

Can we detect deep axisymmetric toroidal magnetic fields in stars?

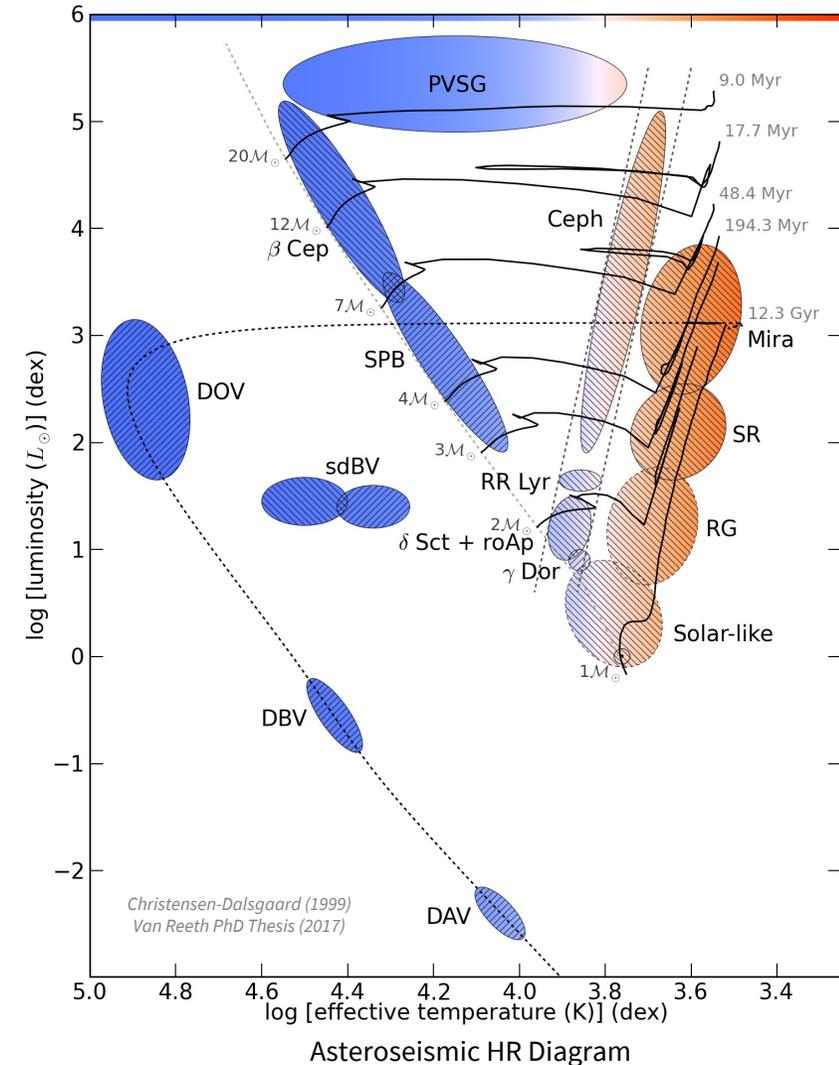
Hachem Dhouib

In collaboration with: S. Mathis, L. Bugnet, T. Van Reeth, C. Aerts

Importance of angular momentum transport in stars

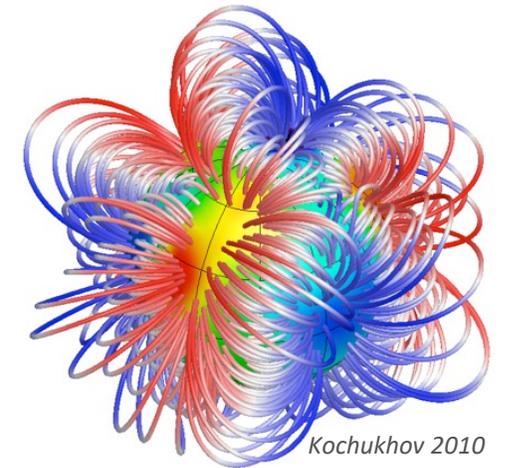
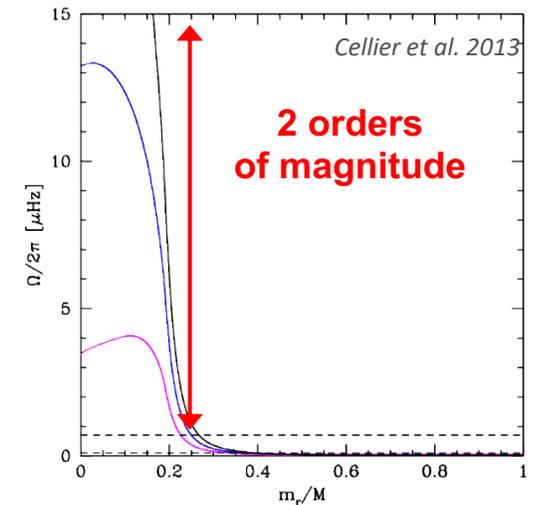
- ❖ Stars: **fundamental units** of the visible Universe
 - Stellar structure and evolution: **cornerstone of Astrophysics**
- **Goal:**
 - Understand stars interiors, dynamics, and evolution
- **Key element:**
 - Transport of angular momentum
- **Why?**
 - Angular momentum transport impacts:
 - ✓ rotation
 - ✓ chemical mixing
 - ✓ evolution
 - ✓ interactions with their environment
 - ✓ magnetism

Angular momentum transport shapes the evolution and the dynamics of stars



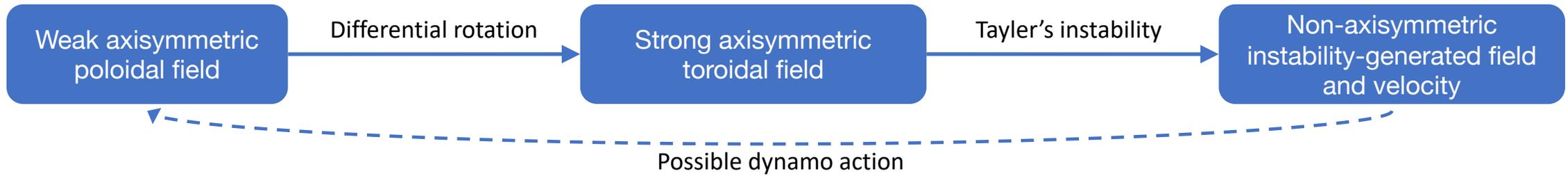
Magnetic fields in radiative zones of early-type stars

