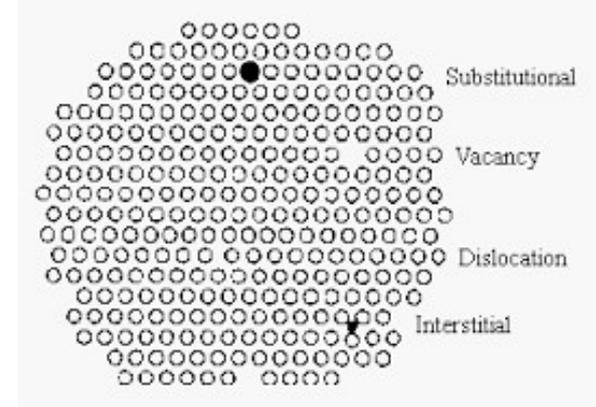
Quantum quasicrystals : The Bose side

# Physics of Disorder : From Condensed Matter to Ultracold Atoms

## Disorder is ubiquitous in condensed matter

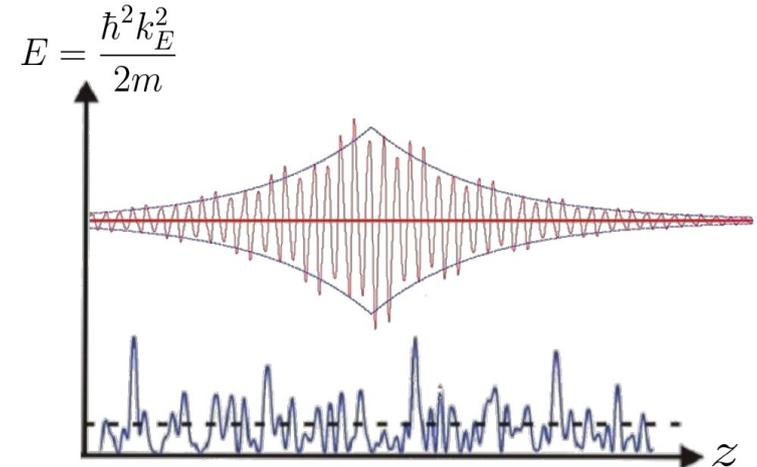
Almost any real-life material is subjected to defects

Disorder may strongly affects the physics of the material



→ Localization of single-particle wavefunctions  
(Anderson, 1958)

→ Wavefunctions are always localized in 1D and 2D ; Localization transition in 3D  
(Anderson *et al.*, 1979)

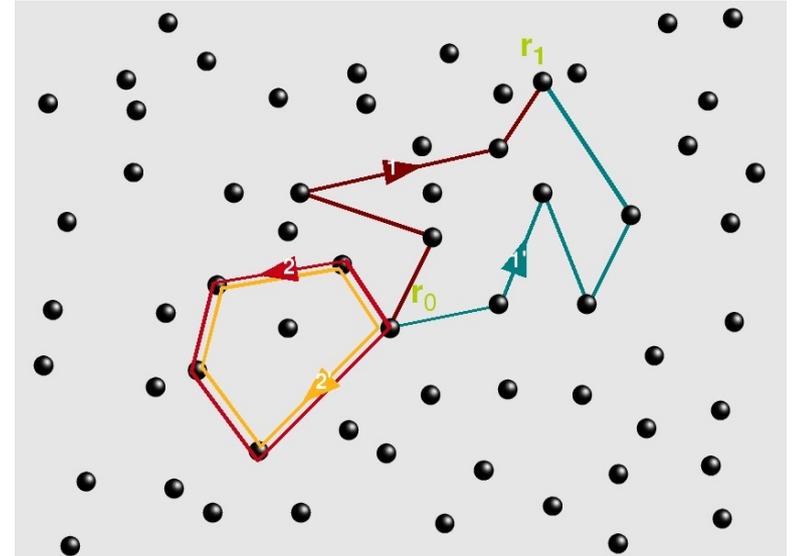


# Physics of Disorder : From Condensed Matter to Ultracold Atoms

## Wave physics and disorder

May be interpreted as an interference effect of waves scattering from the impurities

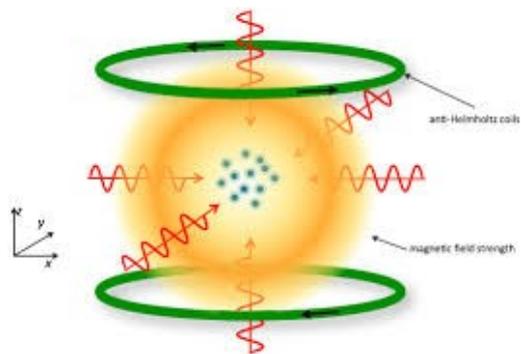
- Single-particle effect : Difficult to observe in traditional condensed matter systems
- May be observed with optical waves?
- Many attempts all refuted up to mid-2000's



## What about correlated quantum matter?

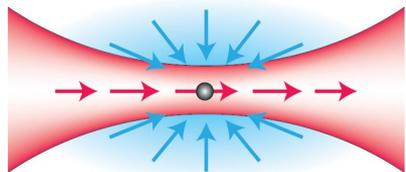
Very important (and difficult) problem in condensed matter (interplay of coherence, interactions, and disorder)

# Ultracold Atoms



Laser cooling ( $\sim 40\text{-}100\mu\text{K}$ )

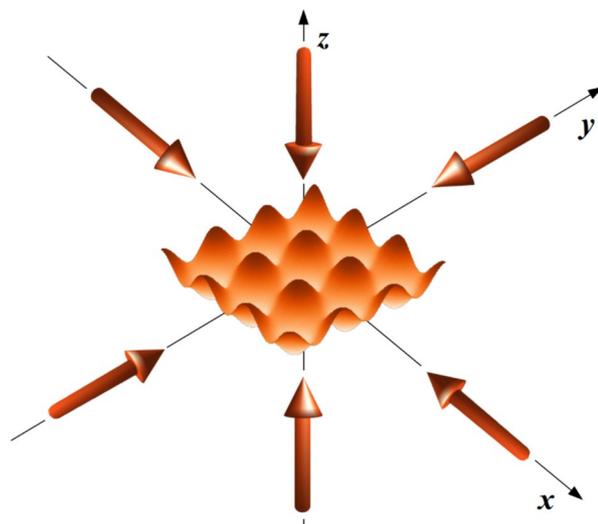
Optical trapping



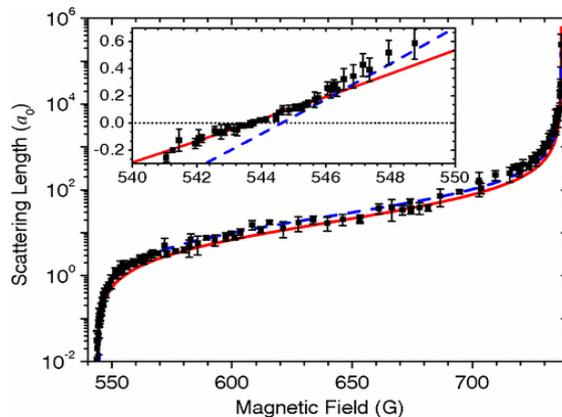
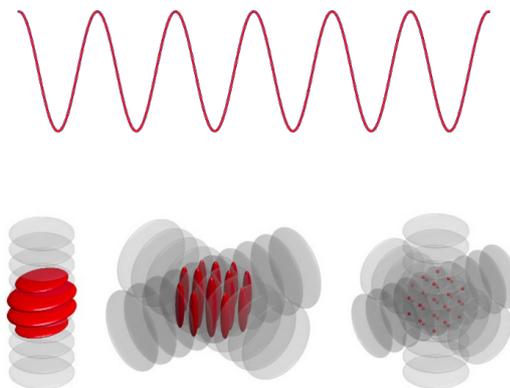
Evaporative cooling ( $< 100\text{nK}$ )

$$\rightarrow \lambda_{\text{dB}} \gtrsim 500 \text{ nm} \sim \lambda_{\text{opt}}$$