- ❑ Current angular momentum transport models unable to reproduce asteroseismic observations (Aerts et al. 2019)
- ❑ One of the best candidates to explain this discrepancy: **Magnetic fields** in radiative zones
- ❑ 2 types of detected magnetic fields at stellar surfaces using spectropolarimetry:
 - Large-scale stable fields with high amplitude
 - ✓ Origin: fossil (Braithwaite & Spruit 2004, Duez & Mathis 2010, Shultz et al. 2019)
 - Small-scale fields with low amplitude
 - ✓ Origin: dynamo action in the thin sub-surface convective layer (Cantiello & Braithwaite 2019) / failed relaxation (Braithwaite & Cantiello 2013) / resulting from non-axisymmetric instabilities (Aurière et al. 2007)

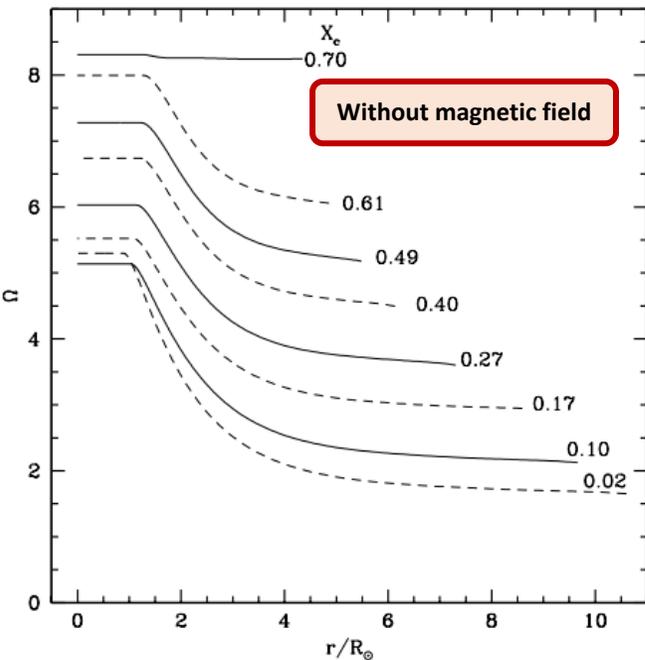


Strong axisymmetric toroidal magnetic field

(Pitts & Tayler 1985, Spruit 2002, Zahn et al. 2007, Fuller et al. 2019, Petitdemange et al. submitted)

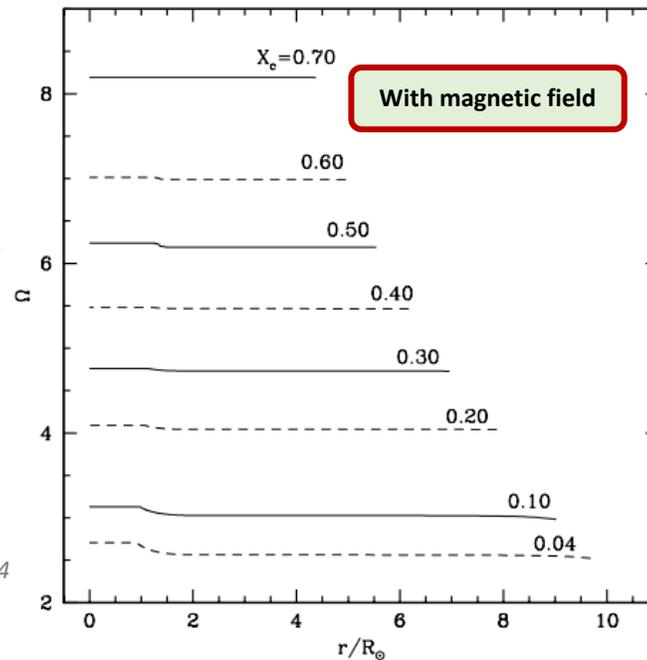


➤ Torque exerted by induced Maxwell stresses: very efficient transport of angular momentum