$$\rightarrow k_{\text{B}} T \lesssim E_{\text{r}}$$



Optical design



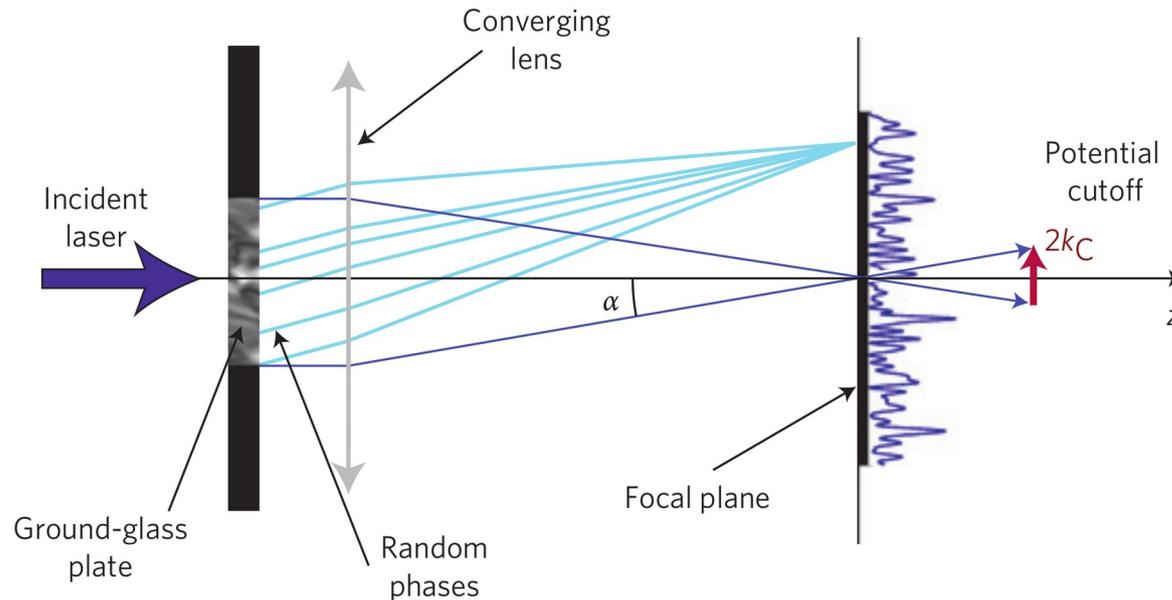
Control of interactions

$\rightarrow$  A wonderful playground  
for basic science  
(quantum simulation)

# Disordered Quantum Gases Under Control

Early proposals :

- Grynberg (ENS Paris, 2000) : Classical particles in disorder
- Lewenstein (Hannover, 2003) : Strongly-correlated quantum gases in disorder

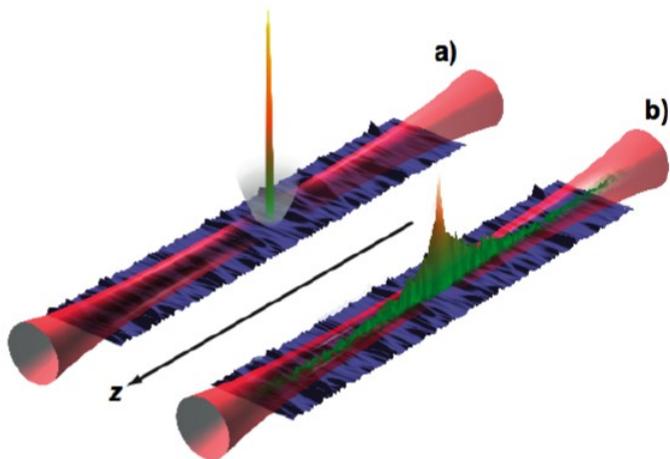


Controlled amplitude  
and correlation  
length/function

# Disordered Quantum Gases Under Control

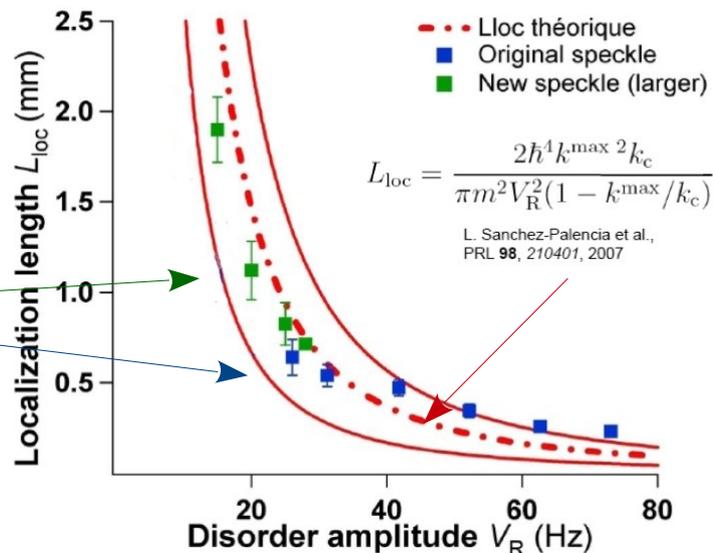
Exponential Anderson localization may be directly observed in expanding condensates and the localization length can be computed exactly !

Sanchez-Palencia *et al.*, Phys. Rev. Lett. **98**, 210401 (2007)



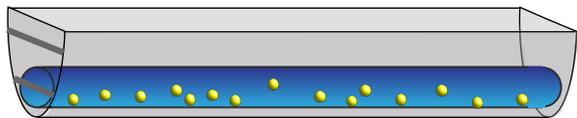
Aspect, Bouyer, Josse, Sanchez-Palencia  
[Nature **453**, 891 (2008)]

→ Agreement without any fitting parameter

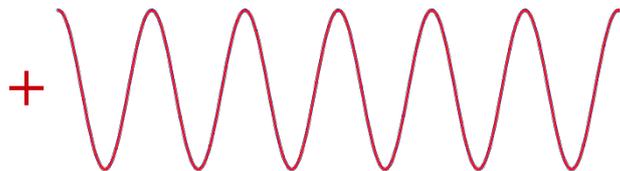


Progress paper : Sanchez-Palencia & Lewenstein, Nat. Phys. **6**, 87 (2010)

# Localization in the Aubry-André Model



→ 1D tubes created by a very strong 2D lattice



→ Tight-binding model created by a strong 1D lattice



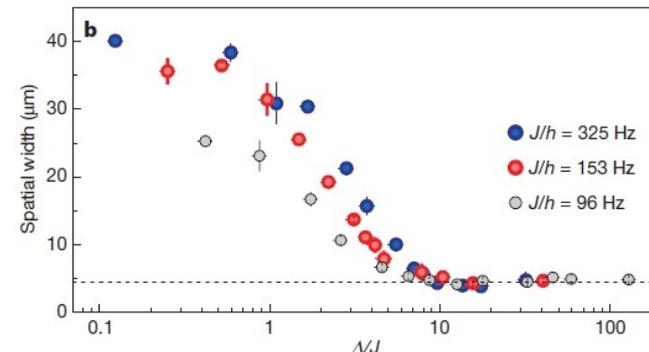
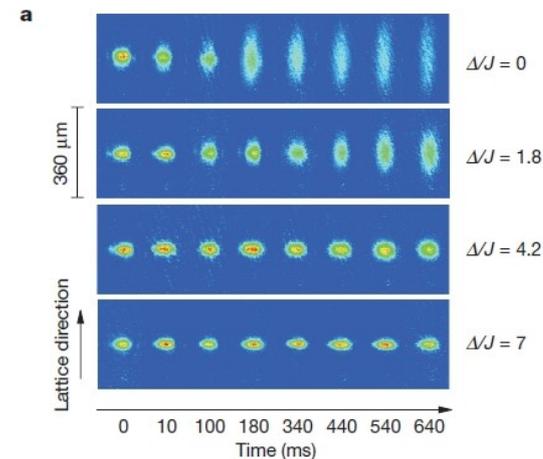
→ Quasi-periodic modulation created by a weak 1D lattice with incommensurate period

## Observation of localization in the Aubry-André model

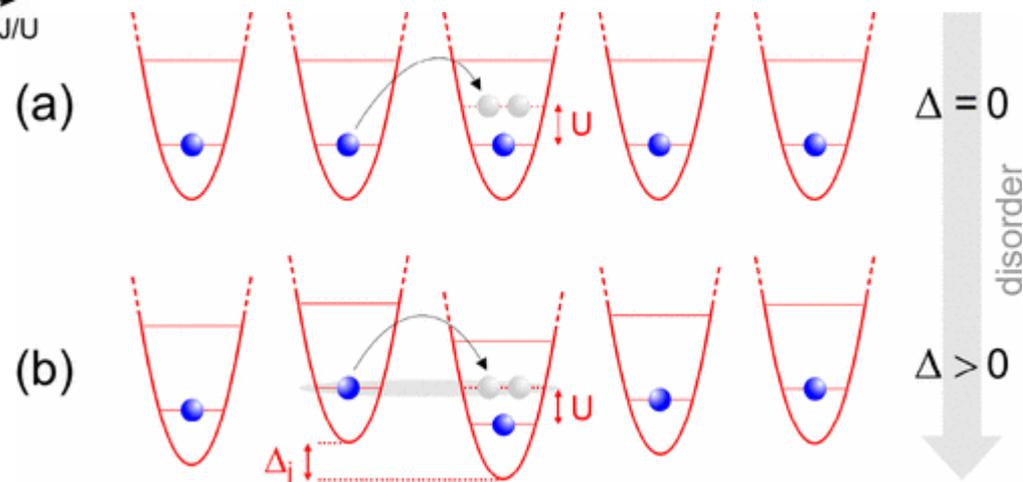
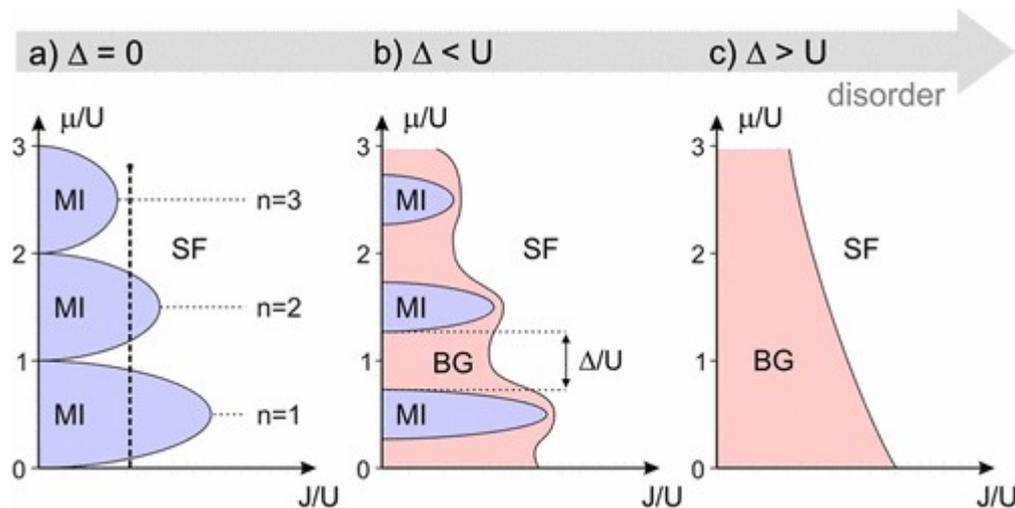
$$\hat{H} = -J \sum_j (\hat{a}_{j+1}^\dagger \hat{a}_j + \text{H.c.}) + \sum_j \Delta \cos(2\pi\beta j + \varphi) \hat{n}_j$$

with  $\beta \in \mathbb{Q}^c$

Inguscio, Modugno [Nature **453**, 895 (2008)]

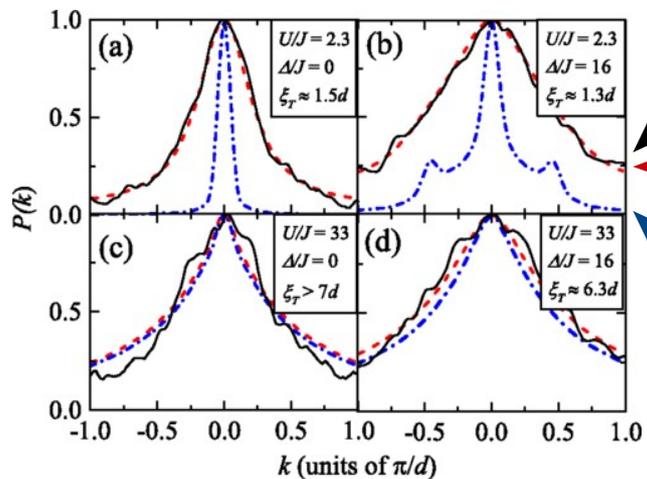
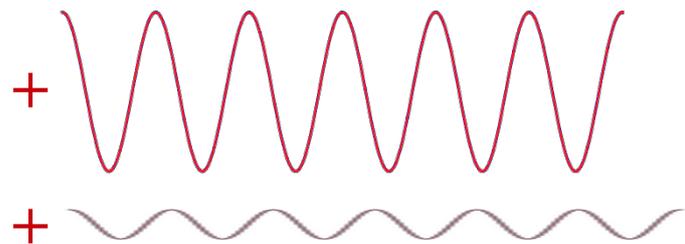
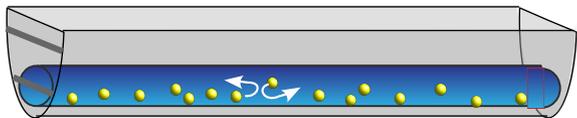


# Bose-Glass Physics in the Hubbard-Aubry-André Model



# Bose-Glass Physics in the Hubbard-Aubry-André Model

Inguscio, Modugno, Giamarchi, Roux [Phys. Rev. Lett. **113**, 095301 (2014)]

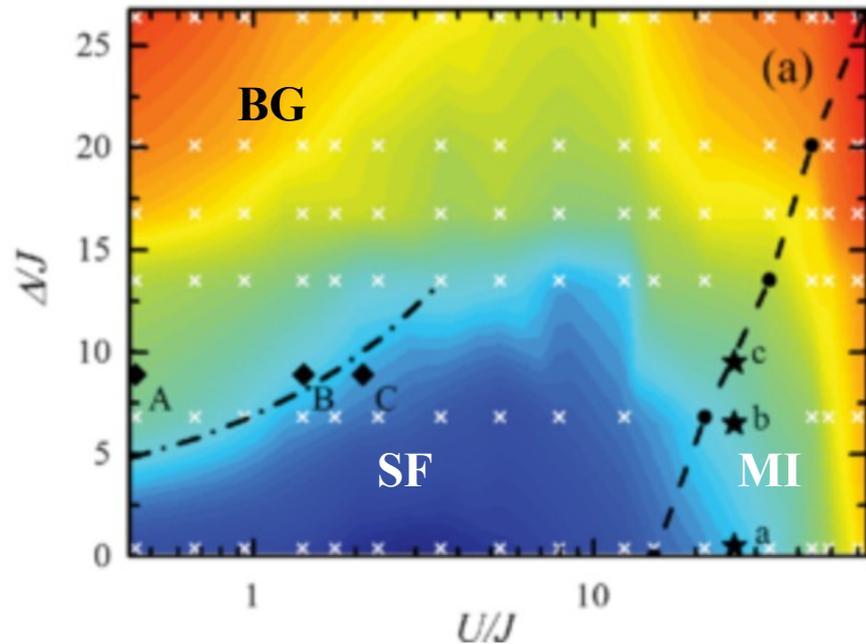


experiment

DMRG+heuristic  
thermal broadening

DMRG

Mobility and coherence measurements



→ Coherence blurred by thermal fluctuations

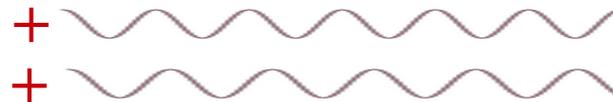
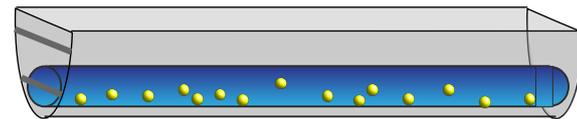
$$J, k_B T \ll E_r \quad \text{but} \quad J \sim k_B T$$

# Localization in a Shallow Quasiperiodic Lattice

Temperature is hard to further decrease

→ Scale up characteristic energies !