Enforce the solid-body rotation of massive MS stars

Maeder & Meynet 2004

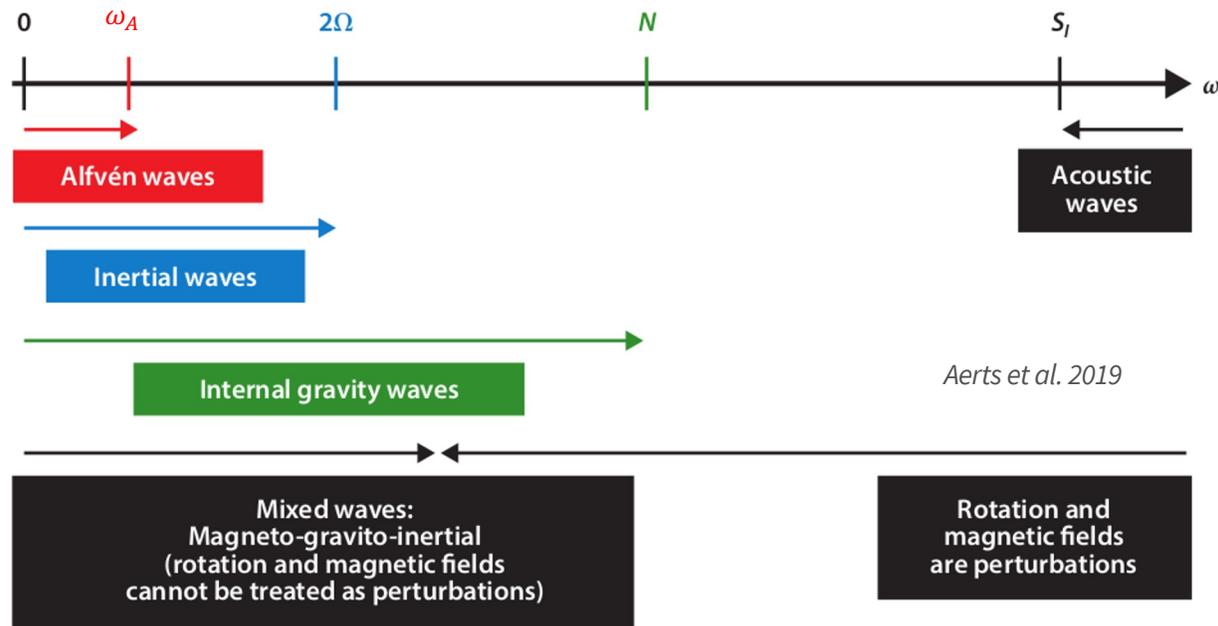


Is it possible to detect this magnetic field?

This magnetic field does not emerge at the stellar surface
 ➔ **spectropolarimetry is blind**

Only window: magneto-asteroseismology
 ➔ look for characteristic magnetic signatures in stellar oscillations

g-mode pulsators and magnetic field



Asteroseismology: only way to probe the internal **structural, chemical, rotational and magnetic properties**

Goal: study the signature of **axisymmetric toroidal magnetic field**

- ❖ Magneto- Gravito-inertial (MGI) modes driven by:
 - ✓ buoyancy: chemical and thermal stratifications (N)
 - ✓ rotation: Coriolis and centrifugal accelerations (2Ω)
 - ✓ magnetic field: Lorentz force (ω_A)

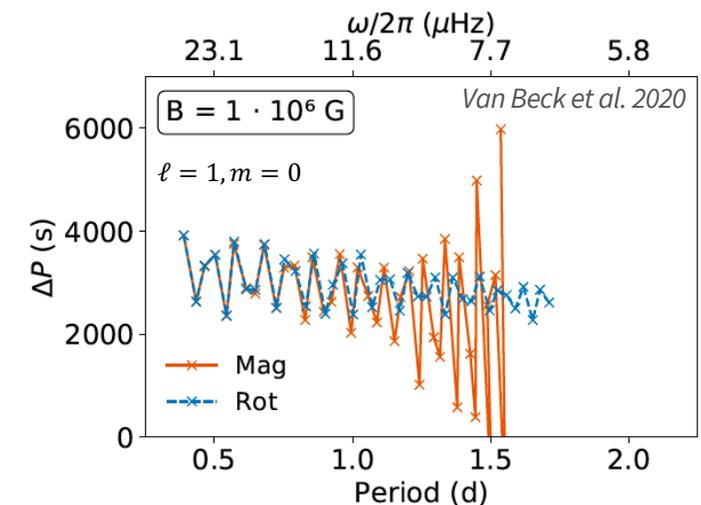
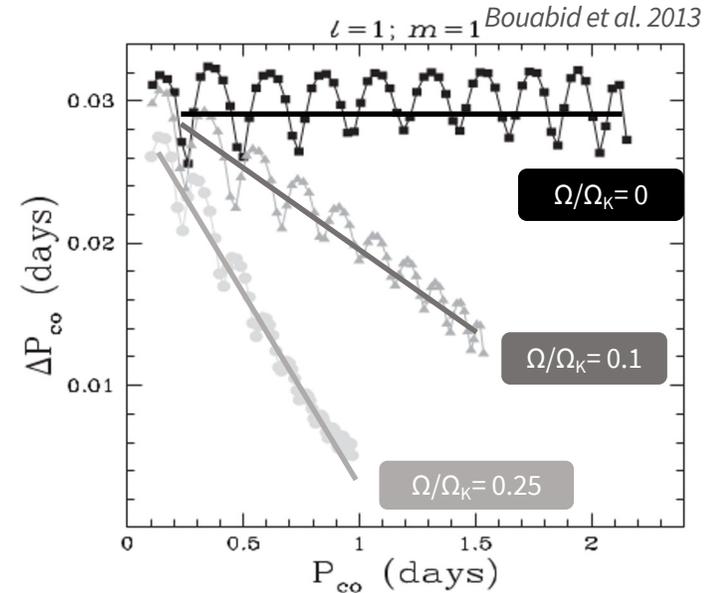
- ❖ Acoustic modes driven by:
 - ✓ Compressibility: pressure force (S_ℓ)

Period spacing pattern: a window into stars' interiors

$$\Delta P = P_{n+1} - P_n: \text{period spacing}$$

- ❑ No rotation → Constant
- ❑ Rotation → Slope
 - Measurement of the slope ⇒ measure of Ω (Bouabid et al. 2013, Van Reeth et al. 2016, Ouazzani et al. 2017)
- ❑ Chemical composition (Degroote et al. 2010, Pedersen et al. 2021) / Differential rotation / Coupling between inertial modes in the convective core and g modes (Ouazzani et al. 2020, Saio et al. 2021) → Dips
- ❑ Mixed magnetic field (poloidal + toroidal) (Prat et al. 2019,2020, Van beck et al. 2020) → Sawtooth pattern
 - ❖ **Perturbative approach**

Derive a non-perturbative formalism to study the impact of strong general axisymmetric toroidal magnetic field



Traditional Approximation of Rotation (TAR)

- Goal: Compute the pulsations of MGI modes
- Issue: Difficulty of performing flexible and intensive seismic modelling with 2D oscillation codes (e.g. Reese et al. 2006)
- Possible solution: The Traditional Approximation of Rotation (TAR) (e.g. Eckart 1960, Bildsten et al. 1996, Lee & Saio 1997)
 - Strengths:
 - ✓ separable dynamics
 - ✓ tractable formalism / efficient for extensive seismic modelling
 - ✓ robust predictions for frequencies when compared to 2D adiabatic modes computation
 - Assumptions:
 - ✓ Strong stratification: $2\Omega \ll N$
 - ✓ Low frequency: $\omega \ll N$
 - Neglect the vertical (horizontal) component of the Coriolis acceleration (rotation vector)
 - ✓ adiabatic limit / Cowling approximation
 - ✓ spherical stars / uniform rotation (e.g. Lee & Saio 1997)
 - ✓ **no magnetic field**
 - Weaknesses:
 - ✓ Can't be used to deal with waves' excitation, damping, and instabilities (e.g. Mathis et al. 2014)
 - ✓ Not applicable in convective regions and in very rapidly rotating stars (e.g. Ogilvie & Lin 2004, Dhouib et al. 2021a,b)

Magnetic TAR

❖ Magnetic Laplace tidal equation (Horizontal structure) :

Assuming: $\omega_A \ll N$

$$\mathcal{L}_{\omega^{\text{in}} m}^{\text{magn}} = \omega^2 \partial_x \left[\frac{1}{\mathcal{A}} \frac{1-x^2}{D_M} \partial_x \right] + m\omega^2 \partial_x \left(\frac{\nu_M x}{\mathcal{A} D_M} \right) - m^2 \frac{\omega^2}{\mathcal{A} D_M (1-x^2)} + m^2 \frac{\omega^2}{\mathcal{A}^2} \frac{x}{D_M} \partial_x \omega_A^2$$

$$\mathcal{L}_{\omega^{\text{in}} m}^{\text{magn.}} \left[w_{\omega^{\text{in}} km}^{\text{magn.}}(r, x) \right] = -\Lambda_{\omega^{\text{in}} km}^{\text{magn.}}(r) w_{\omega^{\text{in}} km}^{\text{magn.}}(r, x)$$

Eigenvalues

Generalised Hough functions

$$\begin{aligned} x &= \cos \theta \\ \mathcal{A} &= \omega^2 - m^2 \omega_A^2 \\ \mathcal{B} &= 2(\Omega \omega + m \omega_A^2) \\ \nu_M &= \mathcal{B} / \mathcal{A} \\ D_M &= 1 - \nu_M^2 x^2 - (1-x^2) \frac{x}{\mathcal{A}} \partial_x \omega_A^2 \end{aligned}$$

Magnetic spin parameter

Alfvén frequency

❖ Dispersion relation + Radial quantification :

$$\int_{r_1}^{r_2} \frac{N \sqrt{\Lambda_{\omega^{\text{in}} km}^{\text{magn.}}}}{r \omega} dr = (n + 1/2) \pi$$

(2D JWKB approximation : Rapidly oscillating waves along the radial direction)

m : azimuthal order

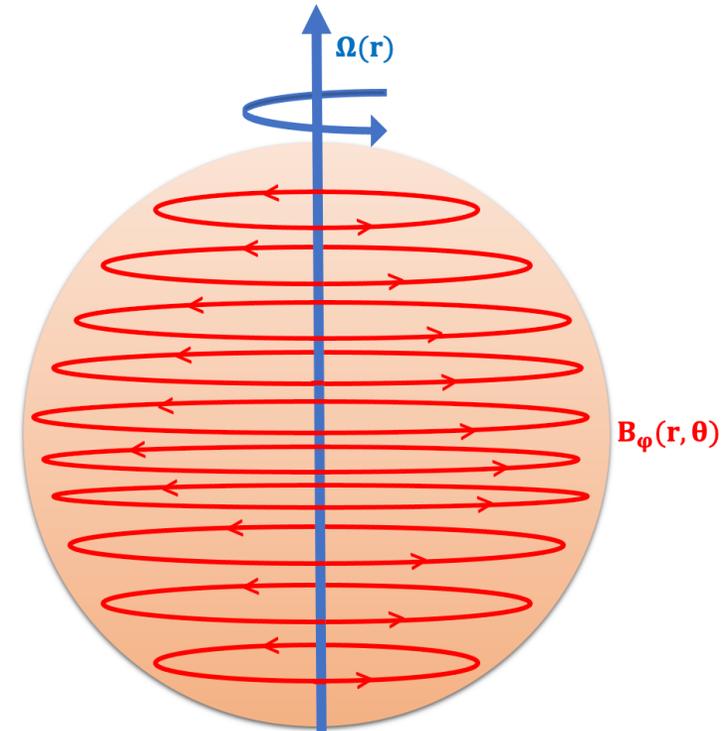
k : latitudinal index ($k = \ell - |m|$)

n : radial order

ℓ : latitudinal order

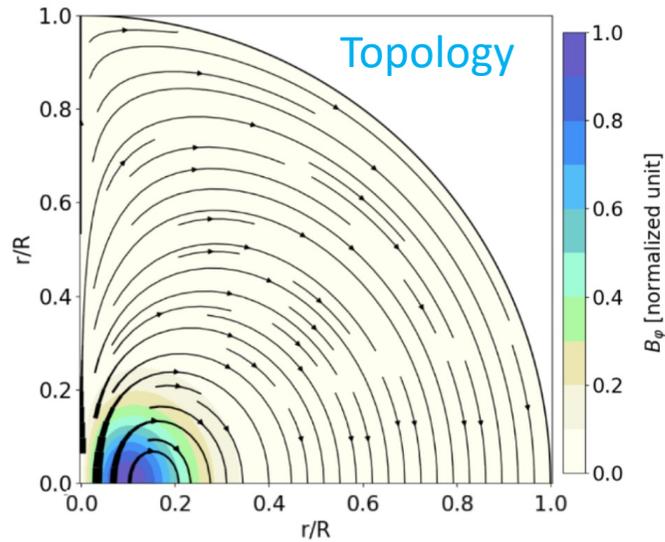
Setup:

- Differential rotation
- General axisymmetric toroidal field



$$B_\phi(r, \theta) = \sqrt{\mu \rho_0} r \sin \theta \omega_A(r, \theta)$$

Case of equatorial toroidal field (1/2)



Initially fossil field

Sufficient radial differential rotation \rightarrow efficient dynamo loop can take place