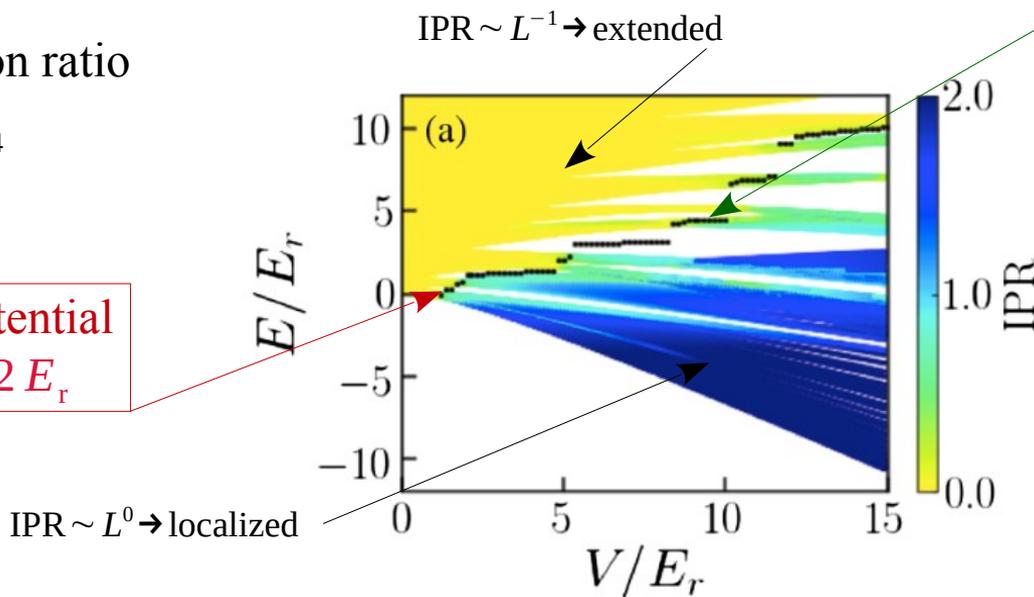
Do particles localize in a shallow quasiperiodic potential?



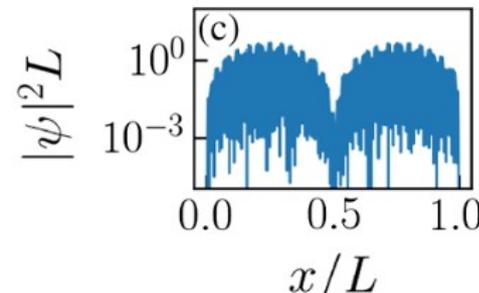
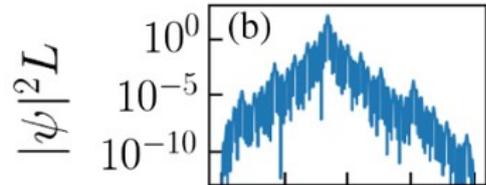
Inverse participation ratio

$$\text{IPR} = \int dx |\psi(x)|^4$$

critical potential  
 $V_c \approx 1.112 E_r$



mobility edge



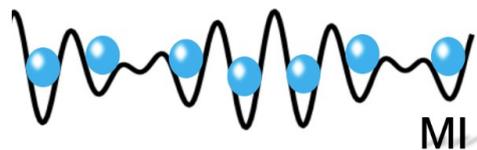
# Lieb-Liniger Gas in a Shallow Quasiperiodic Lattice

Let now add inter-particle interactions

$$\mathcal{H} = \sum_{1 \leq j \leq N} \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_j^2} + V(x_j) \right] + g \sum_{j < \ell} \delta(x_j - x_\ell)$$

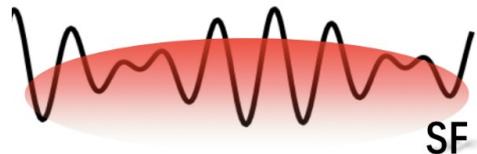


Quantum Monte Carlo in continuous space  
(compressibility and superfluid fraction)



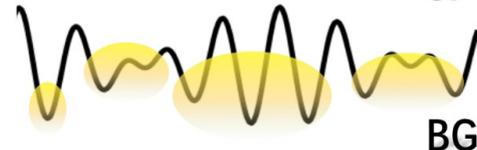
MI

Incompressible  
Non superfluid



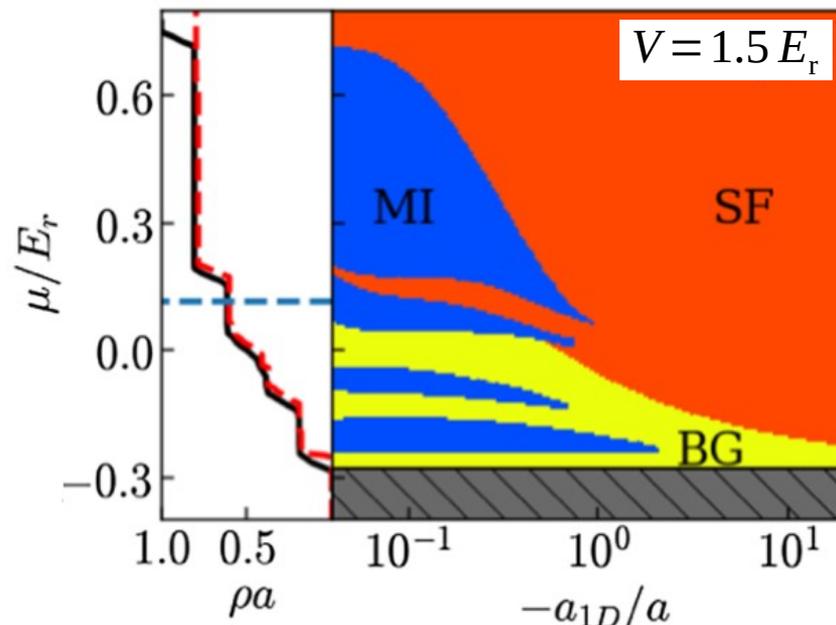
SF

(Compressible)  
Superfluid

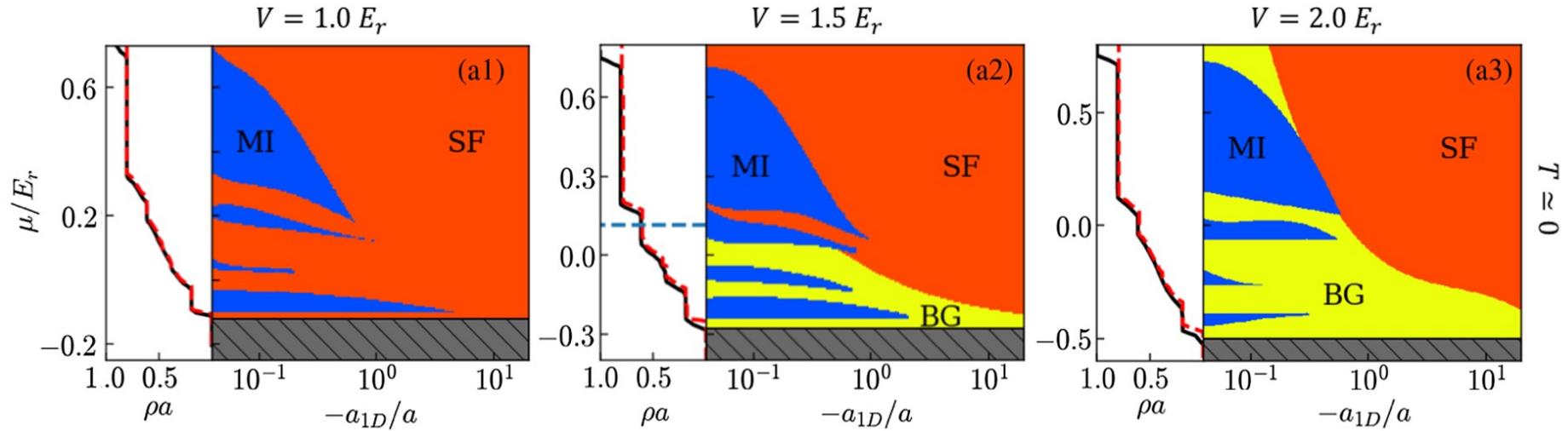


BG

Compressible  
Non superfluid



# Lieb-Liniger Gas in a Shallow Quasiperiodic Lattice

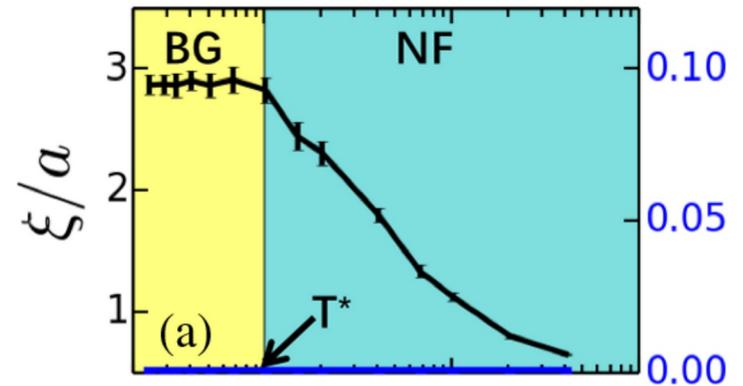


What about finite temperatures?

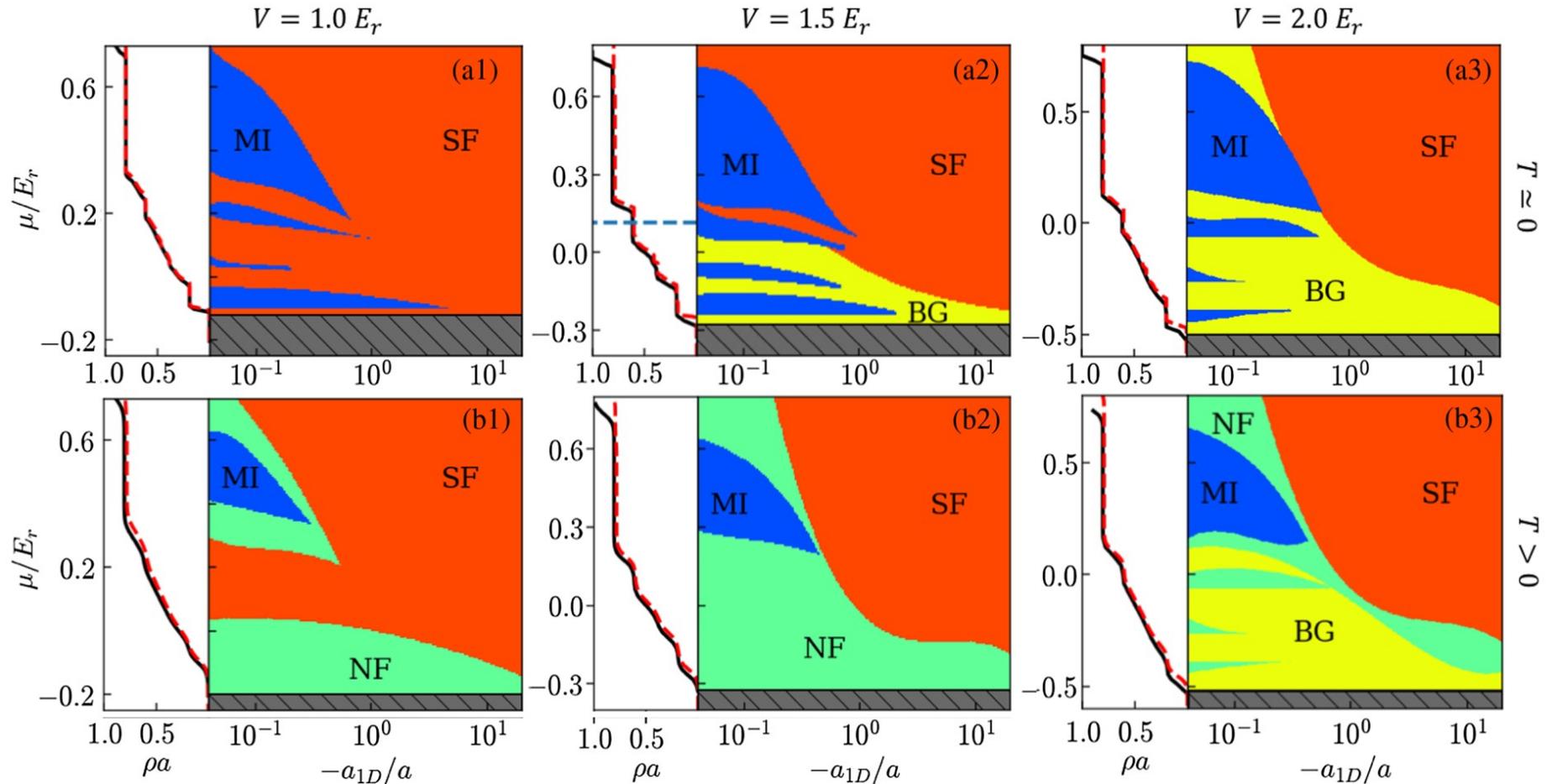
A new regime : the normal fluid

It is a compressible insulator (just as the Bose glass)

→ Bose glass as long as no sizable finite-temperature effects are visible



# Lieb-Liniger Gas in a Shallow Quasiperiodic Lattice



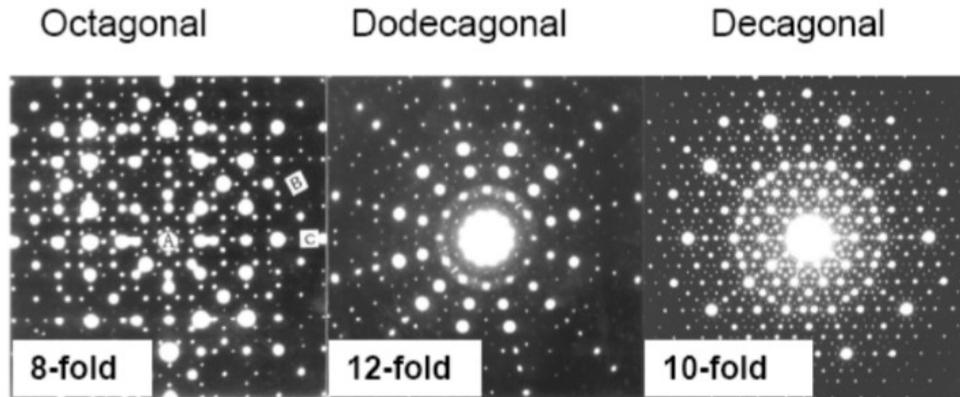
# Optical Quasicrystals : From 1D to 2D

## Quantum simulation of quasicrystals

Solids with long-range but nonperiodic order

Discovered by chance using Bragg spectroscopy

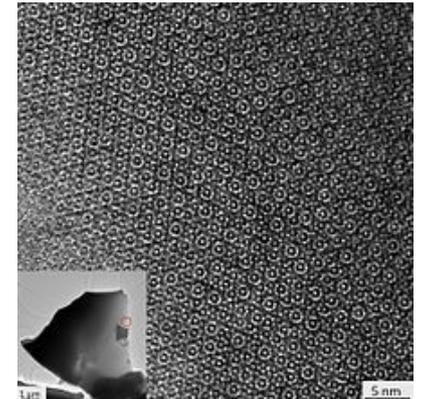
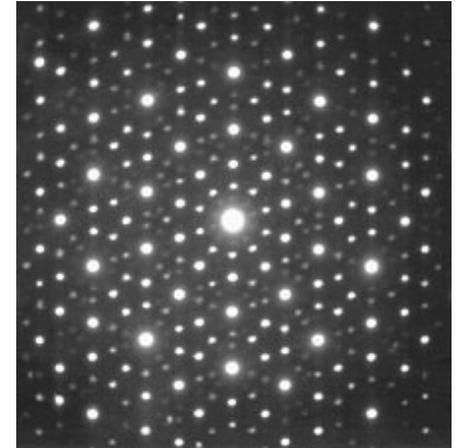
Many intriguing physical properties



Lab synthesis of artificial quasicrystals

→ From correlated electrons to bosonic matter

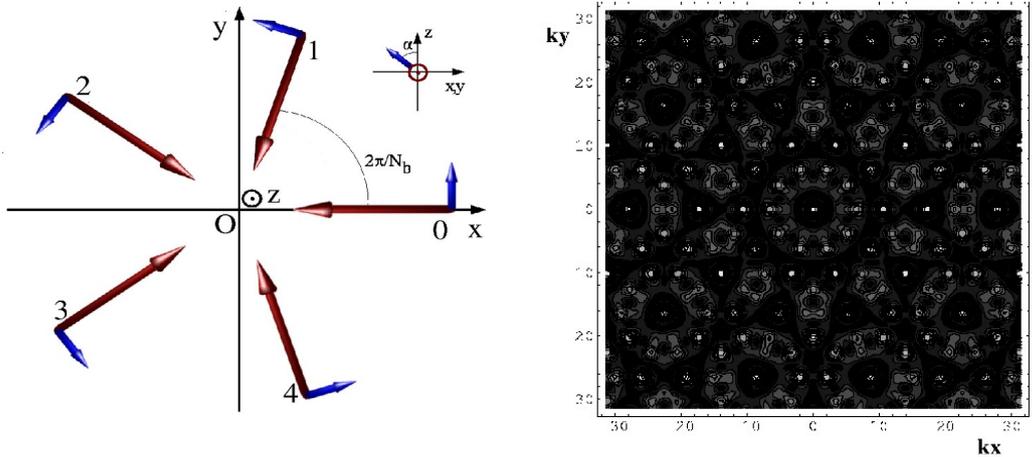
Bragg spectroscopy :  
long-range order  
but nonperiodic