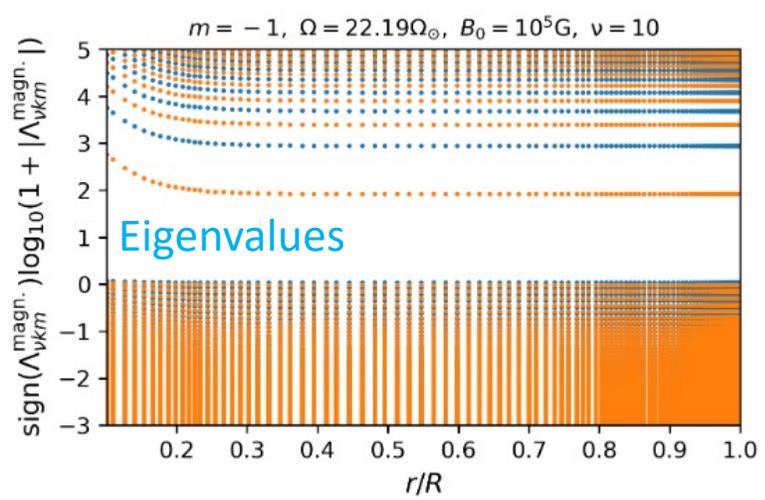
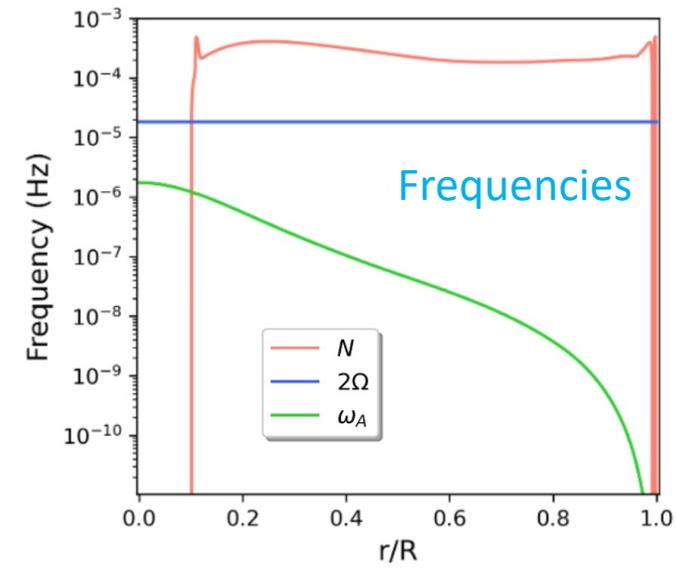
γ Dor model