$\text{Al}_{71}\text{Ni}_{24}\text{Fe}_5$  quasicrystal  
in a Khatyrka meteorite

# Optical Quasicrystal and Emergence of Quasiperiodic Order

Sanchez-Palencia & Santos, Phys. Rev. A 72, 053607 (2005)

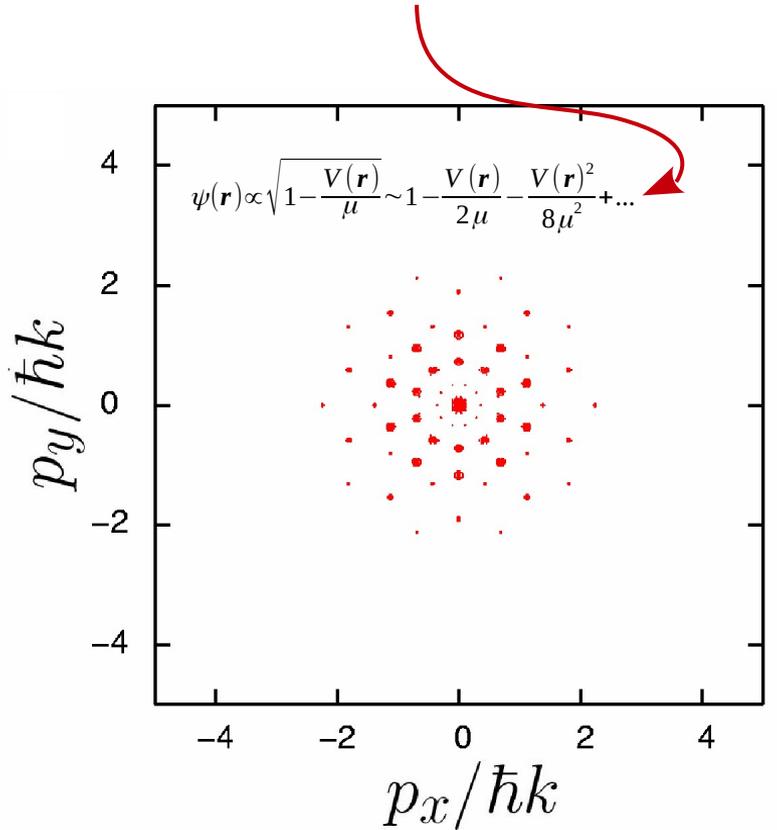


A potential with 10-fold rotational symmetry

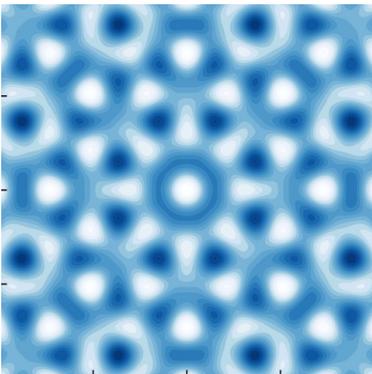
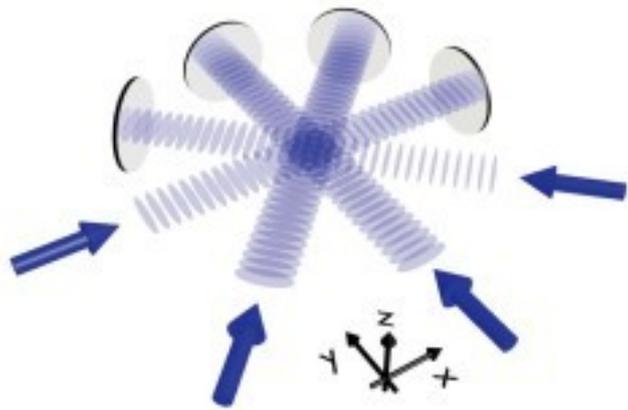
$$V(\mathbf{r}) = V_0 \left| \sum_{1 \leq j \leq 5} e^{i \mathbf{k}_j \cdot \mathbf{r}} \right|^2$$

The condensate contains all harmonics, owing to nonlinearity (interactions)

$$\mu \psi(\mathbf{r}) = \frac{-\hbar^2 \nabla^2}{2m} \psi(\mathbf{r}) + V(\mathbf{r}) \psi(\mathbf{r}) + g |\psi(\mathbf{r})|^2 \psi(\mathbf{r})$$

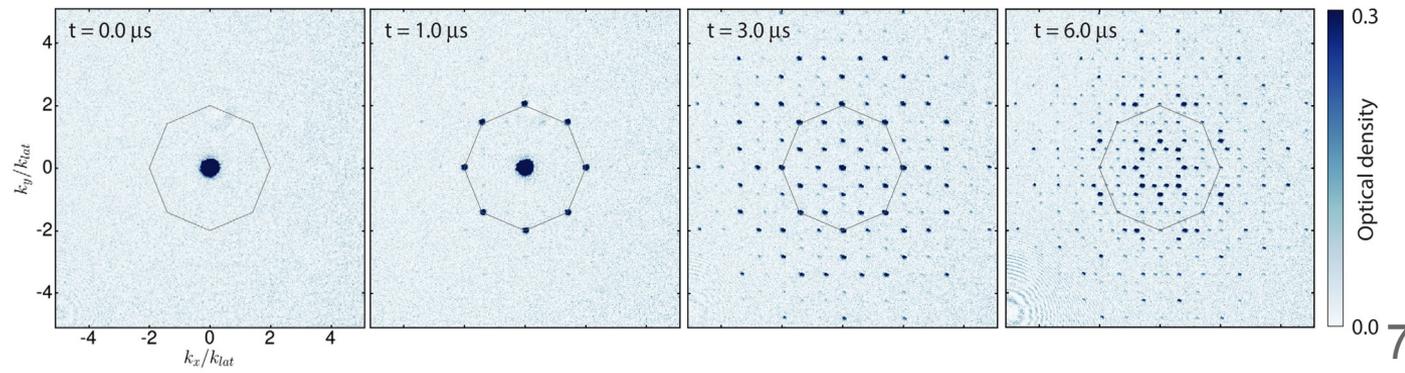


# Emerging Quasiperiodic Order from Ultracold Matterwaves



Jagannathan *et al.*, Europhys. Lett. **104**, 66003 (2013)  
Schneider's group, Phys. Rev. Lett. **122**, 110404 (2019)