$1.6 M_\odot$

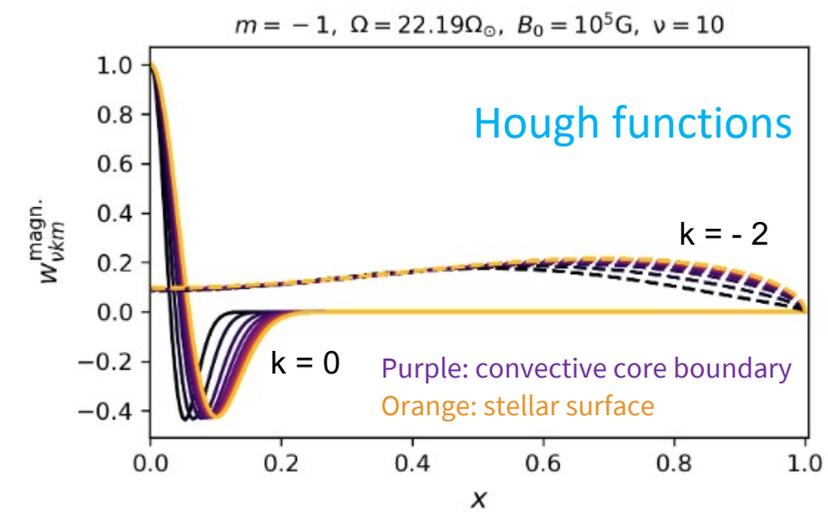
near ZAMS

$B_0 = 10^5 \text{ G}$

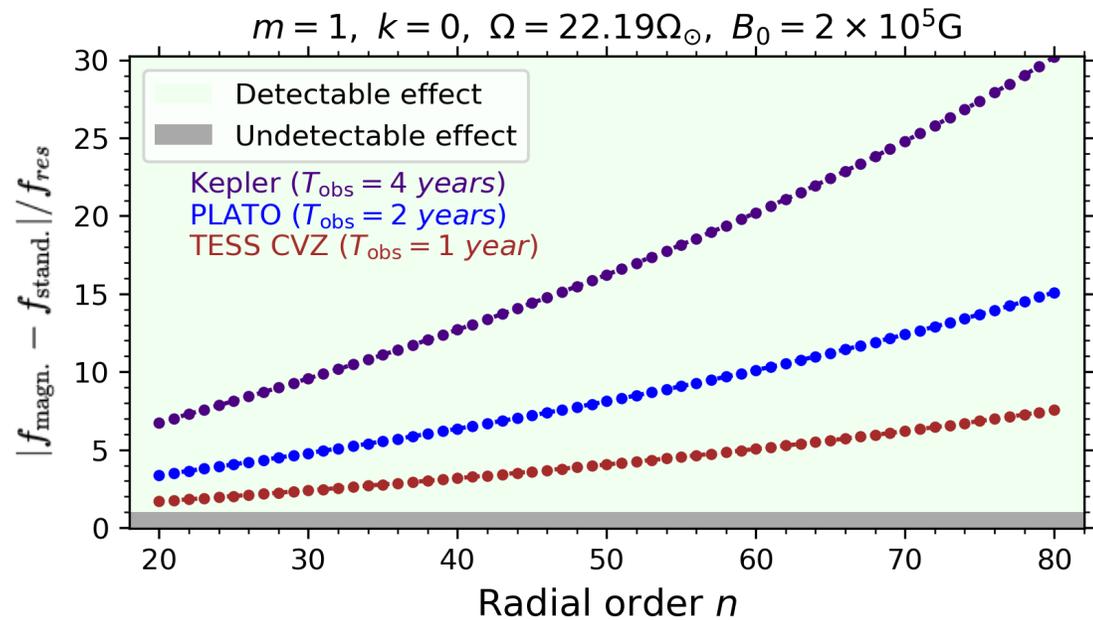
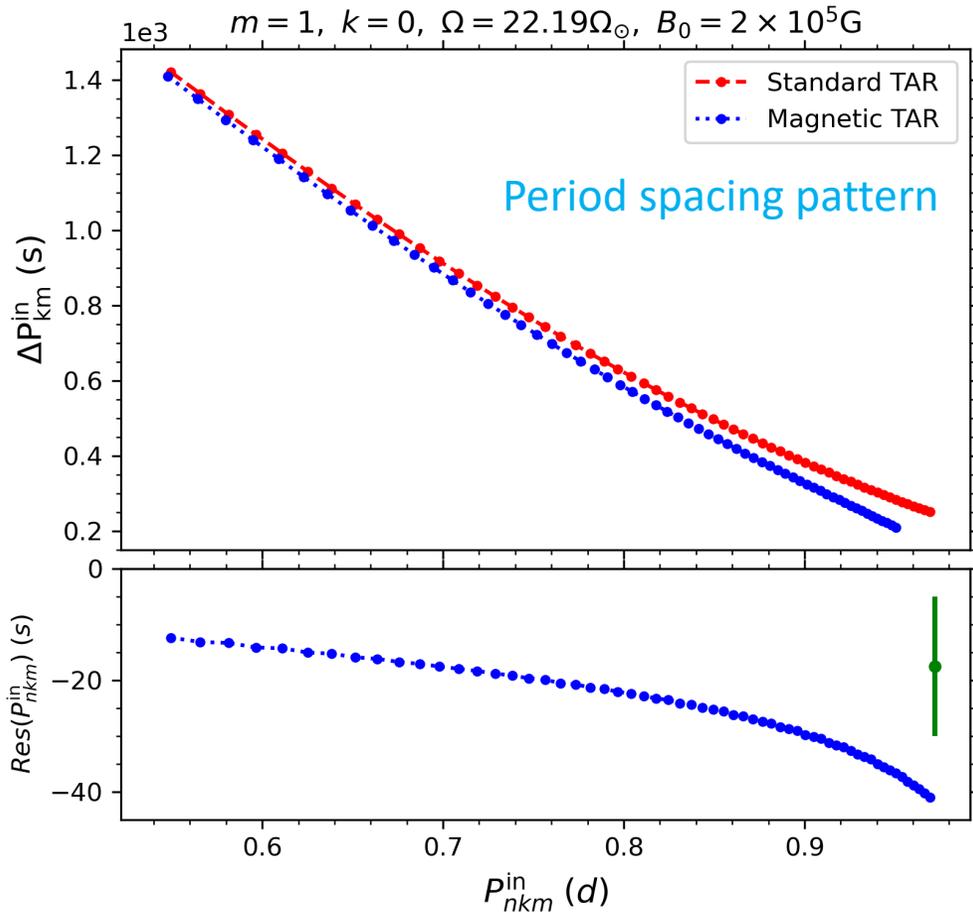
$\Omega = 22.19 \Omega_\odot$



\rightarrow Strong radial dependence in the inner region of the star

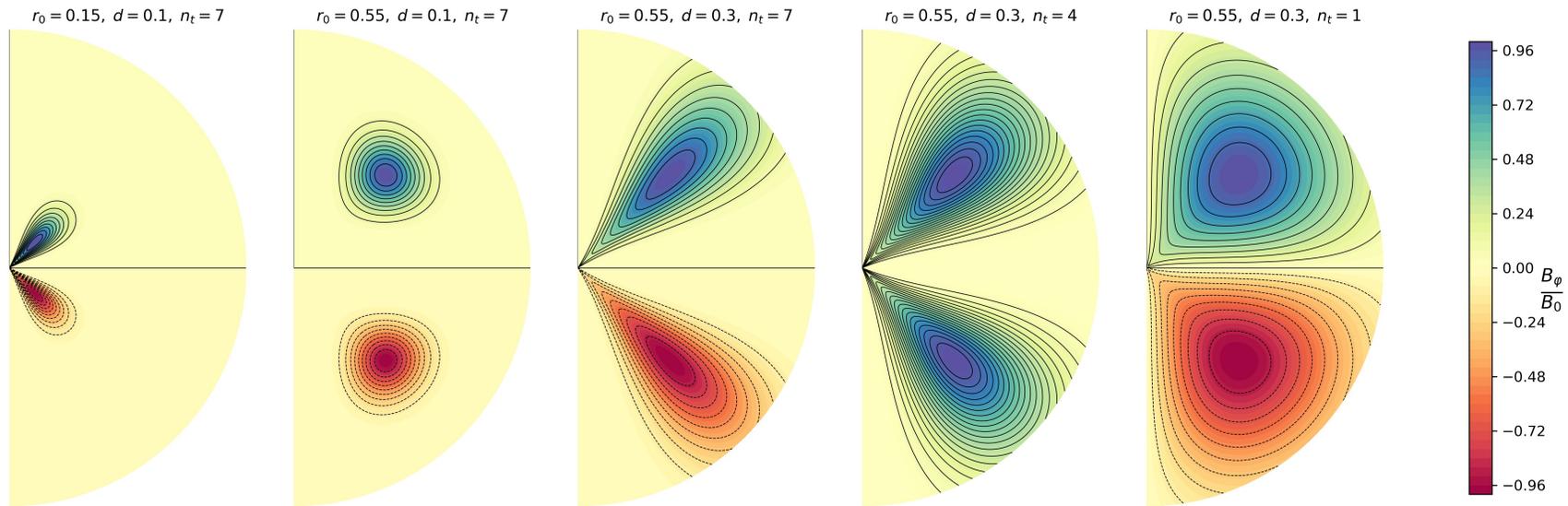


Case of equatorial toroidal field (2/2)



- A deep buried equatorial toroidal field can be detected
- ❖ Signature similar to rotation
- ❖ No sawtooth pattern seen with perturbative method
(Prat et al. 2019, 2020, Van Beeck et al. 2020)
- ✓ Possible cause:
 - absence of poloidal component?
 - asymptotic approach?