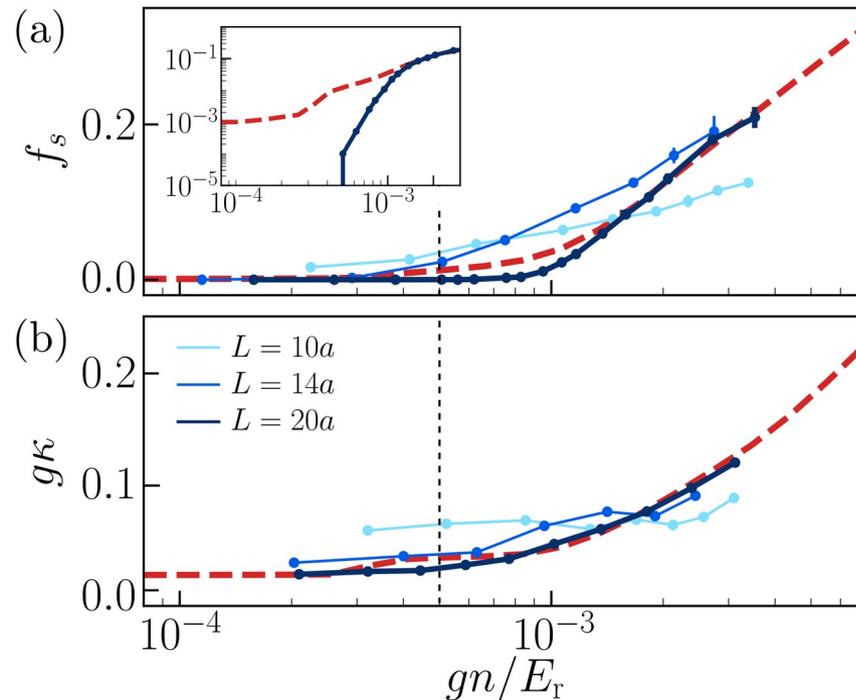
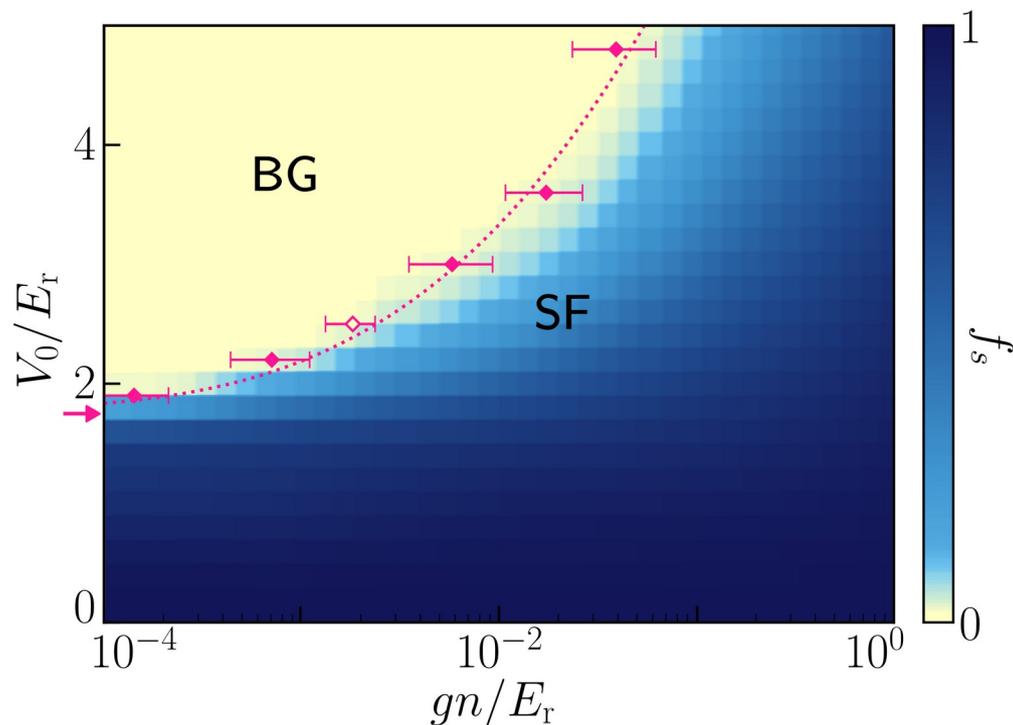
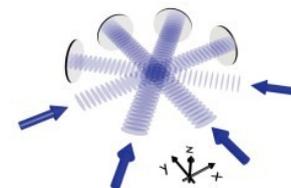
Kapitza-Dirac diffraction  
Quantum random walk in momentum space



# Weakly-Interacting Bosons in a 2D Quasicrystal Lattice

R. Gautier, H. Yao & L. Sanchez-Palencia, Phys. Rev. Lett. **126**, 110401 (2021)

Combining meanfield (Gross-Pitaevskii) and quantum Monte Carlo calculations



# Weakly-Interacting Bosons in a 2D Quasicrystal Lattice

## Coherence measurement

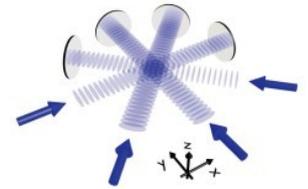
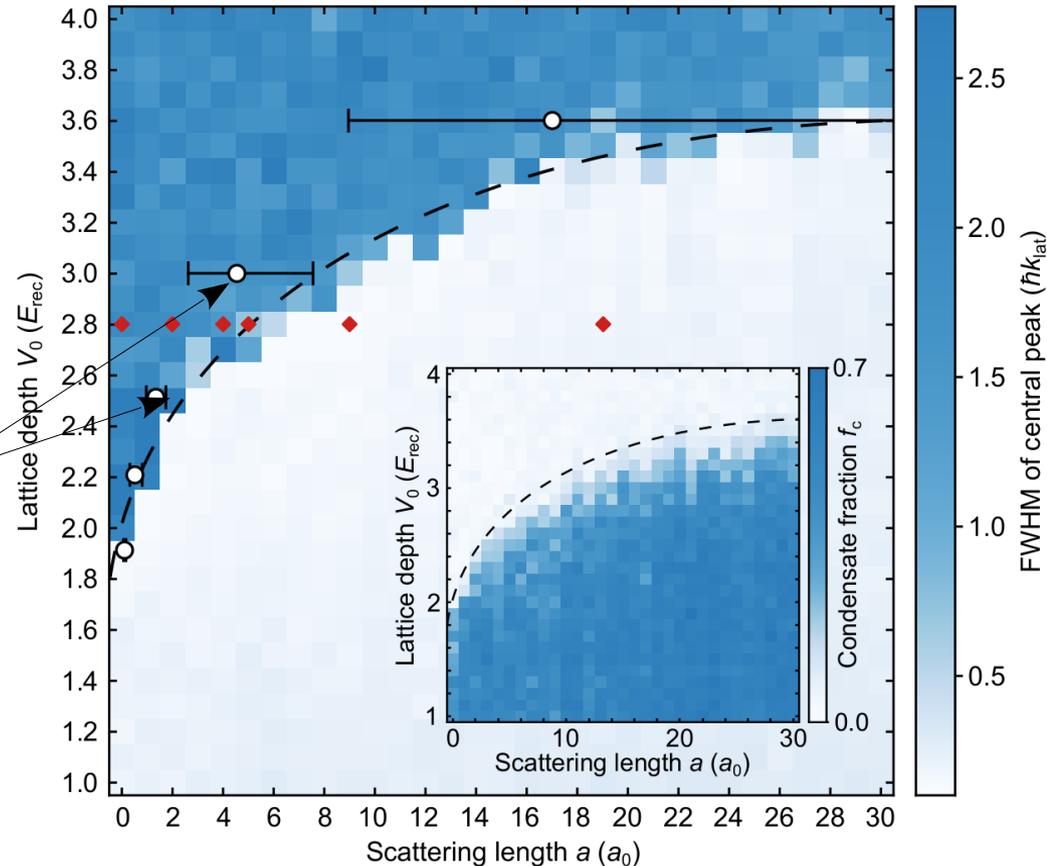
Schneider's group

[arXiv:2303.00737 (2023)]