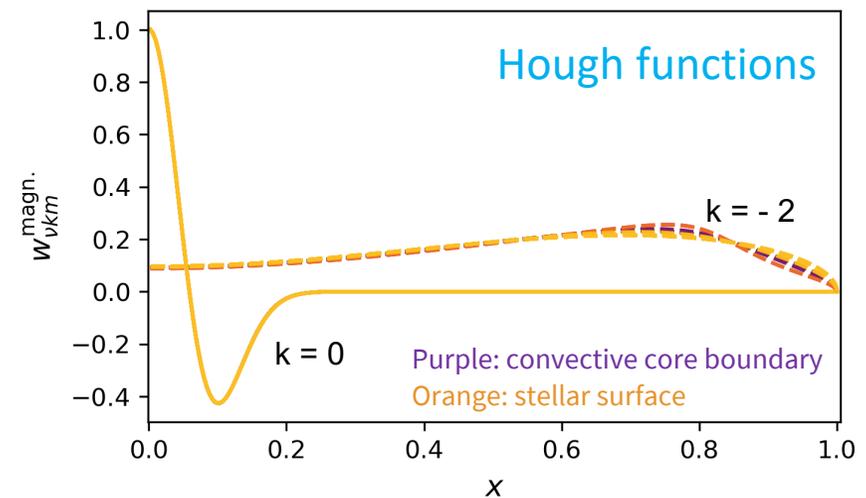
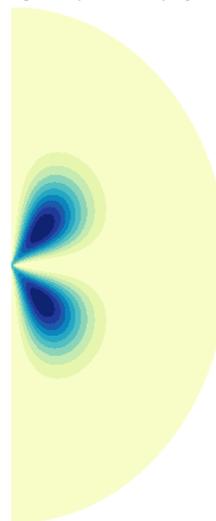
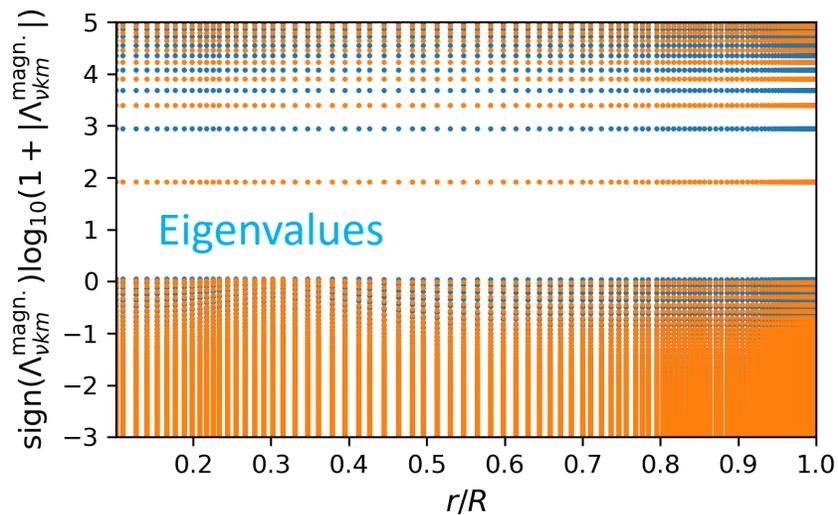
Case of hemispheric toroidal field (1/2)



$$B_\phi(r, \theta) = B_0 \exp\left[-\frac{(r - r_0)^2}{2d^2}\right] \sin^{n_t}(2\theta)$$

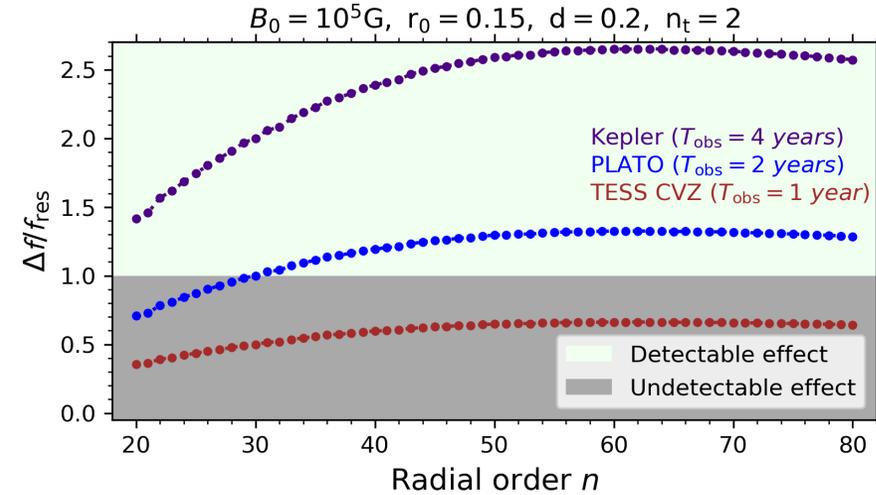
Depend on the history of the differential rotation (*Jouve et al. 2020*)

$m = -1, B_0 = 10^5 \text{G}, r_0 = 0.2, d = 0.15, n_t = 2, \nu = 10$



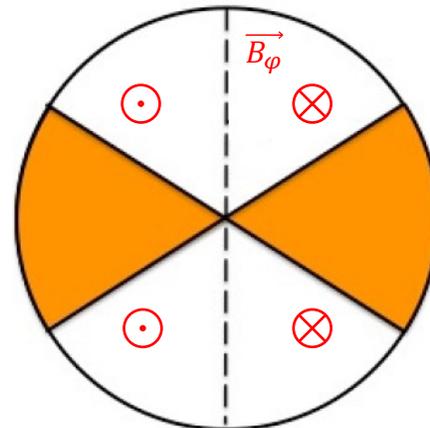
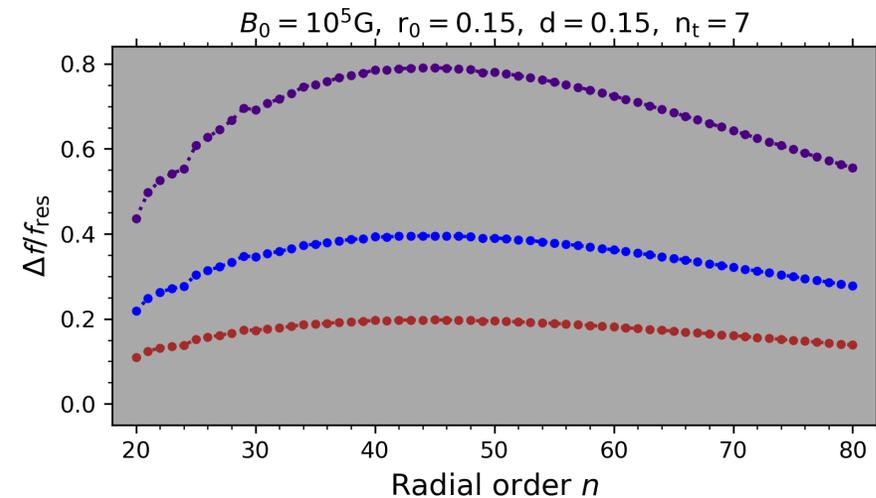
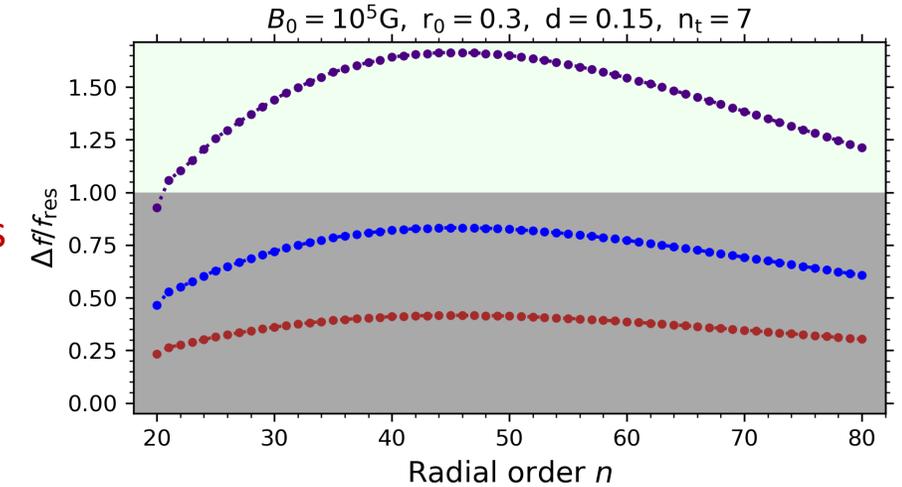
➤ Weak radius dependence

Case of hemispheric toroidal field (2/2)

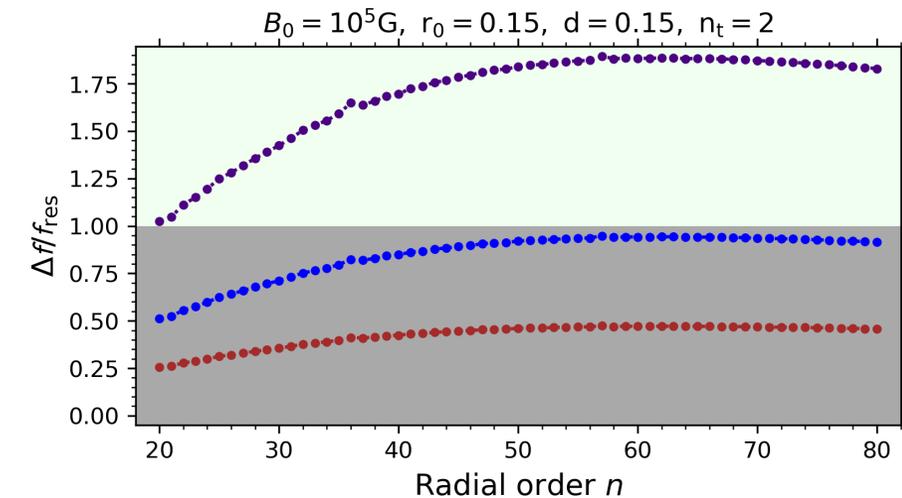


➤ Deep buried hemispheric fields difficult to detect

❖ MGI modes propagate in an equatorial belt

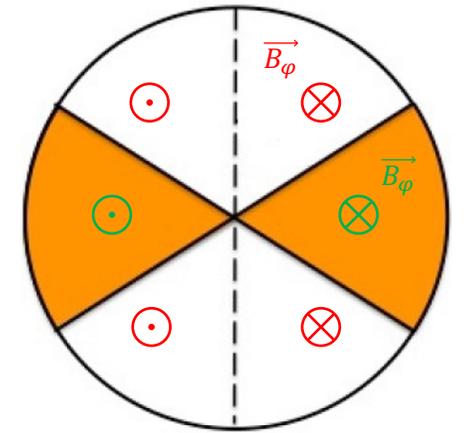


Mathis & de Brye 2012



Conclusions and perspectives

- ✓ Derivation of a **new generalisation of the TAR**
- ✓ Better **magneto-asteroseismic modelling**
- ✓ **Dependency** of the detectability of magnetic fields **on their configurations**
- ✓ **Possible detection** of equatorial fields
- ✓ Signature **similar to rotation**
- ✓ **Greater difficulty** in detecting fields located far from the equator



- ❖ Apply the magnetic TAR to **real stars** → improve **forward modelling?**
- ❖ Go beyond the asymptotic approach: **incorporate the magnetic TAR in GYRE**
- ❖ Go beyond the TAR: **3D MHD oscillation code**