## Theoretical prediction

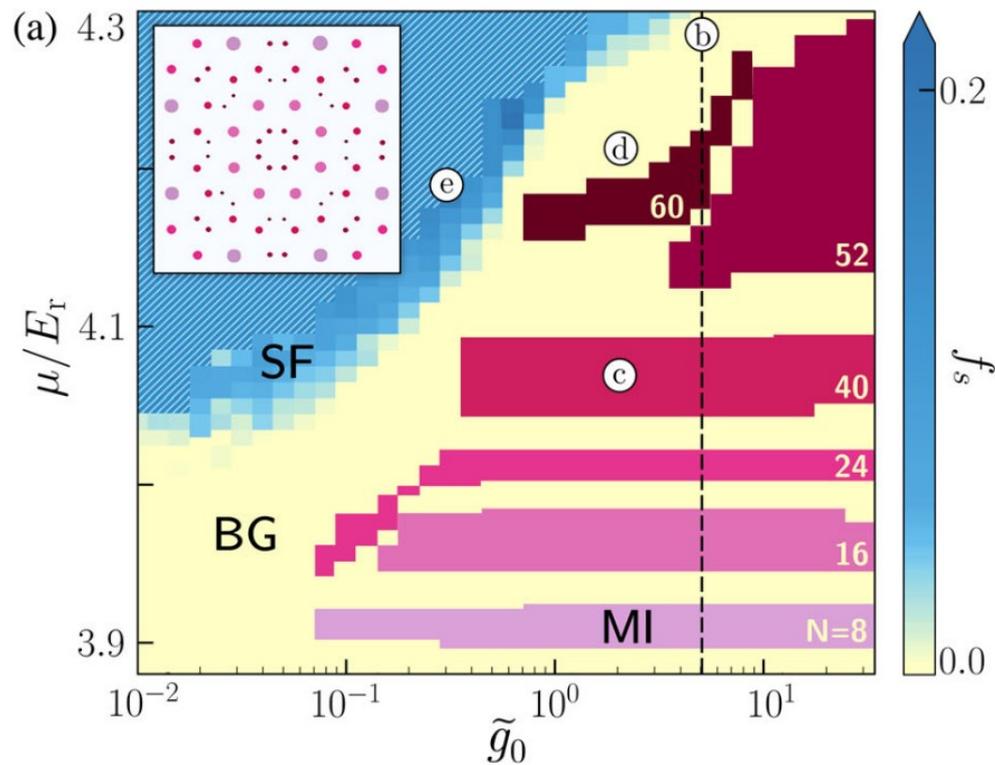
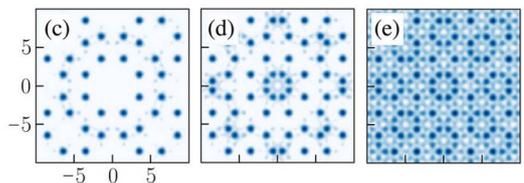
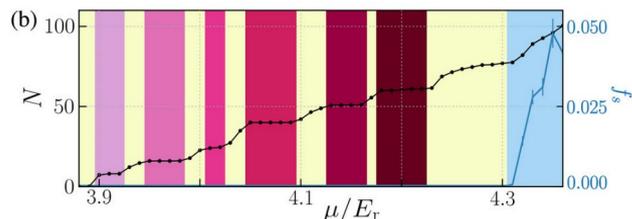
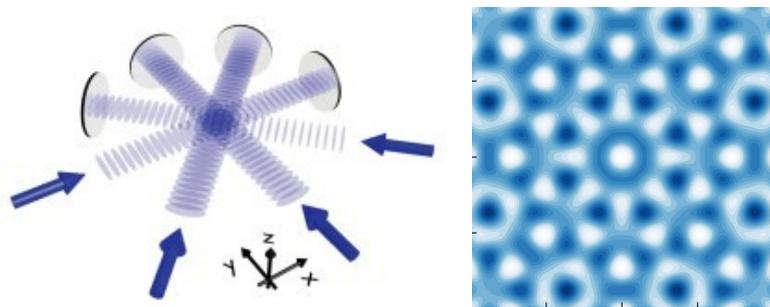
Gautier *et al.*,

Phys. Rev. Lett. **126**, 110401 (2021)



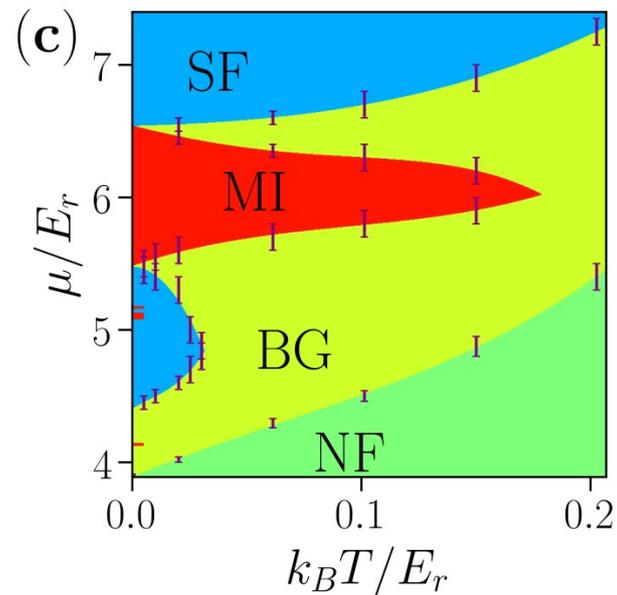
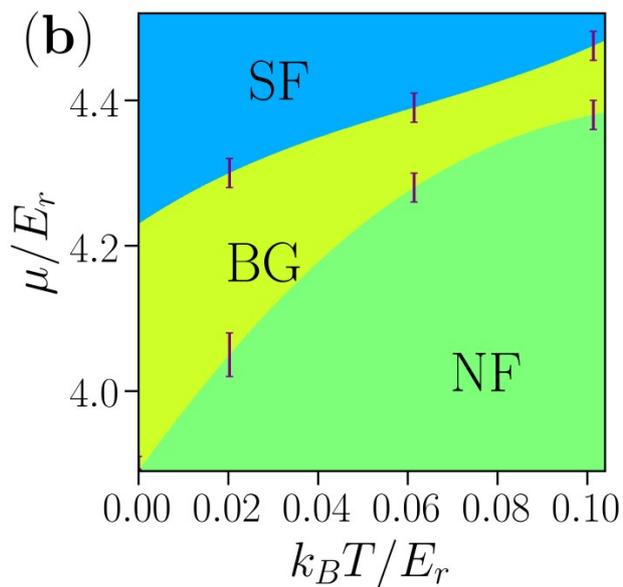
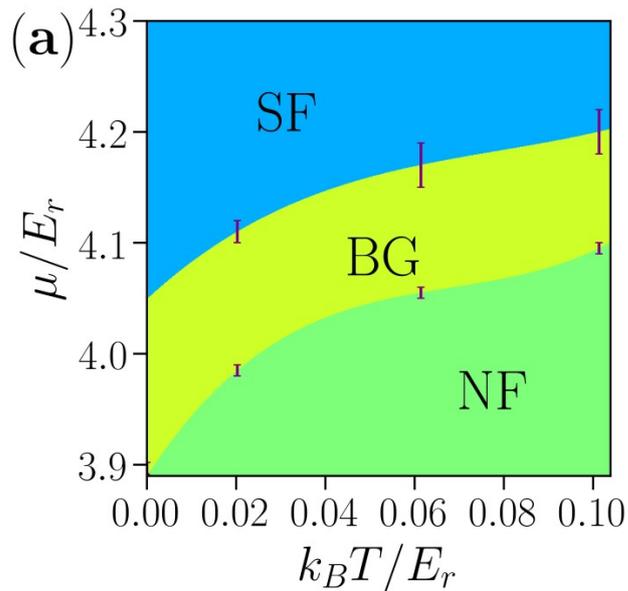
# Strongly-Interacting Bosons in a 2D Quasicrystal Lattice

R. Gautier, H. Yao & L. Sanchez-Palencia, Phys. Rev. Lett. **126**, 110401 (2021)



# Strongly-Interacting Bosons in a 2D Quasicrystal Lattice

Z. Zhu, H. Yao & L. Sanchez-Palencia, Phys. Rev. Lett. (2023)



Quantum phases survive to thermal fluctuations, up to significant temperatures

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Lih-King Lim

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Luca Pezzé

Marie Piraud

Pierre Lugan

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Steven J. Thomson

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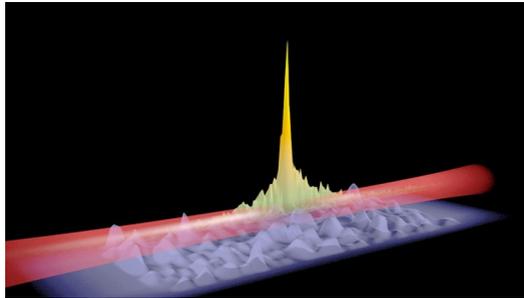
Kim Vallée

Louis Villa

Hepeng Yao



# Some of our (Recent) Activities



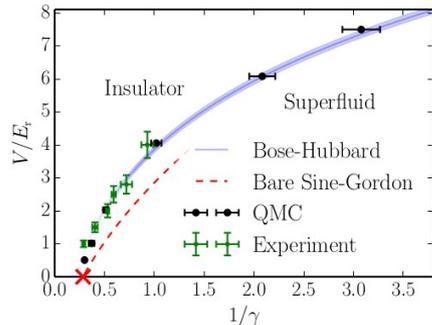
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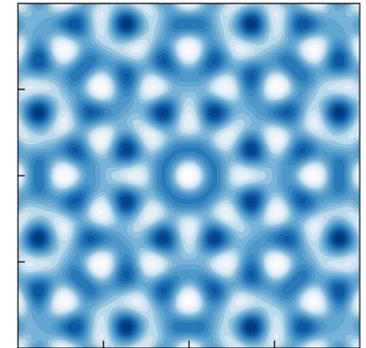
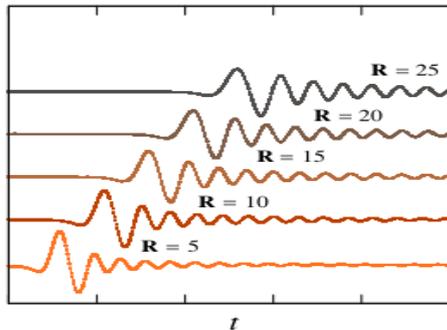


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